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# A new dodecanuclear manganese single-molecule magnet from the arrangement of manganese triangles

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#### ABSTRACT

Bridging-azido ligands are used to create a new dodecanuclear manganese cluster  $[Mn_{12}O_4(salox)_{12}(N_3)_4$ (MeOH)<sub>6</sub>(H<sub>2</sub>O)<sub>2</sub>]·5MeOH·2H<sub>2</sub>O (**1**·5MeOH·2H<sub>2</sub>O) from the arrangement of manganese triangles,  $[Mn_3O]$ . Single-crystal x-ray analysis shows that **1**·5MeOH·2H<sub>2</sub>O contains a  $[Mn^{III}_{12}(\mu_3-O)_4(\mu_2-N_3)_4]^{24+}$  core, constructed from four  $[Mn^{III}_3O]^{7+}$  building units, which is a new structural type in the family of  $Mn_{12}$  complexes. Variable temperature magnetic susceptibility and magnetization measurements of the complex **1**·5MeOH·2H<sub>2</sub>O were carried out. The magnetic data indicate that complex **1**·5MeOH·2H<sub>2</sub>O exhibits antiferromagnetic with a non-zero ground state spin value of S = 4. The frequency dependence of out-of-phase component in alternating current magnetic susceptibility and magnetic hysteresis loop measurements indicate that complex **1**·5MeOH·2H<sub>2</sub>O is a new single-molecule magnet with a significant magnetoaniso-tropy of  $U_{eff} = 70$  K.

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Polynuclear manganese clusters often exhibit large, and sometimes abnormally large, spin values in the ground state, and combined with a large magnetic anisotropy have led some of these species to be single-molecule magnets (SMMs). SMMs are attracting extensive attention because they represent nanoscale magnetic particles of a well-defined size [1,2]. They display sluggish magnetization relaxation phenomena such as magnetization hysteresis loops and frequency-dependent out-of-phase alternating current (AC) magnetic susceptibility [3,4]. The remarkable magnetic properties of an SMM arise from its high-spin ground state (S) split by a large negative axial zero-field splitting (ZFS, D) which results in an anisotropy energy barrier of  $KV = |D|S_Z^2$  [5]. The development of rational syntheses of multinuclear metal complexes, especially containing high-spin manganese(III) ions, has been magnetically interesting due to the pronounced magnetic anisotropy provided by the Jahn-Teller (JT) distortion. Thus, many current routes to multinuclear SMMs culminate in manganese(III) containing complexes, to achieve the associated large single-ion anisotropy [6,7]. The use of salicylaldoxime ligands in manganese chemistry has proven to be extremely successful in the synthesis of new polynuclear complexes [8–12]. Some SMMs and single-chain-magnets (SCMs), using salicylaldoxime (H<sub>2</sub>salox) and its derivatives, were reported [13–17]. Moreover, the azide bridging ligand often mediates ferromagnetic exchange between paramagnetic centers, and this property has been recently exploited in the preparation of some SMMs [18–20] and SCMs [21–24]. In the present study, a polynuclear manganese complex with azide and salicylaldoximate, which produces a new dodecanuclear manganese cluster,  $[Mn_{12}O_4(salox)_{12}(N_3)_4(MeOH)_6(H_2O)_2] \cdot 5-MeOH \cdot 2H_2O$  (1 $\cdot$ 5MeOH  $\cdot 2H_2O$ ), is investigated. Complex 1 $\cdot$ 5MeOH  $\cdot 2H_2O$  (n reveals a novel structural arrangement of  $[Mn^{II}_{3}O]^{7+}$  and displays an obvious frequency-dependent out-of-phase component in AC magnetic susceptibilities and magnetic hysteresis loop measurements, which are indicative of SMM behavior.

 $Mn(ClO_4)_2 \cdot 6H_2O$ , NEt<sub>3</sub>, NaN<sub>3</sub>, and H<sub>2</sub>salox in a 1:1:1:1 ratio leads to a dark green methanol solution. Ether was slowly diffused into the solution to yield dark green-brown needle crystals of  $[Mn_{12}O_4(salox)_{12}]$  $(N_3)_4(MeOH)_6(H_2O)_2]$  · 5MeOH · 2H<sub>2</sub>O (1 · 5MeOH · 2H<sub>2</sub>O) in a 23% yield [25]. The infrared (IR) spectrum of complex **1** shows a simple  $v_{as}$ absorption at 2057  $\text{cm}^{-1}$  and a moderately intense band at 1597  $\text{cm}^{-1}$ in the carbonyl stretching region, with enhanced intensity and complexity in the azido and C = N stretching regions. X-ray crystallographic results of the product readily explain the origin of the spectroscopic modifications. Complex 1.5MeOH.2H<sub>2</sub>O crystallizes in triclinic space groups  $P_{\overline{1}}$  with a = 12.4812(4) Å, b = 14.2985(4) Å, c = 19.0003(6) Å,  $\alpha = 85.4138(10)^\circ$ ,  $\beta = 99.608(2)^\circ$ ,  $\gamma = 66.4047(11)^\circ$ ,  $V = 2968.21(16) \text{ Å}^3$ , Z = 1.  $R(R_w) = 0.0603$  (0.1297), and the molecule lye in the central position and the asymmetric unit contains the whole cluster, five MeOH and two water molecules [26]. The ORTEP representation of 1.5MeOH.2H<sub>2</sub>O is shown in Fig. 1. Complex 1.5MeOH $\cdot$ 2H<sub>2</sub>O can be described as two [Mn<sup>III</sup><sub>6</sub>O<sub>2</sub>(salox)<sub>6</sub>(N<sub>3</sub>)<sub>2</sub> (MeOH)<sub>3</sub>(H<sub>2</sub>O)] subunits bridged by two oximato O atoms (O6,

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Fig. 1. (a) Crystal structure drawing of complex 1.5MeOH $\cdot 2$ H<sub>2</sub>O and (b) arrangement of the four [Mn<sub>3</sub>O] triangles. Hydrogen atoms have been omitted for clarity.

