

# On the entropy minimization problem in Statistical Mechanics 

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## A R T I C L E I N F O

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#### Abstract

In many works on Statistical Mechanics and Statistical Physics, when deriving the distribution of particles of ideal gases, one uses the method of Lagrange multipliers in a formal way. In this paper we treat rigorously this problem for Bose-Einstein, Fermi-Dirac and Maxwell-Boltzmann entropies and present a complete study in the case of the Maxwell-Boltzmann entropy. Our approach is based on recent results on series of convex functions.


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## 1. Introduction

In Statistical Mechanics and Statistical Physics, when studying the distribution of the particles of an ideal gas, one considers the problem of maximizing

$$
\begin{equation*}
\sum_{i}\left[n_{i} \ln \left(\frac{g_{i}}{n_{i}}-a\right)-\frac{g_{i}}{a} \ln \left(1-a \frac{n_{i}}{g_{i}}\right)\right] \tag{1.1}
\end{equation*}
$$

with the constraints $\sum_{i} n_{i}=N$ and $\sum_{i} n_{i} \varepsilon_{i}=E$, where, as mentioned in [5, pp. 141-144], $\varepsilon_{i}$ denotes the average energy of a level, $g_{i}$ the (arbitrary) number of levels in the $i$ th cell, and, in a particular situation, $n_{i}$ is the number of particles in the $i$ th cell. Moreover, $a=-1$ for the Bose-Einstein case, +1 for the Fermi-Dirac case, and 0 for the (classical) Maxwell-Boltzmann case. Even if nothing is said explicitly about the set $I$ of the indices $i$, from several examples in the literature, $I$ is (or may be) an infinite countable set; the examples

$$
\begin{align*}
\varepsilon_{l} & =l(l+1) h^{2} / 2 I, \quad g_{l}=(2 l+1) ; \quad l=0,1,2, \ldots  \tag{1.2}\\
\varepsilon_{v K} & =\varepsilon_{0}+h \omega\left(v+\frac{1}{3}\right)+h^{2} K(K+1) / 2 I ; \quad v, K=0,1,2, \ldots \tag{1.3}
\end{align*}
$$

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$$
\begin{equation*}
\varepsilon\left(n_{x}, n_{y}, n_{z}\right)=\frac{h^{2}}{8 m L^{2}}\left(n_{x}^{2}+n_{y}^{2}+n_{z}^{2}\right) ; \quad n_{x}, n_{y}, n_{z}=1,2,3, \ldots \tag{1.4}
\end{equation*}
$$

\]

are considered in [3, p. 76], [4, p. 138] and [5, p. 10], respectively.
Relation (1.1) suggests the consideration of the following functions defined on $\mathbb{R}$ with values in $\overline{\mathbb{R}}$, called, respectively, Bose-Einstein, Fermi-Dirac and Maxwell-Boltzmann entropies:

$$
\begin{gather*}
E_{B E}(u):= \begin{cases}u \ln u-(1+u) \ln (1+u) & \text { if } u \in \mathbb{R}_{+}, \\
\infty & \text { if } u \in \mathbb{R}_{-}^{*},\end{cases}  \tag{1.5}\\
E_{F D}(u):= \begin{cases}u \ln u+(1-u) \ln (1-u) & \text { if } u \in[0,1], \\
\infty & \text { if } u \in \mathbb{R} \backslash[0,1],\end{cases}  \tag{1.6}\\
E_{M B}(u):= \begin{cases}u(\ln u-1) & \text { if } u \in \mathbb{R}_{+}, \\
\infty & \text { if } u \in \mathbb{R}_{-}^{*},\end{cases} \tag{1.7}
\end{gather*}
$$

where $0 \ln 0:=0$ and $\mathbb{R}_{+}:=\left[0, \infty\left[, \mathbb{R}_{+}^{*}:=\right] 0, \infty\left[, \mathbb{R}_{-}:=-\mathbb{R}_{+}, \mathbb{R}_{-}^{*}:=-\mathbb{R}_{+}^{*}\right.\right.$. We have that

$$
\left.E_{B E}^{\prime}(u)=\ln \frac{u}{1+u} \forall u \in \mathbb{R}_{+}^{*}, \quad E_{F D}^{\prime}(u)=\ln \frac{u}{1-u} \forall u \in\right] 0,1\left[, \quad E_{M B}^{\prime}(u)=\ln u \forall u \in \mathbb{R}_{+}^{*} .\right.
$$

Observe that $E_{B E}, E_{M B}, E_{F D}$ are convex (even strictly convex on their domains), derivable on the interiors of their domains with increasing derivatives, and $E_{B E} \leq E_{M B} \leq E_{F D}$ on $\mathbb{R}$. The (convex) conjugates of these functions are

$$
E_{M B}^{*}(t)=e^{t} \forall t \in \mathbb{R}, \quad E_{F D}^{*}(t)=\ln \left(1+e^{t}\right) \forall t \in \mathbb{R}, \quad E_{B E}^{*}(t)= \begin{cases}-\ln \left(1-e^{t}\right) & \text { if } t \in \mathbb{R}_{-}^{*}, \\ \infty & \text { if } t \in \mathbb{R}_{+}\end{cases}
$$

Moreover, for $W \in\left\{E_{B E}, E_{M B}, E_{F D}\right\}$ we have that $\partial W(u)=\left\{W^{\prime}(u)\right\}$ for $u \in \operatorname{int}(\operatorname{dom} W)$ and $\partial W(u)=\emptyset$ elsewhere; furthermore,

$$
\begin{equation*}
\left(W^{*}\right)^{\prime}(t)=\frac{e^{t}}{1+a_{W} e^{t}} \quad \forall t \in \operatorname{dom} W^{*}, \tag{1.8}
\end{equation*}
$$

where (as above)

$$
a_{W}:= \begin{cases}-1 & \text { if } W=E_{B E}  \tag{1.9}\\ 0 & \text { if } W=E_{M B} \\ 1 & \text { if } W=E_{F D}\end{cases}
$$

The maximization of (1.1) subject to the constraints $\sum_{i} n_{i}=N$ and $\sum_{i} n_{i} \varepsilon_{i}=E$ is equivalent to the minimization problem

$$
\operatorname{minimize} \sum_{i} g_{i} W\left(\frac{n_{i}}{g_{i}}\right) \quad \text { s.t. } \sum_{i} n_{i}=N, \quad \sum_{i} n_{i} \varepsilon_{i}=E
$$

where $W$ is one of the functions $E_{B E}, E_{F D}, E_{M B}$ defined in (1.5), (1.6), (1.7), and $g_{i} \geq 1$.
In many books treating this subject (see [4, pp. 119, 120], [3, pp. 15, 16], [5, p. 144], [1, p. 39]) the above problem is solved using the Lagrange multipliers method in a formal way.

Our aim is to treat rigorously the minimization of Maxwell-Boltzmann, Bose-Einstein and Fermi-Dirac entropies with the constraints $\sum_{i \in I} u_{i}=u, \sum_{i \in I} \sigma_{i} u_{i}=v$ in the case in which $I$ is an infinite countable set. Unfortunately, we succeed to do a complete study only for the Maxwell-Boltzmann entropy. For a short description of the results see Conclusions.

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