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Periodic transmission peaks in non-periodic disordered one-dimensional photonic structures

Ilka Kriegel^{a,b}, Francesco Scotognella^{a,b,c,*}

^a Dipartimento di Fisica, Politecnico di Milano, Piazza Leonardo da Vinci 32, 20133 Milano, Italy

^b Center for Nano Science and Technology@PoliMi, Istituto Italiano di Tecnologia, via Giovanni Pascoli, 70/3, 20133 Milano, Italy

^c Istituto di Fotonica e Nanotecnologie CNR, Piazza Leonardo da Vinci 32, 20133 Milano, Italy

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ABSTRACT

A better understanding of the optical properties of a device structure characterized by a random arrangement of materials with different dielectric properties at a length scale comparable to the wavelength of light is crucial for the realization of new optical and optoelectronic devices. Here we have studied the light transmission of disordered photonic structures made with two and three different materials, characterized by the same optical thickness. In their transmission spectra a formation of peaks, with a transmission of up to 75%, is evident. The spectral position of such peaks is very regular, which is a result of the constraint that all layers have the same optical thickness. This gives rise to a manifold of applications such as new types of bandpass filters and resonators for distributed feedback lasers.

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1. Introduction

The periodic modulation of the dielectric constant, at a length scale comparable to the wavelength of light, is the key property of photonic crystals [1–4]. Such materials are employed for the realization of diverse optical devices, as for example distributed feedback laser [5–7] and optical switches [8,9]. The addition of defects in the periodic alternation, or the realization of completely random sequences, results in disordered photonic structures [10–13]. In the case of one-dimensional disordered photonic structures, very interesting physical phenomena have been theoretically predicted or experimentally observed. In a pioneering work by Faist et al. [14], it was calculated the reflectivity of a GaAs/ $(GaAs)_n(AlAs)_m$ multilayer in which a layer thickness randomness was introduced, and reflectivity was significant for a wide energy region (wider with respect to the bandgap of the periodic structure). Later, an extension of the band gap with the introduction of disorder in multilayers has been theoretically and experimentally demonstrated [15,16]. Generally, the introduction of disorder results in new optical properties that can be very useful for new types of optical device components.

In this work, we have found the formation of a transmission peak at a precise energy region for disordered photonic structures

E-mail address: francesco.scotognella@polimi.it (F. Scotognella).

in which all the layers have the same optical thickness. The photonic structures are comprised of two and three different materials, and for both the formation of transmission peaks with a transmission of about 75% is evident. The periodic repetition of those peaks is a result of the precondition that all layers have the same optical thickness.

2. Methods

We considered periodic and disordered photonic structures made with two (A and B) and three (A, B, and C) different materials. The periodic structure composed of two materials repeated the lattice cell AB, whereas the disordered structure randomly arranged the layers A and B along the multilayer. Similarly, the periodic structure composed of three materials repeated the lattice cell ABC, while the disordered structure randomly arranged the A, B, and C layers along the multilayer. Very important is the condition applied to the thickness of the layers, which requires them to have the same optical thickness of, in this case, $d_i = c/n_i$ nm, with i = A, B, C and n_i being the refractive index of the corresponding *i*th layer. In the calculations, we have employed various refractive indexes and different total number of layers and chose the value *c* to be 310 nm.

For the calculations of the transmission spectra of the periodic and disordered photonic structures, we have employed the transfer matrix method [17,18]. The spectra were calculated with a step of 1 meV.





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^{*} Corresponding author at: Dipartimento di Fisica, Politecnico di Milano, Piazza Leonardo da Vinci 32, 20133 Milano, Italy.

