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## Magnetic properties and magnetocaloric effect of Sr-doped  $Pr_{0.7}Ca_{0.3}MnO_3$ compounds

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#### 1. Introduction

Recently, perovskite manganites  $Ln_{1-x}A_xMnO_3$  (Ln = La, Pr, Nd ...;  $A = Ca$ , Sr, Ba ...; be referred as manganites) have been an interested topic in the solid state physics community due to their rich electrical magnetic properties [\[1](#page--1-0)–4]. The close relationship between structural, electrical, and magnetic properties has opened new aspects in studying of manganites in the fundamental and application viewpoints. It has been reported  $LnMnO<sub>3</sub>$  as an insulator antiferromagnetic (anti-FM) with Néel temperature to be about 135 K [[5\]](#page--1-1). A partial replacement of  $Ln^{3+}$ by  $A^{+2}$  introduces numbers of Mn<sup>3+</sup> ions to be converted to Mn<sup>4+</sup> ions, leads to a mixed valence state  $Mn^{3+}/Mn^{4+}$ . Therefore, the double-exchange interactions of  $Mn^{3+}-Mn^{4+}$  ions appeared. Material thus exhibits the ferromagnetic (FM)-paramagnetic (PM) phase transition taking place at Curie temperature  $(T_C)$ . Changing the  $Ln^{3+}/A^{2+}$  ratio, the  $Mn^{3+}/Mn^{4+}$  ratio will be modified, and affect structural and physical properties of materials [\[6](#page--1-2)–8]. Basically, electrical and magnetic properties of manganites can be explained by such mechanisms i) the coexistence and competition of the  $Mn^{3+}-Mn^{4+}$  double-exchange interactions and the super-exchange interactions of  $Mn^{3+}-Mn^{3+}$  and  $Mn^{4+}-Mn^{4+}$  pairs [\[6](#page--1-2)[,9\]](#page--1-3), ii) the strong electron-phonon coupling (called the Jahn-Teller effect) [[10,](#page--1-4)[11\]](#page--1-5), and iii) the change of the average cationic radius [\[8,](#page--1-6)[12\]](#page--1-7).

pounds around their T<sub>C</sub> have been also analyzed. Results suggested a coexistence of the ferromagnetic short- and long-range interactions in samples. Moreover, Sr-doping favors establishing the short-range interactions.

> Motivated by the fact that several manganites exhibit a discontinuous variation in the structure and a sharp magnetic phase transition, which is known as the first-order phase transition (FOPT) [13-[15\]](#page--1-8). Among them,  $Ln_{0.7}A_{0.3}MnO_3$  systems usually exhibit the strongest physical effects [[9](#page--1-3)[,16](#page--1-9)–18]. However, the FM-PM phase transition in La<sub>0.7</sub>Ca<sub>0.3</sub>MnO<sub>3</sub> (Ln = La) single-crystal or polycrystalline bulk is the first-order  $[12,19-22]$  $[12,19-22]$  $[12,19-22]$  $[12,19-22]$ , whereas, other cases (e.g.  $A = Sr$ , Ba, Pb) are the second-order phase transition (SOPT) [[12,](#page--1-7)[22](#page--1-11)–24]. Studies on the critical properties around the FM-PM transition of  $La_{0.7}Ca_{0.3}$ .  $x(Sr, Ba)xMnO_3$  have shown that the partial replacement of Ca<sup>2+</sup> by  $Sr^{+2}$  results in the change from the long-to short-range FM order [\[12](#page--1-7)]. This is the opposite for the case of  $Ca^{2+}$  replacement with  $Sr^{2+}$ [[12](#page--1-7)[,22](#page--1-11)[,24](#page--1-12)].

> The previous reports have pointed that  $Pr_{1-x}A_xMnO_3$  compounds  $(Ln = Pr)$  are interesting systems with a rich phase diagram  $[25-28]$  $[25-28]$ . For  $A = Ca$ ,  $Pr_{1-x}Ca_xMnO_3$  exhibits a FM insulator in the range of  $x = 0-0.3$  and the anti-FM charge-ordering behavior when  $x = 0.3-0.5$ .

