

Two-dimensional phase plane structure and the stability of the orbital motion for space debris in the geosynchronous ring

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Abstract

Based on the orbital resonance model, we study the two-dimensional phase plane structure of the motion of space debris orbiting the geosynchronous ring under the combined effects of the tesseral harmonics J_{22} , J_{31} and J_{33} of the Earth's gravitational field. We present the main characteristic parameters of the two-dimensional phase plane structure. We also analyze the stability of the two-dimensional phase plane structure with numerical method. Our main findings indicate that the combined effects of the tesseral harmonics J_{22} , J_{31} and J_{33} fully determine the two-dimensional phase plane structure of the space debris, and it remains robust under the effect of the Earth's actual gravitational field, the luni-solar perturbations and the solar radiation pressure with the normal area-to-mass ratios.

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1. Introduction

A distinctive dynamical phenomenon of the space debris orbiting the geosynchronous ring is orbital resonance with the Earth rotation. Its principal resonant term arises from the tesseral harmonic J_{22} associated with the Earth's equatorial ellipticity causing libration of the space debris orbiting the geosynchronous ring around the nearest stable position either about 75°E or about 105°W (Blitzer et al., 1962; Schildknecht, 2007). The Ideal Resonance Problem of the uncontrolled satellites (or space debris) orbiting the

geosynchronous ring was first developed by Garfinkel (1966) with the principal zonal harmonic J_2 and tesseral harmonic J_{22} of the Earth's gravitational field only. Many scientists studied this orbital resonance problem arising from tesseral harmonic J_{22} (Jupp, 1969; Dallas and Diehl, 1977; Nacozy and Diehl, 1982; Liu et al., 1991; Zhao and Liu, 1991; Valk et al., 2009). Liu et al. discussed the two-dimensional (2D) phase plane structure and its main characteristic parameters of the 1:1 resonance of the satellites orbiting the geosynchronous ring based on the ideal resonance model in detail, and presented the analytical expressions of the libration period and libration width of the 1:1 resonance of the satellites (Liu et al., 1991; Zhao and Liu, 1991).

However, because of asymmetry of the Earth's gravitational field due to the higher degree and order tesseral harmonics, especially the third degree terms, the space debris (or uncontrolled satellites) orbiting the geosynchronous ring can also move in the way of complex ('long') libration around both stable positions (we can call it double libration) besides the way of simple libration around the nearest

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stable position (either about 75°E or about 105°W) (Floh-
rer et al., 2011). Vashkov'yak and Lidov first found the
effect of splitting separatrix of the geostationary satellite
under the influence of the higher degree and order tesseral
harmonics of the Earth's gravitational field, which leads to
the appearance of the libration zone including both stable
positions; and obtained the estimates of the libration peri-
ods in the regions surrounding stable positions: 740 days in
the neighborhood of the position with a longitude of 75°E,
930 days in the neighborhood of the position 105°W and
the minimal libration period with respect to both stable
positions 2900 days (Vashkov'yak and Lidov, 1973; Kuz-
netsov and Kaiser, 2007). The Ideal Resonance Problem
of the objects orbiting the geosynchronous ring cannot
explain this double libration phenomenon. Giacaglia pre-
sented an Extended Resonance Problem exhibiting the case
of double libration (Giacaglia, 1970). Garfinkel defined
another form of an Extended Resonance Problem with
the phase-plane trajectories exhibiting a double libration,
constructed the first order solution for the case of double
libration by a generalization of the procedure used in solv-
ing the Ideal Resonance Problem, and presented a General-
ized Ideal Resonance Problem with the double libration
illustrated by the horseshoe-shaped orbits enclosing two
libration centers (Garfinkel, 1975, 1976). These Extended
Resonance Problems or Generalized Ideal Resonance
Problem stated in Giacaglia (1970); Garfinkel (1975,
1976) are generalized theoretical analyses and were not
applied to the actual motion of the objects orbiting the geo-
synchronous ring. Lara and Elife generated families of pla-
nar periodic orbits emanating from the geostationary
points (both stable and unstable) and showed that even
for the unstable points, it is possible to have stable periodic
orbits, which also reflects the existence of the double libra-
tion (Lara and Elife, 2002).

