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Influence of hydrogen content and network connectivity on the coefficient of thermal expansion and thermal stability for a-SiC:H thin films

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ABSTRACT

Thin films of a-SiC:H are of interest for numerous applications in optoelectronic, nanoelectronic, and nanoelectromechanical devices due to a large optical band gap and excellent chemical inertness. In many cases, a low coefficient of thermal expansion (CTE) and high temperature stability are also desired. In this regard, we report an investigation of the influence of hydrogen content and average bond coordination ($\langle r \rangle$) on the thermal stability and CTE of a-SiC:H thin films deposited by plasma enhanced chemical vapor deposition. The a-SiC:H hydrogen content determined by combined nuclear reaction analysis and Rutherford backscattering measurements (NRA–RBS) was skewed from 25 to 60% resulting in films with $\langle r \rangle$ ranging from 2.0 up to 3.2. We show that increased hydrogen content is accompanied by changes in the thermal stability and CTE for the a-SiC:H films. The most dramatic changes were observed to occur for films with hydrogen content >45% where $\langle r \rangle$ is near the critical coordination number ($\langle r \rangle_c$) that delineates over and under constrained materials. For a-SiC:H films above $\langle r \rangle_c$, relatively little change was observed in CTE with $\langle r \rangle$. However for films near or below $\langle r \rangle_c$ CTE was observed to increase with decreasing $\langle r \rangle$ and significant hydrogen loss was observed after the annealing measurements.

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1. Introduction

Hydrogenated amorphous silicon carbide (a-SiC:H) is a material with excellent material properties that makes it of interest for use in numerous optoelectronic [1,2], biomedical [3,4], nanoelectronic [5,6], tribological [7,8], and nanoelectromechanical [9,10] devices. Specific applications include use as an amorphous semiconductor [11–13], a surface passivation layer [14], a hermetic moisture or metal diffusion barrier [15–18], a corrosion and wear resistant protective coating [19–22], an insulating intermetal dielectric [23,24], an etch stop [25–27], and an x-ray hard mask [28,29]. In addition, a-SiC:H offers the added ability to tune the physical properties via varying the Si/C ratio or the total hydrogen content [30–32]. The former allows the realization of materials with properties spanning those of a-Si and a-C [33–37] while the latter allows the preparation of materials with a still wider range of properties [32,38–41].

In many of the above applications, a mix of material properties sometimes at disparate extremes is highly desired. One example would be the use of a-SiC:H as a low dielectric constant (i.e. low-*k*) material in nanoelectronic metal interconnect structures [42] where a low dielectric constant is desired in combination with a high Young's modulus and low coefficient of thermal expansion (CTE) [43]. The low dielectric constant is desired to reduce resistance–capacitance delays in the metal interconnect whereas a low CTE and high elastic modulus is needed to minimize thermally generated strains and to provide sufficient rigidity to survive downstream processing and packaging processes [42–45]. These requirements are at odds with one another as low dielectric constants are typically achieved via incorporation of significant amounts of hydrogen that disrupts the Si–C network bonding [46,47], whereas a high modulus and low CTE are typically achieved at low hydrogen content and high Si–C network connectivity [40,48].

In light of the above conflicting requirements for a-SiC:H in low-*k* dielectric applications, we have elected to perform a detailed investigation of the thermal–mechanical properties of plasma enhancedchemically vapor deposited (PECVD) a-SiC:H films as a function of hydrogen content and network connectivity. In previous studies [40,46], we have demonstrated that increased hydrogen content has a significant influence on the a-SiC:H network connectivity and properties such as Young's modulus, hardness, and thermal conductivity. In this study, we have specifically chosen to investigate the thermal stability and both the in-plane and out-of-plane CTE for a-SiC:H materials with varying

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hydrogen content and average coordination ($\langle r \rangle$). In contrast to the previous studies, we find that hydrogen content has a negligible influence on CTE and thermal stability for films where $\langle r \rangle$ is greater than the critical coordination number ($\langle r \rangle_c$) predicted by bond constraint/percolation theory. For $\langle r \rangle < \langle r \rangle_c$, we find that CTE increases significantly and the a-SiC:H films become increasingly thermally unstable as exhibited by significant changes in the film thickness and the loss of hydrogen after extended annealing at 400 °C (note this is substantially less than the melting point of single crystalline SiC (T_{melt} \cong 2800 °C [49])). For $\langle r \rangle > \langle r \rangle_c$, the combined in-plane and out-of-plane CTE measurements were found to be invariant with $\langle r \rangle$ and yield similar results confirming the expected isotropic properties for the as deposited a-SiC:H films.

2. Experimental

2.1. Film deposition

All the a-SiC:H films analyzed in this study were deposited using standard commercially available parallel plate capacitance PECVD tools and various silane- and methylsilane-like sources diluted in gases such as H₂ or He. All the films were deposited on 300 mm diameter, double-sided polished (100) silicon substrates. The details of the a-SiC:H deposition have been provided previously [46,47,50]. Briefly, most films were deposited at temperatures on the order of 400 °C, however, one highly porous film (a-SiC:H#1) was deposited at a lower temperature of 250 °C. For a-SiC:H#1, an organic pore building precursor (i.e. "porogen") was intentionally added to the deposition process to facilitate the creation of significant levels of interconnected nanopores in the a-SiC:H film and to produce an ultra low-k a-SiC:H film [38]. In this case, the porogen was removed post-deposition using an industry standard electron beam cure at ~400 °C [51]. The hydrogen content for the a-SiC:H films was further controlled by adjusting the deposition pressure, high (13.56 MHz) and low (200-400 kHz) frequency applied power, and reactant/diluent ratio.

