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Intermediate structural phases in rare-earth substituted BiFeO₃

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1. Introduction

Despite considerable efforts undertaken in a rapidly developing area of multiferroic research [1,2], the effect of temperature [3,4], high pressure [5,6], and chemical substitution [7,8] on crystal structure and physical properties of bismuth ferrite is still a matter of intensive research. At ambient conditions, the compound has a rhombohedrally distorted structure with a space group R3c [9]. Asite substitution with lanthanides (Ln) is known to induce a number of structural transitions in BiFeO₃, but symmetry of the intermediate structural forms existed between polar R3c and nonpolar Pnma phases characteristic of the end members of $(BiFeO_3)_{1-x}$ - $(LnFeO_3)_x$ series [9,10] is not unambiguously determined yet. In one of the pioneering work on this topic, the $R3c \rightarrow P1 \rightarrow C222 \rightarrow C222_1 \rightarrow Pn2_1a \rightarrow Pnma$ sequence of the composition-driven transitions (at $x \approx 0.06$, 0.24, 0.40, 0.55, and 0.7, respectively) was proposed for $Bi_{1-x}La_xFeO_3$ single crystals [11]. With decreasing Ln³⁺ ionic radius, the concentrational boundaries of the structural phases were claimed to shift towards smaller x [11]. More recent investigation of $Bi_{1-x}Gd_xFeO_3$ polycrystalline samples confirmed existence of the intermediate polar structure $Pn2_1a$ (0.1 < x < 0.2) [8,12]. Study of $Bi_{1-x}Ln_xFeO_3$ (Ln = Nd, Sm) compounds revealed a substitution-induced transition to antipolar PbZrO₃-like phase at $x \sim 0.15$ [13,14]. A coexistence of R3c, Imma, and *Pnma* phases was found for $Bi_{1-x}Ln_xFeO_3$ (Ln = La, Nd) samples

ABSTRACT

The room-temperature crystal structure and local ferroelectric properties of Bi_{0.85}RE_{0.15}FeO₃ (RE = Sm, Gd, Dy) polycrystalline samples have been investigated by X-ray diffraction (XRD) and piezoresponse force microscopy (PFM) techniques. XRD measurements show that the rare-earth substitution causes the transformation of *R3c* structure typical of BiFeO₃ to yield orthorhombic phases with $\sqrt{2a \times 2\sqrt{2a \times 2a}}$ and $\sqrt{2a \times 2a \times \sqrt{2a}}$ superlattices (*a* is the parameter of the cubic perovskite subcell). Results of PFM investigations imply that both the orthorhombic phases are polar. The models describing crystal structure of the phases are discussed.

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in a broad concentration range [15]. Such a variety of the structural models is probably explained by a strong tendency of the $Bi_{1-x}Ln_xFeO_3$ multiferroics to phase separation and formation of impurity phases [15], that hampers true crystal structure identification. To reveal the intermediate phases of $Bi_{1-x}Ln_xFeO_3$ which can be easily stabilized during a conventional solid-state reaction and to define the basic properties, which can help to determine the crystal symmetry of the phases, we performed synthesis of $Bi_{0.85}RE_{0.15}FeO_3$ (RE = Sm, Gd, Dy) compounds at moderate temperature. Taking into account that formation of the intermediate phases in $Bi_{1-x}Ln_xFeO_3$ seems to be the result of the cation diffusion between the easily forming phases with $BiFeO_3$ -type (*R3c*) and $LnFeO_3$ -type (*Pnma*) structures [15], the moderate-temperature synthesis is expected to give us a "composition sweep" of the possible structural forms.

2. Experimental

Polycrystalline samples with nominal chemical compositions $Bi_{0.85}RE_{0.15}FeO_3$ (RE = Sm, Gd, Dy) were prepared by the conventional solid-state reaction method using Bi_2O_3 , Sm_2O_3 , Gd_2O_3 , Dy_2O_3 , and Fe_2O_3 oxides. The synthesis was carried out in air at moderate temperature of 850 °C [12–15] for 6 h with a heating/ cooling rate of 5 °C/min. The crystal structure of the samples was studied by X-ray diffraction technique using an automated Philips PW 1050/35 diffractometer with Cu K_{\alpha} radiation. XRD patterns were collected over the angular range $20 \le 2\theta \le 100^\circ$ with step of 0.02° and exposition of 25 s/step. Local ferroelectric properties of the samples were investigated with piezoresponse force micros-

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copy using a commercial setup Multimode NanoScope IIIA (Veeco) equipped with a lock-in amplifier (SR-830A, Stanford Research) and a function generator (FG-120, Yokagawa). A commercial tip-cantilever system ArrowTM Silicon SPM Sensor (NanoWorld) was used. Domain visualization was performed under an applied *ac* voltage with the amplitude $V_{ac} = 2.5$ V and frequency f = 50 kHz. Local piezoelectric hysteresis loops were measured inside individual grains by applying the consecutive voltage pulses and measuring the piezoelectric response as a function of the voltage.