O6'). Each Mn<sub>6</sub> subunit consists of six Mn ions arranged as two  $[Mn^{III}]_3$  $(\mu_3 - 0^{2-})$ <sup>7+</sup> triangles bridged by two end-on azides and one oximato with an  $\sim 41^{\circ}$  canting angle between the two Mn<sub>3</sub> planes (Fig. 1S). All manganese atoms are assigned as trivalent based on (i) the clear presence of the JT axial elongation, as expected for a high-spin d<sup>4</sup> ion in near-octahedral geometry and (ii) bond valence sum (BVS) calculations (Table 2S) [27-29]. The Mn ions can be divided into two groups: four Mn<sup>III</sup> ions (Mn1, Mn3, Mn4 and Mn5) with six-coordination and two Mn<sup>III</sup> ions (Mn2 and Mn6) with square pyramidal five-coordination. Peripheral ligations around each Mn<sub>6</sub> subunit are provided by four  $\eta^1:\eta^1:\eta^1:\mu_2$  salox<sup>2-</sup> ions, two  $\eta^1:\eta^1:\eta^2:\mu_3$  salox<sup>2-</sup> ions, and two  $\eta^2:\mu_2$  N<sub>3</sub> ions. Three MeOH and one H<sub>2</sub>O molecules terminally complete the sixth coordination position around Mn3, Mn4 and Mn5. In addition, there is one hydrogen bonding interaction which involves the oximato O atom and the O atom of a terminal MeOH from a neighboring Mn<sub>12</sub> molecule to form a 1-D supramolecular chain [O11···O18' 2.790 Å] (Fig. 2S).

Solid-state, variable-temperature magnetic susceptibility measurements were performed on fresh microcrystalline samples of complex **1**·5MeOH·2H<sub>2</sub>O suspended in eicosane to prevent torquing; the results are shown in Fig. 2. The direct current (DC) magnetic susceptibility data were collected in the 1.8–300 K range in a 1.0 kG magnetic field. The  $\chi_M T$  value of complex **1**·5MeOH·2H<sub>2</sub>O decreases steadily from 31.9 emu mol<sup>-1</sup> K at 300 K with decreasing temperature, reaching 11.6 emu mol<sup>-1</sup> K at 10 K, and then decreases sharply to 2.5 emu mol<sup>-1</sup> K at 1.8 K. The  $\chi_M T$  value at 300 K for complex **1**·5MeOH·2H<sub>2</sub>O is lower than 36.0 emu mol<sup>-1</sup> K, the value expected for a Mn<sup>III</sup><sub>12</sub> complex with noninteracting metal centers with g = 2.0, which may be due to the antiferromagnetic interactions between Mn ions. The decrease of the  $\chi_M T$  value with decreasing temperature also indicates antiferromagnetic coupling within complex



**Fig. 2.** Plots of  $\chi_M T$  vs. *T* for **1** · 5MeOH · 2H<sub>2</sub>O in an applied field of 1 kG from 2.0 to 300 K.

1.5MeOH.2H<sub>2</sub>O. The rapid decrease observed in the  $\chi_{\rm M}T$  value at 2.0– 10 K range is most likely due to zero-field splitting effects and perhaps weak intermolecular interactions mediated by the hydrogen bonding in the crystal structure. A theoretical treatment of the susceptibility data using the Kambe approach is not feasible due to the topological complexity and low symmetry of 1.5MeOH.2H<sub>2</sub>O [30]. Instead, our efforts focused on the determination of the S value of 1.5MeOH.2H2O from variable-temperature and variable-field magnetization (M) measurements in the 2.0-4.0 K and 1-5 T ranges, respectively (Fig. 3). The fits of the data using the program ANISOFIT [31], assuming that only the ground state is populated, were poor and suggest complications from low-lying excited states even at these relatively low temperatures. The measurements of the AC magnetic susceptibility were performed on fresh microcrystalline samples of complex  $1.5MeOH.2H_2O$  in the temperature range 1.8-20 K in a zero DC field and a 3.5 G AC field oscillating at frequencies in the 10-10,000 Hz range. The results for a representative sample of complex 1.5MeOH.2H<sub>2</sub>O are shown in Fig. 4. If the magnetization vector can relax fast enough to keep up with the oscillating field, there is no imaginary (out-of-phase) susceptibility signal ( $\chi_M$ ), and the real (in-phase) susceptibility ( $\chi_M$ ) is equal to the DC susceptibility. In contrast, if the barrier to magnetization relaxation is significant compared to the thermal energy (*kT*), there is a nonzero  $\chi_M$ " signal and the in-phase signal decreases. In addition, the  $\chi_M$  signals will



**Fig. 3.** Plot of reduced magnetization ( $M/N\beta$ ) vs. H/T for complex 1 · 5MeOH · 2H<sub>2</sub>O at the indicated field.

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