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Besides, this system also reveals a coexistence of the FM and anti-FM interactions, and spin-glass behaviors in samples with low Ca-doping level  $(0.0 < x < 0.2)$  [\[27](#page--1-14)]. A non-monotonic change of Curie temperature has been also observed in the range of  $0.2 < x < 0.3$ , in which a maximum value of  $T_c$  appears at  $x = 0.25$  [[26\]](#page--1-15). The complex magnetic properties in this system can be understood via detailed analyses of the critical behaviors. A recent report indicated that the critical exponents ( $\beta$  and  $\gamma$ ) of Pr<sub>1-x</sub>Ca<sub>x</sub>MnO<sub>3</sub> compounds are very sensitive to Ca-doping concentration ( $\beta = 0.351$  and  $\gamma = 1.372$  for  $x = 0.25$ ,  $\beta = 0.521$  and  $\gamma = 0.912$  for  $x = 0.29$  [\[29](#page--1-16)]. It suggested that Ca-doping favors establishing the FM long-range order in this system. For  $A = Sr$ . Knizek et al. [\[28](#page--1-17)] reported that  $Pr_{1-x}Sr$  MnO<sub>3</sub> with  $0.3 \le x \le 0.4$  is a FM conductor at low temperature, whereas  $Pr_{0.5}Sr_{0.5}MnO_3$  is a charge-ordered anti-FM. However, using the isothermal magnetization  $M(H, T)$  data, Phan et al. [[30\]](#page--1-18) indicated a large magnetocaloric effect (MCE) in  $Pr_{0.63}Sr_{0.37}MnO_3$  single crystal, which exhibits the negative slopes of the  $H/M$  versus  $M^2$  plots. This reflects that  $Pr_{0.63}Sr_{0.37}MnO_3$  single crystal is not a conventional ferromagnet, which is known as the FOPT [[30](#page--1-18)]. In our previous report [\[31](#page--1-19)], we observed an opposed behavior for  $Pr_{1-x}Sr_xMnO_3$  compounds with  $x = 0.3 - 0.45$ . All compounds in the range of Sr-doping undergo the SOPT with the critical exponent  $\beta$  value decreasing monotonically as Sr concentration increasing [[31\]](#page--1-19). Presently, very few studies regard on the influences of A-site substitution on the MCE and the critical behaviors of  $Pr_{0.7}Ca_{0.3}MnO_3$  compound. The electrical magnetic properties of Ca and Sr co-doping in  $Pr<sub>0.65</sub>A<sub>0.35</sub>MnO<sub>3</sub>$  compound have been reported by Biswas et al. [\[32](#page--1-20)]. They pointed out that  $Pr_{0.65}(Ca_{0.7}Sr_{0.3})_{0.35}MnO_3$ compound has two large MCEs taking place at the charge order transition ( $T_{\text{CO}}$  = 220 K) and the spontaneous destabilization of charge ordered state ( $T_M$  = 120 K). In the earlier works, the electronic, magnetic, and thermal properties of  $Pr_{0.65}Ca_{0.35-x}Sr_xMnO_3$  compounds have been done [[33,](#page--1-21)[34\]](#page--1-22). Mollah et al. [\[33](#page--1-21)[,34](#page--1-22)], reported a phase separate induced significant enhancement of magnetoresistance 10% at 100 Oe and 99% at 5 kOe around the metal-insulator transition. However, detailed analyses related to the magnetic critical region, and the magnetic interactions in these compounds have not been carried out yet and therefore this issue needs a more deep assessment.

To further understand the influences of Ca and Sr co-substitution, we prepared  $Pr_{0.7}Ca_{0.3-x}Sr_xMnO_3$  with  $x = 0.0-0.3$  and investigated their critical properties in the vicinity of the magnetic phase transition. Basing on a phenomenological model proposed by Hamad [\[35](#page--1-23)] [\[36](#page--1-24)], and using magnetization versus temperature  $M(T)$  data measured at several applied fields, we have predicted the MCE and the FM orders of  $Pr_{0.7}Ca_{0.3-x}Sr_xMnO_3$  compounds through their typical parameters, such as the magnetic entropy change  $(\Delta S_m)$ , the relative cooling power (RCP), the variation of the specific heat change ( $\Delta C_p$ ), and the critical exponent  $\beta$  corresponding to the spontaneous magnetization ( $M_{\rm S}$ ). The results obtained reveal all the samples undergoing the SOPT together high values of  $\Delta S_{\text{m}}$ . Further, the FM short-range order in Pr<sub>0.7</sub>Ca<sub>0.3-</sub>  $xSr_xMnO_3$  compounds is favored when Sr concentration increases.

