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Bayesian statistical modeling of microcanonical melting times at the superheated regime



PHYSICA

STATISTICAL MECHANIC

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ABSTRACT

Homogeneous melting of superheated crystals at constant energy is a dynamical process, believed to be triggered by the accumulation of thermal vacancies and their self-diffusion. From microcanonical simulations we know that if an ideal crystal is prepared at a given kinetic energy, it takes a random time t_w until the melting mechanism is actually triggered. In this work we have studied in detail the statistics of t_w for melting at different energies by performing a large number of Z-method simulations and applying state-of-the-art methods of Bayesian statistical inference. By focusing on a small system size and short-time tail of the distribution function, we show that t_w is actually gamma-distributed rather than exponential (as asserted in a previous work), with decreasing probability near $t_w \sim 0$. We also explicitly incorporate in our model the unavoidable truncation of the distribution function due to the limited total time span of a Z-method simulation. The probabilistic model presented in this work can provide some insight into the dynamical nature of the homogeneous melting times from simulation into the Z-method in order to correct the effect of short simulation times.

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1. Introduction

It is well known that, under carefully controlled conditions, a solid can be heated above its melting temperature without triggering the melting process. This is known as superheating, and in it the solid enters a metastable state with interesting kinetic properties [1]. It is also possible to achieve such an homogeneous melting by ultrafast laser pulses [2]. As a complement to experiments, observation and characterization of this superheated state is especially clear in microcanonical simulations, particularly within the framework of the Z-method [3]. By using this methodology, it has been possible to establish the role of thermal vacancies [4,5] and their self-diffusion [6–8].

The Z-method, in fact, relies on the superheating metastable state to determine the melting temperature T_m . It establishes that in microcanonical simulations (i.e. with constant volume, constant internal energy, constant number of particles) there is a maximum energy E_{LS} (*LS* stands for Limit of Superheating), related to a temperature T_{LS} , that can be given to the crystal before it spontaneously melts. Beyond this point the solid unavoidably melts, showing a sudden, sharp drop in temperature, corresponding to the kinetic energy consumed as latent heat of melting. When a system melts at an energy $E_{LS} + \Delta E$ (with $\Delta E \rightarrow 0$), the final temperature *T* approaches T_m . Despite the current lack of a complete theory of first-order phase

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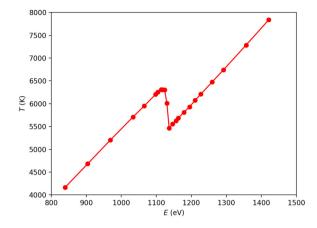


Fig. 1. Isochoric (Z) curve obtained from 23 different 50 ps simulations of high-density Ar, with initial temperatures T_0 ranging from 8000 to 17000 K. From this Z-curve we can determine $T_m \sim 5465$ K and $T_{LS} \sim 6307$ K. A clear drop of temperature is observed between $T_0 = 12500$ K and $T_0 = 12600$ K.

transitions and metastable states, there is plenty of accumulated evidence of the accuracy of the Z-method [9–11]. However, some issues remain concerning the uncertainties associated to T_m and T_{LS} due to the use of short simulation times. One important source of uncertainty is the random distribution of the elapsed time t_w from the beginning of the simulation until the melting process is triggered, even in a microcanonical setting. In previous work, Alfè et al. [12] reported statistics of t_w obtained from molecular dynamics (MD) simulations and postulated that t_w is exponentially distributed, and therefore the most probable value of t_w is close to zero for all temperatures.

In this work, we performed a large number of microcanonical simulations of a Lennard-Jones solid in order to generate statistically independent samples of t_w in a wide range of initial temperatures. We focused particularly in shorter times and a small system size and found that these "waiting times" are not exponentially distributed but gamma-distributed. Accordingly, we present a precise model that reproduces the gamma shape and scale parameters as functions of the initial temperature.

The rest of the paper is organized as follows. In Section 2 we give a detailed description of the simulations that were used to determine the melting temperature and the samples of waiting times (t_w). Section 3 describes the probability models proposed to represent the statistical distribution of t_w and assesses their likelihood given our simulated data. Next, Sections 4 and 5 present a concrete model for the distribution of t_w as a function of temperature fitted to our simulation data. Finally, Section 6 discusses the implications of our results for the process of melting in the superheating state for homogeneous solids.

2. Computational procedure

As a simple model of solid we considered the Lennard-Jones potential with $\epsilon/k_B = 119.8$ K and $\sigma = 3.41$ Å, as used in previous works on the Z-method [3,4]. A high-density, face-centered cubic (FCC) crystal, having 500 atoms and lattice constant a = 4.2 Å, was used as the starting point for all simulations. The use of such a high density (and therefore pressure) increases the superheating effect, revealing the melting process in a more salient way. All simulations were performed using the LPMD [13] molecular dynamics package. For every case we use a total simulation time $\tau_0 = 50$ ps and a timestep $\Delta t = 0.5$ fs with a Beeman integrator for Newton's equations of motion. Our aim is to completely characterize the statistical distribution of waiting times t_w for a small system size N and limited window of time, so that we can perform predictions for larger simulation times at the same N.

A set of 400 simulations with temperatures ranging from 50 K to 20000 K have been used to draw an isochoric T(E) curve for this high-density argon model, in a standard application of the Z-method. We will refer to this T(E) curve as the Z-curve, although the same name has been used in previous works to refer to the isochoric T(P) curve. We prefer the use of the internal energy as the independent variable instead of the pressure is to be preferred because of the possibility of a non-monotonic trend of pressure and temperature, as observed for instance in the case of lithium [14]. In those cases where the melting curve has a region with negative $(\partial T / \partial P)_E$, the T(P) isochore reverts direction while the T(E) isochore is unaffected, as C_v is non-negative for solid and liquid equilibrium branches.

These results are shown in Fig. 1.

From the isochoric T(E) curve, we have estimated the melting temperature T_m and the superheating limit T_{LS} . The sharp inflection at the higher temperature (6307 K) corresponds to T_{LS} and the lower inflection (5465 K) corresponds to T_m .

Once T_{LS} is estimated, and in order to collect statistically independent samples of waiting times, we performed a large number of classical MD runs for each of the 24 different initial temperatures T_0 in the interval from $T_0 = 12075$ K to $T_0 = 13300$ K. This was done as follows. For the temperatures between 12000 K and 12250 K, 10000 simulations per

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