



## Original Research Paper

## Numerical simulation of multilayer deposition in an obstructed channel flow

Gregory Lecrivain<sup>a,\*</sup>, Drapeau-Martin Sevan<sup>a</sup>, Barth Thomas<sup>a</sup>, Uwe Hampel<sup>a,b</sup><sup>a</sup> Helmholtz-Zentrum Dresden-Rossendorf, Institute of Fluid Dynamics, Bautzner Landstraße 400, 01328 Dresden, Germany<sup>b</sup> AREVA Endowed Chair of Imaging Techniques in Energy and Process Engineering, Dresden University of Technology, 01062 Dresden, Germany

## ARTICLE INFO

## Article history:

Received 19 January 2013

Received in revised form 19 March 2013

Accepted 13 May 2013

Available online 2 June 2013

## Keywords:

Particle dispersion

Multilayer deposition

Dry granular mechanics

Obstructed channel flow

## ABSTRACT

Simulation of multilayer deposition of dry aerosol particles in turbulent flows has gained a growing interest in various industrial and research applications. The multilayer deposition of carbonaceous aerosol particles in a turbulent channel flow obstructed by a succession of square ribs is here numerically investigated. The multilayer particle bed growth on the various wall surfaces affects the air flow, which in turn affects the overall deposition rate. An iterative numerical procedure is therefore suggested to simulate the evolution of the graphite layer. The iterative process used to reproduce the layer build-up is decomposed as follows: Reynolds-Averaged Navier Stokes is employed to generate the flow field. The turbulent dispersion of the particles is reproduced through the use of a continuous random walk model. After statistically sufficient deposition of particulate matter, the layer build-up is computed using mechanics of dry granular material. The layer build-up model shows good agreement with data obtained from experimental tests carried out on-site.

© 2013 The Society of Powder Technology Japan. Published by Elsevier B.V. and The Society of Powder Technology Japan. All rights reserved.

## 1. Introduction

## 1.1. Motivation

The deposition of micrometre-sized aerosol particles spans across a wide range of applications such as to name a few examples, exposure to hazardous particulate matter, sedimentation of atmospheric particles and surface obscuration. The deposition of aerosol particles has also gained increasing interest in the field of nuclear safety research [1]. During operation of a graphite moderated high temperature pebble-bed reactor, the motion of the spherical fuel elements within the core initiates the production of carbonaceous dust. The graphite is then conveyed by the coolant gas and deposit in the primary circuit over the course of years. The dust layer build-up as a result of graphite deposition in the cooling system eventually becomes a major source of ionising radiation [2,3].

## 1.2. Particle deposition in a turbulent flow

Deposition denotes here the process of attaching suspended aerosol particles from a gas in turbulent motion to a surface. Friess and Yadigaroglu [4] distinguished two types of deposition: monolayer deposition and multilayer deposition. Monolayer deposition

indicates that particles adhere directly to a wall surface and are not in contact with one another. Multilayer deposition indicates however that particles densely accumulate on the wall surface to form a multilayer particle bed henceforth referred to as dust layer.

Whether monolayer or multilayer, turbulent diffusion is a chief agent responsible for deposition of micrometre-sized aerosol particles. The results of deposition experiment are frequently plotted as curves of particle deposition velocity  $u_d$  against particle response time  $\tau_p$ . The deposition velocity is defined by the surface particle mass flux  $J$  divided by the airborne particle mass concentration  $C_0$ . As illustrated in Fig. 1, two deposition regimes typically exist [5]. For fine particulate matter (diameter typically lower than  $10\ \mu\text{m}$ ), the deposition velocity increases with particle size. Fine particles at the low end of the regime are stopped by strong drag forces in the viscous sublayer whereas fine particles at the very end of the regime possess sufficient energy to overcome the drag forces and eventually collide with the wall surface [6]. For large particulate matter (diameter typically greater than  $10\ \mu\text{m}$ ), the regime becomes ballistic [7], i.e. large particles have greater inertia so that the trajectories are less affected by near wall turbulence. In this regime the deposition velocity reaches a maximum and remains fairly constant.

The diffusion of fine and large particulate matter is largely governed by the unsteady turbulent gas fluctuations. There are currently three well-established approaches to numerically reproduce the transport and deposition of particles. These are Direct Numerical Simulation (DNS), Large Eddy Simulation (LES)

\* Corresponding author. Tel.: +49 351 260 3768; fax: +49 351 260 1 3768.

E-mail address: [g.lecrivain@hzdr.de](mailto:g.lecrivain@hzdr.de) (G. Lecrivain).