#### 2. Experimental

Four compounds of  $Pr_{0.7}Ca_{0.3-x}Sr_xMnO_3$  with  $x = 0.0, 0.1, 0.2,$  and 0.3 were prepared by a solid-state reaction method. Firstly, stoichiometric masses of  $Pr<sub>6</sub>O<sub>11</sub>$ , CaCO<sub>3</sub>, SrCO<sub>3</sub> and Mn powders with high purity (99.9%) were ground and mixed, and then calcinated at 1200 °C in air for 24 h. After that these mixtures were reground and pressed into pellets under a pressure 5000 kg/cm<sup>2</sup>, and then sintered at 1300 °C in air for 48 h. X-ray diffraction (XRD) pattern was recorded at room temperature on an X-ray diffractometer Bruker AXS (D8 Discover) using a Cu- $K_{\alpha}$  radiation. The Win-metric software was used to calculate the lattice parameters. Temperature dependences of magnetization under various magnetic fields ranging  $H = 0-10$  kOe were measured on a vibrating sample magnetometer (VSM), according to the increasing direction of temperature with an increment of 2 K.  $T_c$  value of samples

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Fig. 1. Room-temperature XRD patterns corresponding to the Miller indices for  $Pr_{0.7}Ca_{0.3-x}Sr_xMnO_3$  compounds.

were determined via the second derivative of magnetization versus temperature  $(d^2M/dT^2 \text{ vs. } T)$  [\[37](#page--1-25)].

#### 3. Results and discussion

Powder XRD patterns of  $Pr_{0.7}Ca_{0.3-x}Sr_xMnO_3$  compounds illustrated in [Fig. 1](#page-1-0) present sharp and intense XRD peaks. These are suitable for a PDF card No. 49–0461 [\[38](#page--1-26)] of the XRD standard database. The Millerindexes of all the XRD peaks belong to an orthorhombic structure with space group Pbnm. This suggests that all samples are single phase, without any trace of secondary phase. The lattice parameters  $(a, b, c, c)$ and V) were calculated and presented in [Table 1.](#page-1-1) It shows that the lattice parameters slightly increase with Sr-doping concentration. The variation of the lattice parameters could be associated to a substitution of  $Sr^{2+}$  ion with a larger ionic radius  $(r_{Sr2+} = 1.44 \text{ Å})$  for  $Pr^{3+}/Ca^{2+}$ ions with the smaller ionic radii ( $r_{\text{Pr3+}} = 1.179 \text{ Å}$  and  $r_{\text{Ca2+}} = 1.34 \text{ Å}$ ) located at A-site in  $ABO<sub>3</sub>$  structure. This also affects the average ionic radius at A-site  $( $r_A >$ ), the dispersion of A-cation radius$  $(\sigma^2 = \sum y_i r_i^2 - \langle r_A \rangle^2)$ , and the Goldsmith tolerance factor  $(\tau_G = (K_H + r_O)/\sqrt{2} (r_B + r_O)$ , with  $r_A$ ,  $r_B$ , and  $r_O$  are the ionic radii located at A-, B-, and O-sites in ABO<sub>3</sub> structure, respectively). It is also known that if  $\tau$ <sub>G</sub> value is in the range of 0.75 <  $\tau$ <sub>G</sub> < 1, ABO<sub>3</sub> compounds has a perovskite structure. These compounds will belong to an orthorhombic structure if  $\tau$ <sub>G</sub> < 0.969. And in an ideal case, their structure is a cubic system corresponding to  $\tau$ <sub>G</sub> = 1 [[39\]](#page--1-27). For our case, the values of  $\langle r_A \rangle$ ,  $\sigma^2$ , and  $\tau_G$  were calculated and listed in [Table 1](#page-1-1). Once can see that these parameters gradually increase with Sr concentration. Among these, the value of  $\tau$ <sup>G</sup> increase from 0.924 for  $Pr_{0.7}Ca_{0.3}MnO_3$  (x = 0.0) to 0.935 for  $Pr_{0.7}Sr_{0.3}MnO_3$  (x = 0.3), suggesting that the distortion in the structure of  $Pr_{0.7}Ca_{0.3-x}Sr_xMnO_3$  decreases with increasing Sr content.

The field-cooled magnetization versus temperature curves, M(T), measured at 100 Oe of samples are performed in [Fig. 2\(](#page--1-27)a). A magnetic phase transition from FM state at low temperature region to PM state at





<span id="page-1-1"></span>Table 1

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