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## Physics Letters A

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## Vanadium sulfide nanoribbons: Electronic and magnetic properties



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#### ARTICLE INFO

Article history:
Received 24 April 2013
Received in revised form 23 September 2013
Accepted 29 September 2013
Available online 4 October 2013
Communicated by R. Wu

Keywords: First-principles calculation Vanadium sulfide nanoribbon Magnetism

#### ABSTRACT

We report the structural, electronic and magnetic properties of zigzag-type 2H-VS<sub>2</sub> nanoribbons based on the first-principles calculations. Our results suggest that the zigzag-type 2H-VS<sub>2</sub> nanoribbons are intrinsic ferromagnetic or ferrimagnetic materials dependent on their edge structures. The S-terminated VS<sub>2</sub> nanoribbons with or without hydrogen saturation at the edges are ferromagnetic, whereas V-terminated VS<sub>2</sub> nanoribbons are ferrimagnetic at their ground states. The average magnetic moment per V atom of VS<sub>2</sub> nanoribbons increases monotonously with their width, but still smaller than that of perfect VS<sub>2</sub> monolayer. These results imply the great potential of VS<sub>2</sub> nanoribbons in spintronics application.

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#### 1. Introduction

Developing spintronics with two-dimensional monolayer materials have attracted much research attention recently to meet the need of miniaturization of devices and massive information storage [1–7]. Much effort has been devoted to graphene and graphene-like materials, such as BN and ZnO monolayer. Although these materials are usually nonmagnetic, intensive studies have predicted that graphene, BN, or ZnO monolayer may possess magnetic properties via various physical or chemical methods [8–13]. For instance, the graphene nanoribbon (NR) with zigzag-type edges may be tuned into half metal by applying external electric field or chemically modified with different functional groups at both edges [8,9]. Introducing native defect, doping, or creating interface with substrate may also induce various magnetic properties in graphene, BN, ZnO, MoTe<sub>2</sub>, MoSe<sub>2</sub>, and WS<sub>2</sub> monolayer [10–16].

As a member of layered transition-metal dichalcogenide compounds, VS<sub>2</sub> has attracted increasing research attention since the first discovery in 1970s for its novel properties and promising applications in spintronics and energy storage [17–24]. There are two possible phases of VS<sub>2</sub> crystal, i.e. 1T- and 2H-phase, where 1T-phase has been observed experimentally with various interesting properties, such as coexistence of the charge density wave state and Mott transition [20–24]. In VS<sub>2</sub> compound, one layer of metal V is sandwiched between two layers of S, and the triple layers stack together through weak Van der Waals interaction. Thus, the

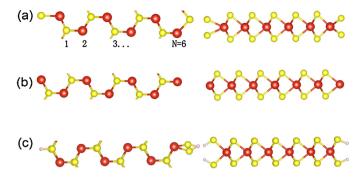
graphene-like VS<sub>2</sub> monolayer may be realized by exfoliating VS<sub>2</sub> crystal. Recently, less than five VS<sub>2</sub> triple layers have been synthesized by intercalation compound precursor of VS<sub>2</sub>.NH<sub>3</sub> into VS<sub>2</sub> ultrathin nanosheet [6]. Interestingly, the theoretical simulation has predicted that 2H-VS<sub>2</sub> monolayer is magnetic metal, suitable for spintronics applications [25].

In this Letter, we report the electronic and magnetic properties of 2H-VS<sub>2</sub> NRs based on the first-principles calculations. Both V and S-terminated zigzag-type VS<sub>2</sub> NRs are considered. Our results suggested that zigzag-type VS<sub>2</sub> NRs are intrinsic ferromagnetic or ferrimagnetic materials dependent on their edge structures. These results open a new door to design VS<sub>2</sub> monolayer based spintronic devices.

#### 2. Materials and methods

All calculations are performed using the Vienna *ab initio* Simulation (VASP) package [26,27]. The frozen-core all-electron projector augmented wave (PAW) method was used to describe electron—ion interaction during the calculation, with the generalized gradient approximation with the PW91 functional for exchange-correlation energy [28–30]. A kinetic-energy cutoff of 400 eV is chosen for the plane wave expansion of the electronic wave function. To simulate one-dimensional VS<sub>2</sub> NRs with infinite length, a tetragonal supercell was used with two adjacent NRs separated by a vacuum region of at least 13 Å. A K-point mesh of  $1 \times 1 \times 8$  grid was employed for geometry optimization, and  $1 \times 1 \times 18$  grid for static electronic structure calculations. The axial lattice constant and atomic positions are fully relaxed until the force on each atom was less than 0.02 eV/Å, and the convergence criteria for energy  $10^{-5}$  eV.

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**Fig. 1.** The optimized structure of zigzag-type  $VS_2$  NR terminated with (a) S, (b) V, or (c) hydrogenated S is illustrated. The red, yellow and small white balls represent V, S, and H atoms, respectively. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this Letter.)

