ARTICLE IN PRESS

Organic Electronics xxx (2014) xxx-xxx

Contents lists available at ScienceDirect

Organic Electronics



5 6

 $\frac{14}{15}$

journal homepage: www.elsevier.com/locate/orgel

Band bending at the P3HT/ITO interface studied 3

by photoelectron spectroscopy

7 Q1 Martin Schneider^a, Alexander Wagenpfahl^b, Carsten Deibel^{b,1}, Vladimir Dyakonov^{b,c}, Achim Schöll^{a,d,*}. Friedrich Reinert^{a,d} 8

9 ^a Universität Würzburg, Experimentelle Physik VII & Röntgen Research Center for Complex Material Systems (RCCM), 97074 Würzburg, Germany

10 ^b Experimental Physics VI, Julius-Maximilians-University of Würzburg, 97074 Würzburg, Germany

11 ^c Bavarian Center for Applied Energy Research, e.V. (ZAE Bayern), D-97074 Würzburg, Germany

12 ^d Karlsruher Institut für Technologie (KIT), Gemeinschaftslabor für Nanoanalytik, 76021 Karlsruhe, Germany

ARTICLE INFO

37 18 Article history:

- 19 Received 30 August 2013
- 20 Received in revised form 10 March 2014
- 21 Accepted 17 March 2014
- 22 Available online xxxx
- 23 Q3 Keywords:
- 24 РЗНТ
- 25 ITO
- 26 Band bending
- 27 Photoelectron spectroscopy
- 28 Defect density 29
- Interface dipole 30
- 43

1. Introduction 44

Poly (3-hexylthiophene) (P3HT) is one of the most 45 important polymer materials used as organic semiconduc-46 47 tor in electronic devices [1,2]. One of the key issues for the 48 device function and performance is the charge transfer across the interface between the organic semiconductor 49 and the electrodes [3,4]. This charge injection or extraction, 50 51 respectively, is determined by the electronic structure in 52 the interfacial regime [5] which thus demands particular 53 attention.

E-mail address: achim.schoell@physik.uni-wuerzburg.de (A. Schöll).

¹ Present address: Institute of Physics, Chemnitz University of Technology, 09107 Chemnitz, Germany.

http://dx.doi.org/10.1016/j.orgel.2014.03.012 1566-1199/© 2014 Elsevier B.V. All rights reserved.

ABSTRACT

The electronic structure of poly (3-hexylthiophene) (P3HT)/indium tin oxide interfaces was 32 studied by photoelectron spectroscopy (PES). In the thickness dependent core and valence 33 PES data we observe a significant band bending of about (0.8 ± 0.1) eV within 40 nm 34 distance of the interface. The thickness dependence of the polaronic state follows an 35 exponential law, thus indicating a Schottky-Mott-like behaviour due to trap states. From 36 37 a drift-diffusion simulation we can determine a p-type doping concentration of $C = (5^{+5}_{-4}) \cdot 10^{15}$ cm⁻³ and a width of the Gaussian density of states of $\sigma = 155$ meV for a 38 vanishing injection barrier. 39

© 2014 Elsevier B.V. All rights reserved.

42

54

55

56

57

58

60

61

62

63

64

65

66

67

68

69

70

40

41

Charge extraction barriers can lead to s-shaped current-voltage characteristics [6] and thus reduce the fill factor and consequently the power conversion efficiency of SCs. Moreover, injection barriers can substantially reduce the open circuit voltage in case of solar cells (SCs) [7–9]. Common electrode materials are metals (such as silver, 59 gold, and aluminum) and in case of optolectronic devices transparent conductive oxides such as indium tin oxide (ITO) are used to establish an optically transparent contact. P3HT/ITO contacts are hence relevant model systems to study the particularities of the energy level alignment between the organic semiconductor and the contact in the interfacial region.

In this paper, we performed PES on P3HT/ITO contacts, varying the polymer thickness to investigate the impact of the electrode on the organic semiconductor energy levels.

1

Q2 * Corresponding author at: Universität Würzburg, Experimentelle Physik VII & Röntgen Research Center for Complex Material Systems (RCCM), 97074 Würzburg, Germany. Tel.: +49 9313185127.

2 May 2014

Q1 2

71 2. Experimental

All X-ray and UV photoelectron spectroscopy measure-72 73 ments (XPS and UPS, respectively) were performed in a 74 UHV system, described in detail elsewhere [10]. The base pressure was below $2 \cdot 10^{-10}$ mbar and the samples were 75 at room temperature (T = 300 K). For XPS the Mg K₂ anode 76 77 of a twin-anode source was used, while the He-I excitation line of a gas discharge lamp was applied for UPS. For work 78 79 function measurements the sample was biased by -9 V.

80 After rinsing with acetone, the ITO coated glass sub-81 strates were oxygen plasma-treated resulting in surfaces with low contamination as verified by XPS. The surface 82 roughness was determined to 2.8-4.7 nm from X-ray 83 diffraction scans and surface profiler measurements (Veeco 84 85 Dektak 150). The plasma-treatment resulted in a very reproducible sample work function of $\phi = (4.78 \pm 0.05) \text{ eV}$ in 86 accordance with literature [11,12]. A representative work-87 88 function measurement on a plasma-treated ITO sample is 89 displayed in Fig. 1. Note that the work function after plasma 90 treatment is significantly larger than reported for nonplasma-treated samples [13,12,14,15] and is furthermore 91 92 stable under X-ray and UV-illumination on the experimen-93 tal time scale of several minutes.