Exactly because the space debris orbiting the geosyn-
chronous ring has the double libration regime besides sim-
ple libration (Floh-
rer et al., 2011), while its conventional
ideal resonance model with only a resonant term arising
from tesseral harmonic J_{22} cannot explain this regime, it
is necessary to establish a new ideal resonance model to
better characterize its 2D phase plane patterns of the
motion. Inspired by Garfinkel's idea: if the sine of an
odd multiple of resonant variable x is added into the con-
ventional ideal resonance model, the topological picture
exhibits the novel feature of double libration enclosing
two libration centers (Garfinkel, 1975), an extended orbital
resonance model of the motion of space debris orbiting the
geosynchronous ring is established in this paper by consid-
ering the combined effects of the zonal harmonic J_2 and
tesseral harmonics J_{22}, J_{31} and J_{33} of the Earth's gravita-
tional field. Based on this extended orbital resonance
model, the 2D phase plane structure (in which the double
libration regime is included) and its main characteristic
parameters are investigated. Under the effect of the Earth's
actual gravitational field, the luni-solar perturbations and
the solar radiation pressure with the normal area-to-mass

ratios (A/M) less than $0.02m^2/kg$, the stability of the 2D
phase plane structure is analyzed with numerical method.

2. An extended orbital resonance model with J_{22}, J_{31} and J_{33}

The Earth's gravitational potential V may be expressed
in terms of geocentric distance r , latitude φ and geographic
longitude λ , as

$$\begin{aligned} V &= -\frac{\mu}{r} \left[1 + \sum_{n=2}^{\infty} \sum_{m=0}^n \left(\frac{a_e}{r} \right)^n P_{nm}(\sin \varphi) (C_{nm} \cos m\lambda + S_{nm} \sin m\lambda) \right] \\ &= -\frac{\mu}{r} \left[1 - \sum_{n=2}^{\infty} J_n \left(\frac{a_e}{r} \right)^n P_n(\sin \varphi) \right. \\ &\quad \left. + \sum_{n=2}^{\infty} \sum_{m=1}^n J_{nm} \left(\frac{a_e}{r} \right)^n P_{nm}(\sin \varphi) \cos m(\lambda - \lambda_{nm}) \right], \quad (1) \end{aligned}$$

where $\mu = G_N M_e$ is the product of the Newtonian constant of
gravitation by the mass of the Earth, a_e is the Earth mean
equatorial radius, C_{nm} and S_{nm} are the spherical harmonic coef-
ficients of degree n and order m of the Earth's gravitational
field, P_n, P_{nm} are the Legendre polynomial and associated
Legendre function respectively, $J_{nm} = \sqrt{C_{nm}^2 + S_{nm}^2}$, $J_n =$
 $-C_{n0}$, and $m\lambda_{nm} = \arctan(S_{nm}/C_{nm})$.

Introduce the following canonical variables

$$\begin{cases} L = \sqrt{\mu a} & l = M + \omega + \Omega - \theta_G, \\ G = \sqrt{\mu a(1 - e^2)} - L & g = \omega, \\ H = \sqrt{\mu a(1 - e^2)} \cos i - L & h = \Omega, \end{cases} \quad (2)$$

which are linear combinations of the Delaunay canonical
variables, where a, e, i, ω, M and Ω are the classical orbital
elements (Ω is measured from the equinox), θ_G is the
Greenwich hour angle satisfying $\theta_G = n_e$, and n_e is the rota-
tional angular velocity of the Earth.

We first focus on the orbital resonance problem due to
geopotential perturbation here, and leave the luni-solar
perturbation effects for the space debris orbiting the geo-
synchronous ring aside. If only the secular and resonant
terms of the dominant zonal harmonic J_2 and main tesseral
harmonics J_{22}, J_{31} and J_{33} are considered, by averaging
short periodic terms and neglecting the second order effects
due to J_2 , the Hamiltonian function of the 1:1 orbital re-
sonance problem of the space debris orbiting the near circu-
lar geosynchronous ring can be expressed in the canonical
variables Eq. (2) as

$$\begin{aligned} F &= \frac{\mu^2}{2L^2} + n_e L + \frac{\mu^4 a_e^2}{L^6} J_2 \left(\frac{1}{2} - \frac{3}{4} \sin^2 i \right) \\ &\quad + \frac{\mu^4 a_e^2}{L^6} J_{22} f_{220}(i) \cos 2(l - \lambda_{22}) + \frac{\mu^5 a_e^3}{L^8} J_{31} f_{311}(i) \\ &\quad \times \cos(l - \lambda_{31}) + \frac{\mu^5 a_e^3}{L^8} J_{33} f_{330}(i) \cos 3(l - \lambda_{33}), \quad (3) \end{aligned}$$

where $f_{220}(i), f_{311}(i)$ and $f_{330}(i)$ are inclination functions
and defined as (Kaula, 1966)

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