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Growth of epitaxial Bi-films on vicinal Si(111)

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ABSTRACT

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1. Introduction

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The semi-metal bismuth (Bi) reveals intriguing physical properties 33 34compared to conventional metals [1-3], which partly even resemble to-35pological insulators [4]. Surface states, e.g., are electrically much more conductive than bulk states, which become readily insulating due to 36 quantum confinement in nano-sized structures. For very thin Bi-films 37 even topologically protected band structures are reported [5]. There-38 fore, details about the surface states in epitaxially grown Bi-films have 39 been studied intensively with transport in the recent past [6–10]. 40

The spin-orbit coupling in Bi is evident from pronounced Rashba-41 splitting of the surface states, which gives rise to spin-polarized trans-4243 port. It limits the allowed channels of electron scattering by defects, so that, e.g., direct back scattering without spin-flip is not possible as 44 long as spin remains a good quantum number. This selection rule can 4546 be lifted by adsorption of magnetic impurities [11–14].

47 Vicinal Bi(111) surfaces reveal further attractive features, as recently shown for Bi(114) crystals. An example is the existence of a one-48 dimensional and spin-split electronic state similarly to edge states in 49 50spin-Hall systems [3,15]. These edge states have been characterized so far only spectroscopically. For probing these states by transport, con-51 trolled epitaxial growth of vicinal Bi-films is mandatory. 52

53Si substrates have been shown to be ideal templates for the growth of well-defined and at the same time electronically decoupled metallic 5455nanostructures. In this respect atomically thin films and wires represent 56the ultimate realization of these concepts [16,17]. For the growth of me-57tallic films, the lattice mismatch as well as surface and interface energies are crucial parameters, which control finally the growth mode and the 58

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film quality. Sticking to the example of growth of Bi on Si(111), the 59 film roughness is larger when grown on a $\sqrt{3} \times \sqrt{3}$ -order wetting 60 layer than for growth on Si(111)-7 \times 7 [19]. In order to explain these 61 findings as well as the film thickness dependence, also allotropy needs 62 to be considered, which is a pronounced effect in semimetals. E.g., the 63 growth of Bi on Si(111) starts in (110) orientation, while (111)-textures 64 appear for larger thickness [18]. Although the preparation of vicinal 65 semi-metallic films is highly desirable in order to study in more detail 66 the properties just mentioned, these considerations show already that 67 the growth of vicinal films is a challenging task. 68

Vicinal semi-metallic Bi-films are expected to reveal topologically protected edge states. In this study the growth 20

of Bi-multilayer structures on Si(557) substrates has been investigated by low energy electron diffraction. There- 21

by, wetting layer structures formed prior to the film deposition on Si(557) surfaces turned out to be crucial for 22

epitaxial growth. Only in the presence of Bi-wetting layers can well-ordered films be grown. In contrast to growth 23

on Si(111), the pseudo-cubic surface of Bi(110) dominates. In addition, Bi(221) surfaces have been obtained only 24

on wetting layers formed by less than a monolayer. The formation of Si(335)-facets during formation of the 25

wetting layers turns out to be essential for the growth of the these structures.

Therefore, as first step towards this field of stepped semi-metallic 69 films, we have systematically varied the structure of the wetting layers 70 formed by various amounts of Bi on Si(557) and studied their influence 71 on further film growth. As it turns out, for Bi coverages ranging from 1/3 72 of a monolayer (ML) to 1 ML not only the local density of Bi atoms of the 73 wetting layer changes, but also the step morphology of the Si(557) sur-74 face is modified. For wetting layers of less than a monolayer of Bi, (335)-75 oriented facets are formed. Stepped Bi(110) films can be grown only on 76 these facet structures. To the best of our knowledge this is the first study 77 which deals with the growth of stepped Bi-films on vicinal Si substrates. 78

2. Experimental setup

The experiments were performed in an ultra-high vacuum system 80 operating below a base pressure of 1×10^{-10} mbar. The system hosts 81 a low energy electron diffraction system with the capability for spot 82 profile analysis (SPA-LEED) in order to analyze the quality of the sub- 83 strate and the films with high resolution. Clean Si(557) substrates 84 have been obtained by several heating cycles to 1370 K for a few sec- 85 onds after the sample had been degassed for hours at 870 K. Bi was 86 evaporated from a Knudsen cell. The amount of Bi was controlled by a 87

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guartz microbalance and the coverage has been calibrated by character-88 89 istic reconstructions on nominally flat Si(111) surfaces, and by oscillations in conductivity due to layer-by-layer growth of Bi on pre-90 91annealed Bi-films. Wetting layers with different concentrations of Bi have been realized by adsorption of 5ML at 300 K followed by anneal-92ing to various temperatures. Annealing for 1 min at 820 K and 620 K re-93 sults in a remaining 1/3 ML (so-called α -phase) and 1 ML (β -phase) 9495coverages, respectively. For intermediate annealing temperatures, 96 e.g. 720 K, a nominal coverage of 2/3 ML remains, which we refer to 97 as γ -phase in this work.

