



Quantum information inequalities via tracial positive linear maps



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ABSTRACT

We present some generalizations of quantum information inequalities involving tracial positive linear maps between C^* -algebras. Among several results, we establish a noncommutative Heisenberg uncertainty relation. More precisely, we show that if $\Phi : \mathcal{A} \rightarrow \mathcal{B}$ is a tracial positive linear map between C^* -algebras, $\rho \in \mathcal{A}$ is a Φ -density element and A, B are self-adjoint operators of \mathcal{A} such that $\text{sp}(-i\rho^{\frac{1}{2}}[A, B]\rho^{\frac{1}{2}}) \subseteq [m, M]$ for some scalars $0 < m < M$, then under some conditions

$$V_{\rho, \Phi}(A) \sharp V_{\rho, \Phi}(B) \geq \frac{1}{2\sqrt{K_{m, M}(\rho[A, B])}} |\Phi(\rho[A, B])|, \quad (0.1)$$

where $K_{m, M}(\rho[A, B])$ is the Kantorovich constant of the operator $-i\rho^{\frac{1}{2}}[A, B]\rho^{\frac{1}{2}}$ and $V_{\rho, \Phi}(X)$ is the generalized variance of X . In addition, we use some arguments differing from the scalar theory to present some inequalities related to the generalized correlation and the generalized Wigner–Yanase–Dyson skew information.

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1. Introduction and preliminaries

In quantum measurement theory, the classical expectation value of an observable (self-adjoint operator) A in a quantum state (density operator) ρ is expressed by $\text{Tr}(\rho A)$. Also, the classical variance for a quantum state ρ and an observable operator A is defined by $V_{\rho}(A) := \text{Tr}(\rho A^2) - (\text{Tr}(\rho A))^2$. The Heisenberg uncertainty relation asserts that

$$V_{\rho}(A)V_{\rho}(B) \geq \frac{1}{4} |\text{Tr}(\rho[A, B])|^2 \quad (1.1)$$

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for a quantum state ρ , where $[A, B] := AB - BA$ is the commutator of two observables A and B ; see [6]. It gives a fundamental limit for the measurements of incompatible observables. A further strong result was given by Schrödinger [15] as

$$V_\rho(A)V_\rho(B) - |\operatorname{Re}(\operatorname{Cov}_\rho(A, B))|^2 \geq \frac{1}{4}|\operatorname{Tr}(\rho[A, B])|^2, \tag{1.2}$$

where the classical covariance is defined by $\operatorname{Cov}_\rho(A) := \operatorname{Tr}(\rho AB) - \operatorname{Tr}(\rho A)\operatorname{Tr}(\rho B)$.

Yanagi et al. [18] defined the one-parameter correlation and the one-parameter Wigner–Yanase skew information (is known as the Wigner–Yanase–Dyson skew information; cf. [10]) for operators A, B , respectively, as follows

$$\operatorname{Corr}_\rho^\alpha(A, B) := \operatorname{Tr}(\rho A^* B) - \operatorname{Tr}(\rho^{1-\alpha} A^* \rho^\alpha B) \quad \text{and} \quad I_\rho^\alpha(A) := \operatorname{Corr}_\rho^\alpha(A, A),$$

where $\alpha \in [0, 1]$. They showed a trace inequality representing the relation between these two quantities as

$$|\operatorname{Re}(\operatorname{Corr}_\rho^\alpha(A, B))|^2 \leq I_\rho^\alpha(A)I_\rho^\alpha(B). \tag{1.3}$$

In the case that $\alpha = \frac{1}{2}$, we get the classical notions of the correlation $\operatorname{Corr}_\rho(A, B)$ and the Wigner–Yanase skew information $I_\rho(A)$. The classical Wigner–Yanase skew information represents a for non-commutativity between a quantum state ρ and an observable A .

Luo [11] introduced the quantity $U_\rho(A)$ as a measure of uncertainty by

$$U_\rho(A) = \sqrt{V_\rho(A)^2 - (V_\rho(A) - I_\rho(A))^2}.$$

He then showed a Heisenberg-type uncertainty relation on $U_\rho(A)$ as

$$U_\rho(A)U_\rho(B) \geq \frac{1}{4}|\operatorname{Tr}(\rho[A, B])|^2. \tag{1.4}$$

These inequalities was studied and extended by a number of mathematicians. For further information we refer interested readers to [3,5,8,17].

Let $\mathbb{B}(\mathcal{H})$ denote the C^* -algebra of all bounded linear operators on a complex Hilbert space $(\mathcal{H}, \langle \cdot, \cdot \rangle)$ with the unit I . If $\mathcal{H} = \mathbb{C}^n$, we identify $\mathbb{B}(\mathbb{C}^n)$ with the matrix algebra of $n \times n$ complex matrices $M_n(\mathbb{C})$. We consider the usual Löwner order \leq on the real space of self-adjoint operators. Throughout the paper, a capital letter means an operator in $\mathbb{B}(\mathcal{H})$. An operator A is said to be strictly positive (denoted by $A > 0$) if it is a positive invertible operator. According to the Gelfand–Naimark–Segal theorem, every C^* -algebra can be regarded as a C^* -subalgebra of $\mathbb{B}(\mathcal{H})$ for some Hilbert space \mathcal{H} . So we may consider elements of \mathcal{A} as Hilbert space operators. We use $\mathcal{A}, \mathcal{B}, \dots$ to denote C^* -algebras. We denote by $\operatorname{Re}(A)$ and $\operatorname{Im}(A)$ the real and imaginary parts of A , respectively. The geometric mean is defined by $A\sharp B = A^{\frac{1}{2}} \left(A^{-\frac{1}{2}} B A^{-\frac{1}{2}} \right)^{\frac{1}{2}} A^{\frac{1}{2}}$ for operators $A > 0$ and $B \geq 0$. A W^* -algebra is a $*$ -algebra of bounded operators on a Hilbert space that is closed in the weak operator topology and contains the identity operator. The C^* -algebra of complex valued continuous functions on the compact Hausdorff space Ω is denoted by $C(\Omega)$.

A linear map $\Phi : \mathcal{A} \rightarrow \mathcal{B}$ between C^* -algebras is said to be $*$ -linear if $\Phi(A^*) = \Phi(A)^*$. It is positive if $\Phi(A) \geq 0$ whenever $A \geq 0$. It is called strictly positive if $A > 0$, then $\Phi(A) > 0$. We say that Φ is unital if \mathcal{A}, \mathcal{B} are unital and Φ preserves the unit. A linear map Φ is called n -positive if the map $\Phi_n : M_n(\mathcal{A}) \rightarrow M_n(\mathcal{B})$ defined by $\Phi_n([a_{ij}]) = [\Phi(a_{ij})]$ is positive, where $M_n(\mathcal{A})$ stands for the C^* -algebra of $n \times n$ matrices with entries in \mathcal{A} . A map Φ is said to be completely positive if it is n -positive for every $n \in \mathbb{N}$. According to [16, Theorem 1.2.4] if the range of the positive linear map Φ is commutative, then Φ is completely positive. It is known (see, e.g., [4]) that if Φ is a unital positive linear map, then

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