Film thickness and refractive index were determined using a J. A. Woollam variable angle spectroscopic ellipsometer (VASE) [52]. The a-SiC:H film stress immediately post-deposition was calculated using Stoney's formula and the optically determined film thickness and change in Si wafer curvature measured using a laser deflection method [53,54]. The percent porosity for these materials was determined by ellipsometric porosimetry (EP) measurements using toluene as the solvent [55–57]. The Young's modulus and hardness of the a-SiC:H films was determined by nano-indentation as previously described [58]. Table 1 summarizes some of the key material properties for the a-SiC:H films investigated in this study and that have been previously reported [46,47,50].

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ummary of select material properties for PECVD a-SiC:H films investigated in this study	umma
D = Not Detected.	√D =

Film	RBS C/Si	% Porosity	Stress (MPa)
a-SiC:H#1	5.0	<12	50
a-SiC:H#2	3.0	<2	10
a-SiC:H#3	1.1	ND	20
a-SiC:H#4	1.1	ND	60
a-SiC:H#5	1.5	ND	0
a-SiC:H#6	1.2	ND	-52
a-SiC:H#7	1.1	ND	-94
a-SiC:H#8	1.1	ND	-182
a-SiC:H#9	1.1	ND	-140
a-SiC:H#10	1.1	ND	-425
a-SiC:H#11	1.05	ND	-400
a-SiC:H#12	1.0	ND	-1100

2.2. In-plane CTE

The in-plane CTE for the a-SiC:H films was measured using a Frontier Semiconductor TC900 laser stress measurement system. Specifically, this tool was utilized to monitor changes in the a-SiC:H/Si wafer curvature as a function temperature using a laser deflection system previously described [59]. The wafer curvature is directly proportional to the intrinsic stress in the a-SiC:H film according to Stoney's formula [60]:

$$\sigma_{\rm f} = E_{\rm s} h_{\rm s}^2 \Delta K / 6 h_{\rm f} \tag{1}$$

where σ_f is the biaxial film stress, E_s is the biaxial modulus of the Si substrate, h_s is the Si substrate thickness, h_f is the film thickness and ΔK is the change in substrate curvature after deposition of the a-SiC:H film. The value of σ_f post-deposition is the sum of thermally generated stresses (σ_t) and intrinsic stresses generated during the deposition process (σ_i) [58]. The thermally generated stress σ_t is a result of the mismatch in thermal contraction between the a-SiC:H film and the Si substrate on cooling from the deposition temperature. The intrinsic stress σ_i is the result of strained bonds formed in the film due to either unoptimized bond orientation or strain imparted to the film by ion bombardment during the plasma deposition process. On heating, the radius of curvature of the a-SiC:H/Si wafer will change with temperature (T) as a result of changes in the a-SiC:H film stress due to changes in the σ_t component. In this case, the CTE of the film (α_f) may be determined using the following relation [61]:

$$\partial \sigma_{\rm f} / \partial T = \left[E_{\rm f} / (1 - \nu_{\rm f}) \right] \left(\alpha_{\rm s} - \alpha_{\rm f} \right) \tag{2}$$

where E_f and v_f are respectively Young's modulus and Poisson's ratio for the a-SiC:H film, and α_s is the CTE for the Si substrate. For the experiments described here, the change in radius of curvature while heating from 23 to 400 °C in a $<10^{-5}$ Torr vacuum was measured using 500 nm thick a-SiC:H films deposited on Si. The maximum temperature of 400 °C was selected as this was the deposition/curing temperature for the films and is generally the maximum temperature allowed in low-k/Cu interconnect fabrication [54]. The heating rate was 5 °C/min and the cooling rate was approximately 10 °C/min. The samples remained at the target temperature of 400 °C for 5 min before cooling. The radius of curvature was monitored on heating and cooling for two cycles. Linear regression analysis was utilized to determine $\delta\sigma_f/\delta T$ over the temperature range of approximately 50–300 °C. To calculate the CTE of the a-SiC:H films, the CTE of the Si substrate was taken to be 2.6 °C/ppm and the a-SiC:H Young's modulus values previously determined by nanoindentation [58] were utilized. The Poisson's ratios for the a-SiC:H films were assumed based on previously reported measurements on identical or similar samples [21,62,63]. Specifically, the Poisson's ratio for a-SiC:H#1 and 2 was taken to be 0.33 \pm 0.03 based on prior Brillouin light scattering measurements of identical samples [62]. For a-SiC:H#3–12, the Poisson's ratio was assumed to be 0.25 \pm 0.03 based on prior measurements by Khakhani [21] and Cros [63] for a-SiC:H films of similar stoichiometry and hydrogen content.

2.3. Out-of-plane CTE

The out-of-plane CTE for identical 100 nm a-SiC:H films deposited on Si was determined by x-ray reflectivity (XRR) thickness measurements performed between room temperature and 400 °C. The details of these measurements have been described previously [64,65]. Briefly, the XRR measurements were carried out on a Bruker D8 Discover diffractometer (Cu K α radiation, $\lambda = 0.15418$ nm) equipped with a DHS 900 domed hot stage attachment. The temperature of the hot stage was controlled via a Paar Physica TCU 150 temperature control unit with N₂ gas continuously flowing through the dome. The samples were heated from the back through thermal contact with the hot Download English Version:

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