In order to access the properties of the P3HT/ITO inter-94 95 facial regime with surface sensitive techniques such as PES the preparation of ultra-thin P3HT films is required. The 96 97 generally applied spin-coating, however, implies a lower 98 limit of the film thickness of a few nm. For lower coverages 99 films become inhomogenous. As a consequence other preparation routes have been applied, such as thermal 100 101 evaporation [16] or electrospray deposition in UHV [17], or post-preparation thinning out by solvents [18]. 102 103 Although these approaches allow performing experiments 104 with P3HT samples down to sub-monolayer coverages, we utilized spin-coating in this work in order to grant film 105 106 structures and morphologies comparable to those used in various other experiments and in applications. 107

The spin-coating was performed at 1500 rpm from solutions between 20 mg/ml and 25 mg/ml of regioregular P3HT (purchased from Rieke Metals) in chlorobenzene onto the ITO substrates. The resulting layer thickness was



Fig. 1. UPS scan of an oxygen plasma-treated ITO sample. The determination of the work function from the secondary electron cut off is illustrated. The sample was biased by -9 V and the energy scale is referenced to the Fermi energy E_F .

118

119

120

121

122

123

124

125

126

127

128

129

130

149

between 5 nm and 152 nm, as determined by a surface profiler for samples thicker than 25 nm and by optical absorption for the ultra-thin films. The P3HT film preparation and sample mounting was performed in a glove box. The transfer time into the UHV system was kept below 116 10 s in order to minimize contamination. 117

3. Experimental results

ARTICLE IN PRESS

M. Schneider et al. / Organic Electronics xxx (2014) xxx-xxx

Fig. 2 displays the result of thickness dependent PES measurements on P3HT thin films prepared on ITO/glass substrates. The film thickness variation allows us to access the evolution of the electronic structure in the interface regime. The C 1s (Fig. 2a) and S 2p (Fig. 2b) data consistently show a shift of the core level positions towards lower binding energy in the vicinity of the P3HT/ITO interface. The same shift is observed for the P3HT HOMO band in Fig. 2c. Here, the delocalized π -band results in the broad feature which spreads over 1.5 eV, followed by a sharper peak at around 3.5 eV, which can be attributed to a localized P3HT π -band [19].

Fig. 3 summarizes the thickness dependent energy shift 131 of the C 1s (Fig. 3a), S 2p_{3/2} (Fig. 3b), and HOMO level 132 (Fig. 3c) derived from Fig. 2. While the energetic positions 133 were straightforwardly derived from the maxima of the C 134 1s peaks in Fig. 2a, the HOMO onset was used as a refer-135 ence in the valence spectra (Fig. 2c). The position of the S 136 $2p_{3/2}$ level was determined from a fit to the experimental 137 data by three Gaussian functions. One Gaussian each was 138 used for the S $2p_{3/2}$ and S $2p_{1/2}$ signals utilizing constraints 139 for the area $(A_{3/2} = 2 \cdot A_{1/2})$, width $(w_{3/2} = w_{1/2})$ and spin-140 orbit splitting ($\Delta E = 1.1 \text{ eV}$). An additional Gaussian was 141 allowed to model the high energy background from satel-142 lites. The fit results are plotted in Fig. 2b as solid lines. In 143 addition, the results of work function measurements on 144 the same samples are plotted in Fig. 3d. The comparison 145 shows that the energetic position of the P3HT core and 146 valence levels in the interface regime immediately follow 147 the sample work function. 148

4. Discussion

The energy level alignment at the P3HT/ITO interface 150 clearly shows a behaviour reminiscent of band bend-151 ing.[20,4] For contacts between inorganic semiconductors 152 and metals this is a well-known effect and is generally 153 described by the Schottky-Mott model. According to this 154 model, a band bending is observed due to the smaller work 155 function of the semiconductor, which leads to a transfer of 156 electrons to the metal and thus to a positive space charge 157 in the semiconductor interface region. The resulting elec-158 trostatic potential "bends" the electronic levels of the 159 semiconductor upwards in the vicinity of the metal con-160 tact. The magnitude of this built-in potential can be quan-161 tified from the difference of the energy position of the 162 valence levels at and far away from the interface. In our 163 case it can be estimated to more than 0.7 eV from the bind-164 ing energy of the P3HT HOMO in thick films and an extrap-165 olation of the HOMO position for very thin samples (see 166 Fig. 2c). 167

Q1 Please cite this article in press as: M. Schneider et al., Band bending at the P3HT/ITO interface studied by photoelectron spectroscopy, Org. Electron. (2014), http://dx.doi.org/10.1016/j.orgel.2014.03.012 Download English Version:

https://daneshyari.com/en/article/10566291

Download Persian Version:

https://daneshyari.com/article/10566291

Daneshyari.com