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Search for large extra dimensions in dimuon and dielectron events in pp collisions at $\sqrt{s} = 7$ TeV ☆

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ABSTRACT

Results are presented from a search for large, extra spatial dimensions in events with either two isolated muons or two isolated electrons. The data are from proton–proton interactions at $\sqrt{s} = 7$ TeV collected with the CMS detector at the LHC. The size of the data sample corresponds to an integrated luminosity of approximately 2 fb^{-1}. The observed dimuon and dielectron mass spectra are found to be consistent with standard-model expectations. Depending on the number of extra dimensions, the 95% confidence level limits from the combined $\mu\mu$ and ee channels range from $M_s > 2.4$ TeV to $M_s > 3.8$ TeV, where M_s characterizes the scale for the onset of quantum gravity.

Models that extend the structure of space–time predict new phenomena beyond the standard model (SM) of particle physics. Additional spatial dimensions, essential for formulating quantum gravity in the context of string theory, have been proposed as a solution to the SM hierarchy problem [1–3]. In this Letter, we present a search for events at large dimuon or dielectron invariant mass due to contributions from virtual-graviton processes in the Arkani-Hamed–Dimopoulos–Dvali (ADD) model [1,2].

The ADD model postulates the existence of compactified extra dimensions. Gravity is assumed to propagate in the entire higher-dimensional space, while particles of the SM are confined to a 3-dimensional slice of the multidimensional space. The resulting fundamental Planck energy scale M_D in the ADD model can be reduced to significantly lower values than suggested by the apparent Planck mass $M_{Pl} \approx 1.2 \cdot 10^{19}$ GeV deduced for 3 spatial dimensions. M_D must be of the order of the scale of electroweak symmetry breaking to provide an explanation of the hierarchy problem. This scenario predicts phenomenological effects that might be observed in proton–proton collisions at the LHC. In this Letter, we adopt the assumption [4,5] that all extra dimensions are compactified on a torus of size r. In this case, M_D is related to M_{Pl} through $M_D^{n+2} = M_{Pl}^2/(8\pi r^n)$, where n is the number of extra dimensions.

The graviton in this $(3+n)$-dimensional formulation can be equivalently expressed as a set of 3-dimensional Kaluza–Klein (KK) modes [6] with different graviton masses. The coupling of the KK modes to the SM energy–momentum tensor leads to an effective theory with virtual-graviton exchange at leading order (LO) in perturbation theory. An ultraviolet (UV) cutoff Λ must be introduced to avoid divergences in the summed contributions from all modes. A phenomenological consequence of the small mass separation between adjacent KK modes is an enhancement in the expected rate of dilepton events at large invariant masses that appears to be non-resonant. Depending on the details of the model, virtual-graviton effects can provide the dominant experimental ADD signature at high-energy colliders [4,5].

Several ways of parameterizing the LO differential cross sections are provided in the literature, including the Han–Lykken–Zhang (HLZ) [4] and the Giudice–Rattazzi–Wells (GRW) [5] conventions. In the GRW convention, the leading-order phenomenology for partonic center-of-mass energies $\sqrt{\hat{s}} \ll \Lambda$ is described by a single parameter Λ_T, which does not depend on $\sqrt{\hat{s}}$ for $n \geqslant 3$. The HLZ convention describes the phenomenology in terms of n and a mass scale M_s, where M_s is related to the selected UV cutoff and reflects the scale for the onset of quantum gravity. Typically, M_s is expected to be of order M_D. The parameter Λ_T can be related to the

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parameters in the HLZ convention [7]. The results of the analysis are interpreted in terms of both the HLZ and the GRW parameter conventions.

The effective theory breaks down at energy scales at which the underlying theory of quantum gravity starts to affect the phenomenology. We assume that the range of validity is characterized by a value $\sqrt{\hat{s}_{\max}}$, roughly corresponding to the mass $M_{\max}$ of the lepton pairs emitted in the decay of the graviton. As no clear prediction for $\sqrt{\hat{s}_{\max}}$ can be made within the ADD model, and to take into account the requirement $\sqrt{\hat{s}_{\max}} \ll \Lambda$, most results in this Letter are presented both for $M_{\max} = M_s$ and for a range of different values of $M_{\max}$.

Constraints on virtual-graviton signatures in the ADD model of extra dimensions have been obtained at HERA [8,9], LEP [10–15], and the Tevatron [16,17]. At the LHC, limits have been presented based on measurements of diphoton events [18–20].

CMS uses a right-handed coordinate system with axes labeled x, y, and z, and the origin at the center of the detector. The z-axis points along the direction of the anticlockwise beam. The azimuthal and polar angles are ϕ and θ, with θ measured from the positive z-axis. The pseudorapidity η is defined by $\eta = -\ln\tan(\theta/2)$.

A main feature of the CMS apparatus is a superconducting solenoid of 6 m internal diameter, providing a magnetic field of 3.8 T. Located within the field volume are silicon pixel and strip inner trackers, an electromagnetic calorimeter (ECAL), and a hadronic calorimeter (HCAL). The ECAL consists of lead-tungstate crystals covering pseudorapidities of $|\eta| < 1.5$ (barrel) and $1.5 < |\eta| < 3.0$ (endcaps). The CMS muon detectors are embedded in the return yoke of the magnet. Muons are measured with detection planes using three different technologies: Drift Tubes, Cathode Strip Chambers, and Resistive Plate Chambers. The first stage of the CMS trigger system employs custom hardware and processes information from the calorimeters and the muon system. The event rate is further reduced by a computer farm using the event information from all detector systems. A detailed description of CMS can be found in Ref. [21].

