



# Scalar explanation of diphoton excess at LHC

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## Abstract

Inspired by the diphoton signal excess observed in the latest data of 13 TeV LHC, we consider either a 750 GeV real scalar or pseudo-scalar responsible for this anomaly. We propose a concrete vector-like quark model, in which the vector-like fermion pairs directly couple to this scalar via Yukawa interaction. For this setting the scalar is mainly produced via gluon fusion, then decays at the one-loop level to SM diboson channels  $gg, \gamma\gamma, ZZ, WW$ . We show that for the vector-like fermion pairs with exotic electric charges, such model can account for the diphoton excess and is consistent with the data of 8 TeV LHC simultaneously in the context of perturbative analysis.

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## 1. Introduction

The first data at the 13 TeV Large Hadron Collider (LHC) was released on December 15 2015 [1,2]. It shows an excess in diphoton final state at the invariant mass  $M \simeq 750$  GeV, with local significance of order  $3.9\sigma$  and  $2.6\sigma$  for ATLAS and CMS, respectively. In contrast, no excesses in the Standard Model (SM) diboson channels such as  $\gamma\gamma, ZZ, WW, ZW$ , dilepton and dijet were seen in the old data of 8 TeV LHC [3–11].

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Table 1

Matters and their SM quantum numbers in the vector-like quark model. Another fermion doublet  $\tilde{\Psi}$  is added to make sure that the model is free of gauge anomaly.

Matters	$SU(3)_c$	$SU(2)_L$	$U(1)_Y$
$\phi$	<b>1</b>	<b>1</b>	0
$\Psi = (\psi_1, \psi_2)^T$	<b>3</b>	<b>2</b>	$q_\psi$
$\tilde{\Psi} = (\tilde{\psi}_1, \tilde{\psi}_2)^T$	$\bar{\mathbf{3}}$	$\bar{\mathbf{2}}$	$-q_\psi$

If the diphoton excess is indeed a hint of some new physics beyond SM, for an on-shell decay to diphoton it should be due to either spin-0 or spin-2 scalar  $\phi$ . To explain the observed excess, the cross section  $\sigma(pp \rightarrow \phi \rightarrow \gamma\gamma)$  is required to satisfy the signal strength of order,

$$\sigma(pp \rightarrow \phi \rightarrow \gamma\gamma) |_{\sqrt{s}=13 \text{ TeV}} \simeq (8 \pm 3) \text{ fb.} \tag{1.1}$$

Such SM singlet scalar which is responsible for the excess has stimulated extensive interests, see Ref. [12–54].

In this paper, we propose a concrete vector-like quark model, in which the vector-like fermion pairs directly couple to  $\phi$  via tree-level Yukawa interaction. Under our setup,  $\phi$  is mainly produced via gluon fusion, then decays at the one-loop level to SM diboson channels  $gg, \gamma\gamma, ZZ, WW$ , with the colored vector-like fermion pair running in the Feynman loop. For the vector-like fermion pairs with exotic electric charges, such model can account for the diphoton excess, and is consistent with the data of 8 TeV LHC simultaneously in the context of perturbative analysis.

This paper is organized as follows. In Sec. 2 we address the matter content in the vector-like quark model, define the parameter space, and summarize the experimental limits on  $\phi$  and vector-like quark at the 8 TeV LHC. In Sec. 3 we explore the parameter space for  $\phi$  either being a real scalar or pseudo-scalar. Finally, we conclude in Sec. 4.

## 2. The vector-like quark model

### 2.1. The model

In order to reproduce the on-shell decay  $\phi \rightarrow \gamma\gamma$ , which is a loop process for the SM singlet  $\phi$ , we directly couple  $\phi$  to a fermion doublet  $\Psi$ , the latter of which is a subsector of vector-like quark model as defined in Table 1. In this table, another fermion doublet  $\tilde{\Psi}$  is added in order to evade the gauge anomaly problem.

For simplicity, we assume that the mass  $M_{\tilde{\Psi}}$  for  $\tilde{\Psi}$  is obviously larger than the mass  $M_\Psi$  for  $\Psi$ . Below the mass scale  $M_{\tilde{\Psi}}$  the effective Lagrangian in the new physics is described by,<sup>1</sup>

$$\mathcal{L}_{\text{BSM}} = \frac{1}{2} (\partial\phi)^2 - \frac{1}{2} m_\phi^2 \phi^2 + i \bar{\Psi} \gamma^\nu D_\nu \Psi - M_\Psi \bar{\Psi} \Psi + \mathcal{L}_{\text{Yukawa}}, \tag{2.1}$$

where

$$\mathcal{L}_{\text{Yukawa}} = \begin{cases} y\phi \bar{\Psi} \Psi, & (\text{scalar}), \\ iy\phi \bar{\Psi} \gamma_5 \Psi, & (\text{pseudo-scalar}). \end{cases} \tag{2.2}$$

<sup>1</sup> The effective Lagrangian as analyzed in the previous version of this manuscript is a simplification of this concrete one.

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