#### 3. Results

Fig. 1 shows the optimized structures of three examples of 2H-VS<sub>2</sub> zigzag-type NRs terminated with V, S, or hydrogenated S atoms, namely VS<sub>2</sub>-V, VS<sub>2</sub>-S, and VS<sub>2</sub>-H NRs, respectively. The width of ribbon is chosen as N = 6, where N indicates the number of zigzag-type atomic lines along the ribbon's direction, as shown in Fig. 1(a). The optimized lattice constant along the ribbon's direction is about 3.19 Å, which is in good agreement with that of previous work [6]. The V-S bond lengths in VS<sub>2</sub> NRs vary remarkably depending on the position of atoms due to the reconstruction of their edge. For example, the V-S bond lengths within the VS<sub>2</sub>-S NR (N = 6) range from 2.36 to 2.38 Å, whereas range from 2.25 to 2.35 Å on the edge. For  $VS_2$ -V NR (N = 6), the V-S bond lengths range from 2.16 to 2.44 Å on the edge. The difference of V-S bond lengths for VS<sub>2</sub>-V and VS<sub>2</sub>-N NRs mainly originates from the number of dangling bonds at their edges. For instance, the edged V atom in the VS<sub>2</sub>-V NR has two dangling bonds, whereas the edged S atom in the VS<sub>2</sub>-S NRs has one dangling bond when cutting from VS<sub>2</sub> monolayer. The optimized S-H bond length of VS<sub>2</sub>-H NR (N = 6) is about 1.36 Å. Also, the widths of zigzag-type VS<sub>2</sub> NRs depend on their edge structure for same N. The optimized widths of  $VS_2$ -V,  $VS_2$ -S, and  $VS_2$ -H NRs (N = 6) are 15.65, 16.87 and 18.68 Å, respectively.

To explore the electronic properties of VS<sub>2</sub> NRs, we calculated the total density of states (DOS) and partial DOS of V's 3d orbital and S's 2p orbital, as illustrated in Fig. 2. Three examples of zigzagtype VS<sub>2</sub> NRs are magnetic metals, independent of the chemical characteristic of their edges. The V's 3d orbital and S's 2p orbitals contribute significantly to the total DOS in the displayed energy windows of [-6.0, 6.0] eV around the Fermi energy level, suggesting the strong hybridization of V's d orbital and S's p orbital. It is

also evident that the magnetism of  $VS_2$  NRs is mainly contributed by the V atoms, whereas S atoms only have a little contribution to the total magnetic moment.

In Fig. 3, we plotted the profiles of spin charge density distribution of three  $VS_2$  NRs with N = 6. The spin charge density mainly distributes on V atoms, consistent with the calculated DOS. However, VS<sub>2</sub> NRs present an edge-dependent magnetic property. For VS<sub>2</sub>-S and VS<sub>2</sub>-H NRs, the ground state is ferromagnetic. The V atoms at the center region of NRs play an important role to the total magnetism, where the V atoms at the edge have a little contribution to the total magnetism. This behavior is different from that of zigzag graphene NR, where the edged carbon atoms have the largest contribution to the total magnetism [8]. The VS<sub>2</sub>-V NR is ferrimagnetic at its ground state, as shown in Fig. 3(b). The total energies of its ferromagnetic, antiferromagnetic, and nonmagnetic states are higher than those of its ferromagnetic state of about 5.2, 12.0, and 254.3 meV per supercell. The V atoms neighbor to the edged V atoms are coupled in antiferromagnetic order with other V atoms. Different from those in VS2-S and VS2-H NRs, the edged V atoms in VS2-V NR have the largest contribution to the total magnetism. Both H and S atoms have negligible contribution to the magnetism.

At last, the effect of ribbons' widths on their magnetic properties has been investigated. As shown in Fig. 4, the spin charge density distribution profiles of VS2-S, VS2-V, and VS2-H NRs with various widths are plotted. Clearly, the magnetic orders of VS<sub>2</sub> NRs are not sensitive to their widths. For instance, the VS2-S and  $VS_2$ -H NRs with N=4 and 10 are ferromagnetic at their ground states, and have similar spin charge density distributions as the ribbon with N = 6. The V atoms at the center regions have the major contribution, whereas the edge ones have minor contribution to the total magnetism. The VS2-V NRs with various widths (N = 4, 5, 7, 8, 10) are ferrimagnetic at their ground states and present the same magnetic order, i.e. the V atoms neighbor to the edged one are coupled with other V atoms in antiferromagnetic order. The edged V atoms present the largest contribution to the total magnetism of VS2-V NRs. Note here we mainly studied the magnetic properties of zigzag-type VS<sub>2</sub> NRs with V, S or H termination. Test calculations were also performed on the zigzag-type VS<sub>2</sub> NRs with both V and S edge (VS2-S-V), as well as the armchair-type VS2 NRs. As shown in Fig. 5, the zigzag-type VS2-S-V NR is ferrimagnetic at its ground state with similar spin charge distribution as that of zigzag-type VS<sub>2</sub>-V NR at V edge. The armchair-type VS<sub>2</sub> NR is ferromagnetic at its ground state and spin charge density mainly locates on the V atoms within the ribbon.

In Fig. 6, we plotted the average magnetic moment of V atom for  $VS_2$  NRs with different edge structures and widths. The average magnetic moment of V atom is defined as the total magnetic moment of  $VS_2$  NR divided by the number of V atoms. For

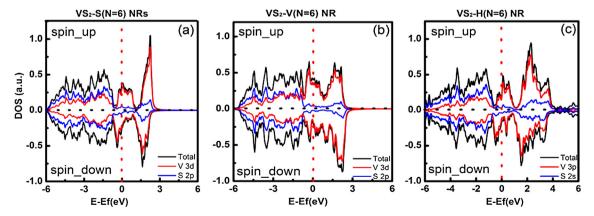


Fig. 2. The total DOS and partial DOS of V's 3d orbital and S's 2p orbital are plotted for (a) VS2-S, (b) VS2-V, and (c) VS2-H NRs (N = 6), respectively.

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