3. Results and discussion

In Fig. 1a we show the transmission spectra of the periodic photonic structure repeating the unit cell AB for 1000 times (i.e. 1000 unit cells). The refractive index of the layer A and B have been chosen to be n_A = 1.8 and n_B = 2. The band gap occurs at 1 eV, and its higher order modes are found at 3 and 5 eV. Due to the condition to keep the same optical thickness among the layers, the peak positions are found at $E_m = \frac{1240}{4c} \cdot (2m-1) \text{ eV}$ with m = 1, 2...,with an energy separation of $\Delta E = \frac{1240}{2c}$ eV. This finding has been reported in a previous work of ours [19] and demonstrates that the photonic band gap position is independent of the chosen refractive index and depends solely on c. As predicted by the given formulation, the photonic band gap is found at $E_1 = 1$ eV and its higher order modes at E_2 = 3 eV and E_3 = 5 eV for the chosen value of c = 310 nm, and thus, thickness of the layers being $(d_i = 310/n_i - 10)$ nm). As demonstrated already previously [19], the gaps at $2E_1$ and $4E_1$ are not present. In the ordered three layer photonic structure (with the number of layers being t = 3) repeating the unit cell ABC with $n_A = 1.8$, $n_B = 2$, and $n_C = 2.2$ for 1500 times, each gap E_m

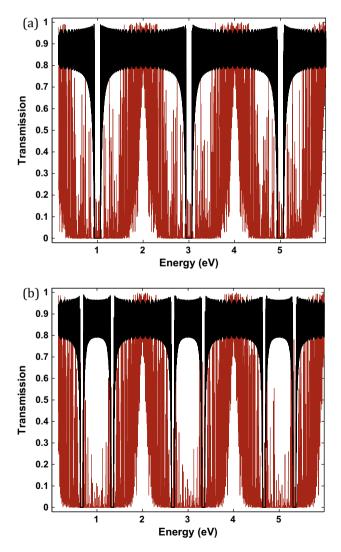


Fig. 1. (a) Periodic two-material photonic crystals (black curves) and random twomaterial photonic structures (orange curves). Both are made with 1000 layers; (b) periodic three-material photonic crystals (black curves) and random three-material photonic structures (orange curves). Both are made with 1500 layers. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

of the two layer photonic structure (band gap and higher order gaps at $E_1 = 1 \text{ eV}$, $E_2 = 3 \text{ eV}$, and $E_3 = 5 \text{ eV}$) is split into two gaps. This is another result from the condition of the layer thickness being $d_i = c/n_i$ reported by us in Ref. [19]. The energy position E_{SG} of the split gaps follows the rule of $E_{SG} = \frac{1240}{4c} \cdot (2m - 2) + \frac{1240}{4c} \cdot \left(\frac{2 \cdot (1.2 \dots t - 1)}{t}\right) \text{ eV}$, with *t* representing the number of layers in the unit cell of the ordered photonic structure with more than two layers $t = 2, 3, 4, \dots$ See black curves in Fig. 1a for two layers and Fig. 1b for three layers.

The calculation of the transmission properties of photonic crystal structures with a disordered arrangement of the layers, is based on a random distribution of the layers AB or ABC for a number of 1000 and 1500 layers, respectively. A similar layer thickness and refractive index as given above is used. The transmission spectra of the disordered structure for two and three materials are the orange curves shown in Fig. 1a and b, respectively. Interestingly, the completely disordered structures show very regular transmission peaks at 2 and 4 eV, which corresponds to the missing bandgaps at $2E_1$ and $4E_1$, i.e. following $E_n = \frac{1240}{4c} \cdot 2n \text{ eV}$ with $n = 1, 2, 3 \dots$ Such transparency windows show a high transmission above 75%. Notably,

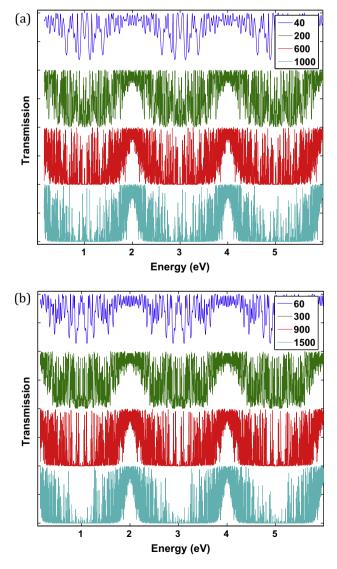


Fig. 2. (a) Random two-material photonic structures made by 40, 200, 600 and 1000 layers; (b) random three-material photonic structures made by 60, 300, 900 and 1500 layers.

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