This analysis uses data samples collected with the CMS detector in 2011, corresponding to an integrated luminosity [22] $\mathcal{L}$ of 2.3 ± 0.1 fb^{-1} (dimuons) or 2.1 ± 0.1 fb^{-1} (dielectrons). The integrated luminosity for the dimuon channel is larger because the muon selection has less stringent requirements on the performance of the calorimeters during data-taking. The muon data sample was collected using a single-muon trigger with a transverse momentum (p_T) threshold which was varied between 15 and 40 GeV over the course of data-taking to allow for changes in instantaneous luminosity. The selection of electron events is based on a trigger requiring the presence of 2 electrons or photons with energy depositions > 33 GeV. Candidate events are required to have a reconstructed interaction vertex with $|z| < 24$ cm, and a radial distance $\sqrt{x^2 + y^2} < 2$ cm. For events passing the complete selection requirements, the trigger efficiencies for signal and SM Drell–Yan (DY) events with large mass are $> 99\%$, with an uncertainty of $< 1\%$.

Muons with $|\eta| < 2.1$ and $p_{T,\mu} > 45$ GeV are selected. The candidates are required to be identified both in the outer muon system and the inner tracker, and the inner track must contain reconstructed energy deposits in at least 1 pixel layer and more than 10 strip-tracker layers. Muon candidates are required to have signals from at least two muon detector layers included in the reconstructed muon track. Muon candidates satisfy the isolation requirement $\sum p_T^i / p_{T,\mu} < 0.1$, where the sum extends over the momenta p_T^i of all charged particle tracks (excluding the muon track) within a cone of size $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.3$ around the muon direction of flight. To reject backgrounds from cosmic-ray muons, we require a transverse impact parameter relative to the primary vertex of < 0.2 cm, and an opening angle of $\alpha_{\mu\mu} < \pi - 0.02$ between the 2 muon momentum vectors. No charge requirement is applied to the muon pairs. However, all selected muon pairs of mass > 450 GeV are found to have opposite charges. Events with 2 muons passing the selection criteria are accepted for analysis.

Electron candidates are reconstructed from energy depositions in the ECAL (superclusters) matched to a track in the silicon tracker. ECAL superclusters are constructed from 1 or more clusters of energy depositions surrounding the crystal with the highest local energy deposition. An associated track is required to contain signals from at least 5 tracker layers. The track must be matched geometrically to the supercluster, and the spatial distribution of energy must be consistent with that expected for an electron. Only electron candidates with a ratio of energy depositions in the HCAL and ECAL below 0.05 are considered. To minimize the contamination from jets, electron candidates are required to be isolated. Candidates with a sum of transverse track momenta $\geqslant 5$ GeV within $0.04 < \Delta R < 0.3$ around the candidate track are rejected. In the ECAL and inner HCAL layer, the deposited transverse energy E_T in a cone $\Delta R = 0.3$ around the electron candidate (excluding the transverse energy $E_{T,e}$ of the electron) must be < 2 GeV $+ 0.03 \times E_{T,e}$ for the barrel, or < 2.5 GeV (< 2.5 GeV $+ 0.03 \times (E_{T,e} - 50$ GeV$)$) for the endcaps if $E_{T,e} < 50$ GeV ($E_{T,e} \geqslant 50$ GeV). Additionally, the E_T deposition in the outer HCAL layer within $0.15 < \Delta R < 0.3$ around the electron position is restricted to < 0.05 GeV. Selected events must contain 2 electrons of opposite charge, each with transverse energy $E_T > 35$ GeV (in the barrel region), or with $E_T >$ 40 GeV (in the endcaps). The explicit charge requirement is found to have negligible influence on the presented results. Events in which both electrons are reconstructed in the endcaps are not used in the analysis, since electrons from the ADD signal would on average be produced at smaller values of η than the SM backgrounds.

The search is performed with a set of events that contains either electron or muon pairs above a mass value $M_{\min}$. The lower bound of the signal region is chosen to maximize the expected upper limits of the ADD model parameter Λ_T in each lepton channel. The optimum value of $M_{\min}$ is found to be 1.1 TeV for both the dimuon and the dielectron channel, based on simulation studies.

In both search channels, the PYTHIA 8.142 [23,24] event generator with the MSTW08 [25] parton distribution function (PDF) set is used to simulate the expected signal. Interference terms between the standard model DY process and the virtual graviton are taken into account in the evaluation of the signal cross sections. Simulated events for both signal and SM backgrounds are passed through a detailed detector simulation based on GEANT4 [26], using a realistic CMS alignment scenario, and the same reconstruction chain as data.

In this analysis, the SM DY process is the dominant background. In the dimuon channel, we use the MC@NLO [27,28] event generator with the CTEQ6.6 [29] PDF set to simulate the DY background. The parton level events from MC@NLO are passed to HERWIG 6 [30] for the simulation of the QCD parton shower and hadronization, PHOTOS [31] for the simulation of the electroweak (EW) parton shower, and JIMMY [32] for the simulation of multiple parton interactions. The simulated reconstruction efficiencies in the chosen region of acceptance, including all selection criteria, are found to be $90\% \pm 3\%$ for the high-mass DY dimuon background and $90\% \pm 4\%$ for the ADD dimuon signal.

Mass-dependent corrections [33] beyond the QCD next-to-leading-order (NLO) predictions implemented in MC@NLO are studied to improve the SM DY estimate in the dimuon channel. EW NLO effects are evaluated by comparing HORACE [34] NLO predictions interfaced to HERWIG 6 with HORACE LO predictions interfaced