



A simple kinetic equation of swarm formation: Blow-up and global existence



Mirosław Lachowicz^{a,b,*}, Henryk Leszczyński^c, Martin Parisot^{d,e,f,g}

^a Institute of Applied Mathematics and Mechanics, Faculty of Mathematics, Informatics and Mechanics, University of Warsaw, ul. Banacha 2, 02-097 Warsaw, Poland

^b School of Mathematics, Statistics and Computer Science, University of KwaZulu-Natal, South Africa

^c Institute of Mathematics, University of Gdańsk, ul. Wita Stwosza 57, 80-952 Gdańsk, Poland

^d INRIA, ANGE Project-Team, Rocquencourt, F-78153 Le Chesnay Cedex, France

^e CEREMA, F-60280 Margny-Lès-Compiègne, France

^f CNRS, UMR 7598, Laboratoire Jacques-Louis Lions, F-75005, Paris, France

^g Sorbonne Universités, UPMC Univ Paris 06, UMR 7598, Laboratoire Jacques-Louis Lions, F-75005, Paris, France

ARTICLE INFO

Article history:

Received 1 December 2015

Received in revised form 18 January 2016

Accepted 19 January 2016

Available online 27 January 2016

Keywords:

Blow-up

Global existence

Kinetic equation

ABSTRACT

In the present paper we identify both blow-up and global existence behaviors for a simple but very rich kinetic equation describing of a swarm formation.

© 2016 Elsevier Ltd. All rights reserved.

1. Introduction

In paper [1] a model of swarming behavior of an individual population was proposed and studied. The main aim was the macroscopic (*hydrodynamic*) limit. The mathematical structure that was proposed seems very reach and interesting from mathematical point of view. Let $f = f(t, x, v)$ be a probability density (p.d.) of individuals at time $t \geq 0$ and position $x \in \mathbb{R}^d$ with velocity $v \in \mathbb{V}$; $\mathbb{V} \subset \mathbb{R}^d$, the set of velocities of the individuals, is a bounded domain. The evolution of populations at the mesoscopic scale is defined by the

* Corresponding author at: Institute of Applied Mathematics and Mechanics, Faculty of Mathematics, Informatics and Mechanics, University of Warsaw, ul. Banacha 2, 02-097 Warsaw, Poland.

E-mail address: lachowic@mimuw.edu.pl (M. Lachowicz).

nonlinear integro-differential Boltzmann-like equation, see [1],

$$\begin{aligned} \partial_t f(t, x, v) + v \cdot \nabla_x f(t, x, v) &= \frac{1}{\varepsilon} Q[f](t, x, v) \\ &= \frac{1}{\varepsilon} \int_{\mathbb{V}} \left(T[f(t, x, \cdot)](w, v) f(t, x, w) - T[f(t, x, \cdot)](v, w) f(t, x, v) \right) dw \end{aligned} \tag{1}$$

with the initial data $f(0, x, v) = f_0(x, v)$. The parameter ε corresponds to the Knudsen number and the macroscopic limit is defined by $\varepsilon \rightarrow 0$. The nonlinear operator Q describes interactions between individuals. The turning rate $T[f](v, w)$ measures the probability for an individual with velocity v to change velocity into w . Macroscopic limit for a simpler (two-velocities) kinetic equation was studied in [2] (see also [3]). In the context of modeling of preferential choice one should mention the paper [4] where a collision model was proposed describing ant trail formation.

In Ref. [1] the following general nonlinear case

$$T[f(t, x, \cdot)](v, w) = \sigma \beta(v, w) f^{\gamma\rho, x}(t, x, w), \tag{2}$$

was considered, where the interaction rate β , the attractiveness coefficient γ , and σ characterize the interaction between the individual agents. Parisot and Lachowicz [1] proposed results of global existence in the space homogeneous case for any set of collision parameters σ and γ except the so-called positive gregarious interaction, i.e. $\sigma = 1$ and $\gamma > 1$. The aim of the present paper is the analysis of simpler (but still rich) equation in this case. More general equation and some details of the present approach will be given in Ref. [5].

2. Mathematical analysis of the space homogeneous case

We focus on the space homogeneous case, i.e. all functions and parameters are assumed to be independent of x . Moreover we assume that σ, γ, β are constants: $\sigma = \beta = 1$, and $\gamma > 1$.

Throughout the paper the L^p -norm in the velocity space \mathbb{V} is denoted by $\|\phi\|_p = \left(\int_{\mathbb{V}} \phi^p dv\right)^{\frac{1}{p}}$. It is easy to see that any solution preserves the nonnegativity of the initial datum and the L^1 -norm of the nonnegative initial datum. Therefore Eq. (1) can be simplified to the following equation

$$\partial_t f = f^\gamma - \|f\|_\gamma^\gamma f, \quad \text{with} \quad f(0, v) = f_0(v), \quad t \geq 0, \quad v \in \mathbb{V}. \tag{3}$$

Let $z(t) = \int_0^t \|f(s, \cdot)\|_\gamma^\gamma ds$. It fulfills

$$d_t z(t) = e^{-\gamma z(t)} \int_{\mathbb{V}} \left(f_0^{1-\gamma} - (\gamma - 1) \int_0^t e^{-(\gamma-1)z(s)} ds \right)^{-\frac{\gamma}{\gamma-1}} dv. \tag{4}$$

Eq. (4) determines global existence or blow-up for Eq. (3). Let $u(t) = \int_0^t e^{-(\gamma-1)z(s)} ds$. The function $u = u(t)$ is increasing and (as we will see) concave. A blow-up occurs for $T > 0$ such that

$$(\gamma - 1) \|f_0\|_\infty^{\gamma-1} u(T) = 1. \tag{5}$$

The ODE for u reads $d_t u(t) = e^{-(\gamma-1)z(t)}$ and we have

$$\begin{aligned} d_t^2 u &= -(\gamma - 1) e^{-(\gamma-1)z} d_t z \\ &= -(\gamma - 1) d_t u (d_t u)^{\frac{\gamma}{\gamma-1}} \int_{\mathbb{V}} \left(f_0^{1-\gamma} - (\gamma - 1) u \right)^{-\frac{\gamma}{\gamma-1}} dv. \end{aligned}$$

By integration we obtain

$$d_t u = \left(\int_{\mathbb{V}} f_0(v) \left(1 - (\gamma - 1) f_0^{\gamma-1} u \right)^{\frac{1}{1-\gamma}} dv \right)^{1-\gamma}. \tag{6}$$

Download English Version:

<https://daneshyari.com/en/article/1707447>

Download Persian Version:

<https://daneshyari.com/article/1707447>

[Daneshyari.com](https://daneshyari.com)