**Nomenclature***Greek*

$\alpha_d$	dynamic angle of repose
$\alpha_p$	drag correction factor
$\alpha_s$	static angle of repose
$\delta t_c$	critical time step
$\delta x$	horizontal grid spacing
$\delta y_c$	critical height
$\epsilon$	dissipation rate
$\eta_d$	deposited fraction
$\nu$	fluid kinematic viscosity
$\phi$	medium porosity
$\tau_p$	particle response time
$\xi_i$	zero mean, unit variance distributed random number

*Latin*

$C_0$	particle mass concentration
$d_p$	particle diameter
$dm/dt$	particle mass flow
$e$	rib height
$e_k$	sharing coefficient
$\mathbf{g}$	gravitational acceleration
$H$	channel height
$J$	particle mass flux
$k$	turbulent kinetic energy
$L$	channel length
$L_{Vi}$	associated length of the recirculation zone $V_i$ in the streamwise direction

$p$	rib pitch
$Re^*$	friction-velocity-based Reynolds number
$Re_b$	bulk Reynolds number
$Re_p$	particle Reynolds number
$S$	particle to fluid density ratio
$t_c$	critical time
$T_L$	Lagrangian integral time scale
$\mathbf{u}$	instantaneous fluid velocity ( $\mathbf{u} = \bar{\mathbf{u}} + \mathbf{u}'$ )
$\mathbf{u}'$	fluctuating fluid velocity
$u^*$	friction velocity
$\mathbf{u}^c$	drift correction velocity
$U_b$	bulk fluid velocity
$u_b$	build-up velocity
$u_d$	deposition velocity
$\bar{u}_d^f$	mean deposition velocity on the floor
$\mathbf{u}^p$	particle velocity
$V1$	recirculation zone at downstream corner of the rib
$V2$	largest recirculation zone downstream of the rib
$V3$	smaller recirculation zone at upstream corner of the rib
$y$	distance to the nearest wall

*Miscellaneous*

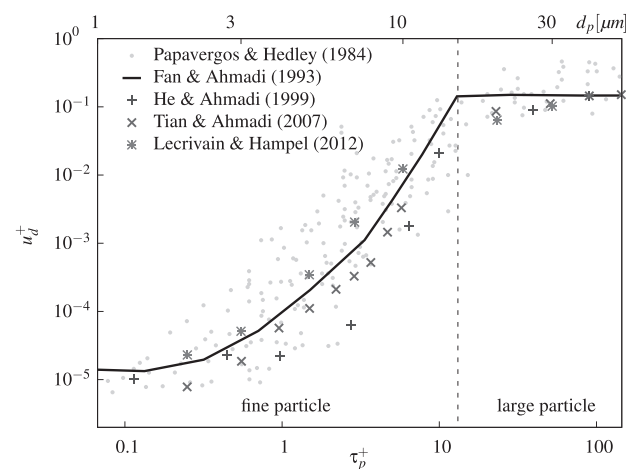
$X^+$	superscript denoting the normalisation of arbitrary quantity $X$ with wall scaling ( $u^*$ for velocity variables, $\nu/u^*$ for length variables, $\nu/u^{*2}$ for time variables and $u^{*3}/\nu$ for acceleration variables)
$\bar{X}$	overbar denoting time averaging of arbitrary quantity $X$

and Reynolds-Averaged Navier Stokes (RANS) simulations. While DNS can precisely resolve all turbulent features down to the smallest scales, its heavy computational cost makes inappropriate for turbulent particle-laden flows in complex geometries. Its use has therefore been largely limited to simple geometries such as channel flows [8]. In the LES approach, the large eddies are directly simulated whereas eddies smaller than the grid scale are modelled. Studies on particle deposition using LES have been performed in more complex turbulent flows such as flows in respiratory tract [9]. RANS simulations have gained a growing interest owing to its relatively small computational cost since only the mean flow field is computed. Information on the turbulence is only available in the statistical terms and therefore a dispersion model must be implemented beforehand to reproduce the turbulent diffusion of particles [10]. The deposition of particles using a turbulent dispersion model combined with a RANS simulation has proven successfully under various conditions. The continuous random walk model of Tian and Ahmadi [11] accurately predicted the deposition of nano- and micro-particles in a turbulent channel flow. The random walk model developed by Berrouk [12] has also shown some reasonable accuracy in predicting deposition rates in a curved piped. Li et al. [13] investigated the deposition of aerosol particles around a single obstruction placed in a two-dimensional channel. His findings showed that the particle deposition rate decreases as the shape of the obstacle becomes more streamlined.

### 1.3. Particle layer build-up in a turbulent flow

A few attempts have been made in the literature to predict layer build-up in turbulent flows. Sarimeseli and Kelbaliyev [14] derived empirical equations to reproduce the multilayer deposition of 10  $\mu\text{m}$  size particles in a horizontal two-dimensional cylindrical pipe. As opposed to vertical flow, the distribution of the deposition

velocity across the pipe circumference was included to account for gravity effects. Stempniewicz et al. [15] later estimated the layer build-up of carbonaceous dust in various parts of a pebble bed modular reactor. Through the use of a fairly simple one-dimensional code, the team simulated a build-up that reached, after years of sedimentation, a height of a few centimetres. In the study of Friess and Yadigaroglu [4], the two-dimensional clustering of micron-sized particles was investigated in greater detail. The deposition and re-entrainment of Lagrangian particles in the turbulent boundary layer over a flat plate was simulated. More recently Bozzi and Passoni [16] combined a two-dimensional LES with Lagrangian particles to reproduce the evolution of a sand heap in a two-dimensional turbulent flow.



**Fig. 1.** Effect of the particle response time  $\tau_p$  on the deposition velocity  $u_d$ .

Download English Version:

<https://daneshyari.com/en/article/144071>

Download Persian Version:

<https://daneshyari.com/article/144071>

[Daneshyari.com](https://daneshyari.com)