Bi films have been grown by deposition of Bi at 200 K followed by 98 annealing to 350 K. The wetting layer coverage is given in monolayers 99 with respect to Si(111) (1 ML = 7.84×10^{14} atoms/cm²). For Bi-films 100 we use two different calibrations: for the (111)-textured films the 101 coverage is given in bilayers (1 BL = 1.14×10^{15} atoms/cm²) while 102 we refer to monolayers in the case of Bi(110) (1 ML = 9.27×10^{14} 103 $atoms/cm^{2}$). The indexing of the Bi-films is given with respect to the 104 rhombohedral notation. For a definition of the basic vectors $\overrightarrow{a_1}$, $\overrightarrow{a_2}$ 105and $\overrightarrow{a_3}$, see Fig. 5c). 106

107 3. Results and discussion

The growth of semi-metallic films depends crucially on details of the 108 wetting layers. Therefore, the different Bi-induced facet structures de-109110 pending on Bi coverage will be presented first. The growth of Bi-films 111 of about 10 nm in thickness is studied in the second part of this paper. We want to emphasize that we restrict ourselves in this study to film 112 thicknesses above 10 nm in order to avoid the influence of allotropic 113 phase transitions, which have been shown to occur for thinner films 114 [18,19]. 115

116 3.1. Wetting layers of Bi on Si(557)

The Si(557) surface is a non-homogeneously stepped surface along the $\left[\overline{1}\ \overline{12}\right]$ -direction. The (7 × 7) reconstruction on small (111) terraces as well as dimerization at the step edges of (112)-oriented facets favor the alternating arrangement of these two facets 9 and 8 atomic units wide, respectively. They are separated by three bilayer steps so that the unit cell is (17 × 3.32Å)² + (3 × 3.13Å)²]^{1/2} = 57Å. More details are described in refs. [21–25].

Upon adsorption of Bi and annealing, the reconstructions of the Si(557) surface changes, as illustrated by the LEED patterns after growth of Bi wetting layers with nominally 1/3, 2/3, and 1 ML shown in Fig. 1. Common to all phases is the appearance of $\sqrt{3} \times \sqrt{3}$ diffraction spots (blue dashed lines) compatible with the trimer-model for Bi adsorption on Si(111) [26]. Compared to the adsorption on Si(111), the LEED patterns shown in 130 Fig. 1 reveal further substructures, which depend on Bi concentration, 131 e.g. splitting of the integer spots ((1×1) unit cell marked by red dotted 132 lines) along the $\left[\overline{1}\,\overline{1}2\right]$ direction. We will start the discussion with the situation after adsorption of 1 ML Bi, i.e. appearance of the Si(557)– 134 γ Bi-phase. The line scan through the (00)-spot shown as inset in 135 Fig. 1c) reveals a spot splitting of 5.8% SBZ (SBZ = surface Brillouin 136 zone), equivalent to a periodicity with 17 atomic units. In other 137 words, the initial step structure remains unchanged. Compared to dif-138 fraction patterns obtained from clean Si(557), only few spots of the 139 step train are visible, similar to recent results obtained by adsorption 140 of 1 ML Ag on Si(557) [27], which is indicative of modified form factors 141 and/or reduced long-range ordering.

When wetting layers on Si(557) are formed with submonolayer covrages of Bi, the spot splitting is fundamentally different from the clean Si(557) surface. In order to quantify the strong refacetting process upon formation of these wetting layers, line scans have been taken along the $\overline{1112}$ direction through the (00)-beam at various scattering conditions 147

S. These were plotted as a gray-scale coded picture (Fig. 2a-c). From 148 such plots (called $(k_{\parallel},k_{\perp})$ -plots in the following) the formation and in- 149 clination of faceted areas can be directly read off. The scattering phase 150 S (S = $\frac{2\pi}{k}$) plotted in Fig. 2 is given with respect to the step height of 151 Si(111) of d = 3.13 Å. k_{\perp} denotes the vertical component of the scatter-152 ing vector. As deduced from the intersections of the rods with the Bragg 153 points in Fig. 2a), (113)- and (335)-facets appear after adsorption of 154 1/3 ML of Bi and corroborate the Bi-induced refacetting of the Si(557) 155 substrate. The (335)-facet has an inclination of 14.4° with respect to 156 the (111)-planes, which is significantly larger than the average tilt 157 angle of the (557)-surface (9.45°) . A sketch of this facet orientation in 158 real space is shown in Fig. 2d). This demands other facet orientations 159 in order to maintain the average inclination, or, alternatively, a widen- 160 ing of the (111)-oriented terraces already present on the (557) surface. 161 The high intensity along the (111)-rods of Fig. 2a indicates that this may 162 have happened. All other orientations, are obviously not periodically ar- 163 ranged and are visible in LEED only as background, as already found ear- 164 lier for Pb/Si(557) [28]. The (113)-facets are a minority species, as 165 evident from the faint intensities. 166

Although the (335)-facets have a width of only 12.1 Å^o $\left(3\frac{2}{3} \times 3.32\text{\AA}\right)$ 167 the facets are covered with Bi. This becomes obvious when analyzing 168 facet rods of the $\sqrt{3}$ -reconstruction spots shown in Fig. 2b). The inclination of the rods is identical to that before for the integer-spots indicating 170 also a strong correlation between reconstructions on adjacent terraces. 171

Long range-ordered arrays with (335)-facets have been found only 172 on the Si(557)– α Bi-phase, i.e. at 1/3 ML Bi coverage (Fig. 2d). This 173 means that there is a Bi-induced minimum of (free) energy for this 174



Fig. 1. Bi-wetting layers on Si(557) with coverage of (a) 1/3 ML (α -phase), (b) 2/3 ML (γ -phase) and (c) 1 ML (β -phase). The electron energies were 140 eV (a, b) and 80 eV (c). The dotted (red) and dashed (blue) rhombi denote the unit cells of the 1 × 1 and $\sqrt{3} \times \sqrt{3}$ structures, respectively. Images were taken at 80 K. The inset in c) shows a line scan taken along the $\begin{bmatrix} \overline{1} & \overline{1} \\ 1 \end{bmatrix}$ direction through the central spot (log-intensity vs. SBZ, $\Delta K_{\parallel} = 5.8$ % SBZ). (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

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