3. Results and discussion

In spite of the rather complex character of the obtained XRD spectra, suggesting a multiphase structural state of the samples, a careful analysis made using the FullProf software package allowed us to identify all the diffraction peaks and to separate out contributions from different structural forms (Fig. 1). For Bi_{0.85}Sm_{0.15}FeO₃ sample, the main contribution is related to an orthorhombic phase O_l with $\sqrt{2a \times 2}\sqrt{2a \times 2a}$ superlattice (a, \sim 4 Å, is the parameter of the cubic perovskite subcell). Crystal structure of the phase was recently studied by Karimi et al. [13,14]. Taking into account a close analogy between the X-ray and electron diffraction data obtained for the $Bi_{1-x}Ln_xFeO_3$ (Ln = Nd, Sm; $x \sim 0.15$) and well-known PbZrO₃ [16], the authors identified this phase as having antiferroelectric structure [13,14]. It was also shown that the simplest unit cell that fitted the electron diffraction and was consistent with the XRD data was similar to that of PbZrO₃ except that the *c* parameter was quadrupled rather than doubled with respect to the fundamental perovskite lattice (i.e. the unit cell should have a $\sqrt{2a} \times 2\sqrt{2a} \times 4a$ metrics) [13,14]. We, however, did not find any additional XRD peaks suggesting the quadrupling the unit cell along the c axis. The same superstructure $\sqrt{2a} \times 2\sqrt{2a} \times 2a$ was found upon the synchrotron X-ray powder diffraction characterization of the structural phase stabilizing in the vicinity of the ferroelectric R3c phase of BiFeO₃ during application of external pressure [6]. The similarity of the structures is unlikely to be accidental: because of the difference in ionic radii between the host and the substituting elements, a chemical substitution could act on interatomic distances in the similar way as an external pressure, thus leading to appearance of the same structures [17]. As expected, two other dominant contributions to the resulting XRD pattern of the Sm-substituted sample were attributed to an orthorhombic phase O_{II} with $\sqrt{2a \times 2a \times \sqrt{2a}}$ superlattice characteristic of LnFeO₃ [10] and a rhombohedral



Fig. 1. Experimental XRD pattern for $Bi_{0.85}Sm_{0.15}FeO_3$ sample. Inset shows parts of the spectra obtained for $Bi_{0.85}RE_{0.15}FeO_3$ samples together with the Bragg reflections allowed by the space groups R3c, Pbam/Pba2, and $Pnma/Pn2_1a$, from top to bottom (see Section 3). *Position of the main peak characteristic of $Bi_{25}FeO_{39}$.

phase *R* typical of BiFeO₃ [9]. Traces of Bi₂₅FeO₃₉ impurity, which often accompanies the formation of the perovskite-like phase of BiFeO₃ [6,12–15], were also detected. The same main structural forms O_I , O_{II} , *R* were found for Gd and Dy-substituted compounds, the volume fraction of the phase O_I being reduced with decreasing ionic radius of substituting element.

To shed light on crystal structure of the discovered phases, we undertook investigation of local ferroelectric properties of the samples with piezoresponse force microscopy. Indeed, if the O_1 phase is really characterized by antiferroelectric PbZrO₃-like structure, it will be immediately confirmed by observation of the local areas showing no piezoresponse contrast, but demonstrating double piezoelectric hysteresis loops [18]. If the O_{II} phase has Pnma symmetry, neither PFM contrast, no local piezoelectric hysteresis loops will be observed [8]. PFM experiment, however, did not reveal a presence of any extended regions with zero piezoresponse. Typical examples of our PFM investigation are presented in Fig. 2. A piezoelectric contrast corresponding to the antiparallel ferroelectric domains was found in all locations tested (Fig. 2(a-c)). A unipolar piezoelectric response after scanning with a dc voltage applied to the tip was also observed. The shape of the PFM hysteresis loops taken on the surface of the samples is typical of ferroelectric materials (Fig. 2(d)). These facts imply that the phases found in the Bi_{0.85}RE_{0.15}FeO₃ composites are ferroelectric and their crystal structures must have a polar symmetry. The data are consistent with the results of investigations of piezoelectric, pyroelectric, and ferroelectric properties of Bi_{0.85}Nd_{0.15}FeO₃



Fig. 2. (a–c) Piezoresponse force microscopy images demonstrating existence of ferroelectric domains in $Bi_{0.85}RE_{0.15}FeO_3$ samples (RE = Sm, Gd, Dy, respectively). (d) Local piezoresponse hysteresis loop obtained for $Bi_{0.85}Sm_{0.15}FeO_3$ sample. (e and f) PFM images for $Bi_{0.8}RE_{0.2}FeO_3$ samples (RE = Sm and Dy, respectively).

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