# A novel nitridogallate fluoride $\mathrm{LiBa}_{5} \mathrm{GaN}_{3} \mathrm{~F}_{5}$ - Synthesis, crystal structure, and band gap determination 

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#### Abstract

$\mathrm{LiBa}_{5} \mathrm{GaN}_{3} \mathrm{~F}_{5}$ was obtained as red crystals by reaction of $\mathrm{Ba}, \mathrm{Ga}, \mathrm{NaN}_{3}$ and $\mathrm{EuF}_{3}$ in a $\mathrm{Na} / \mathrm{Li}$ flux at $760{ }^{\circ} \mathrm{C}$ in weld-shut tantalum crucibles. The crystal structure (Pnma (no. 62), $a=15.456(3), b=5.707(1)$, $c=12.259(3) \AA, Z=4)$ was solved on the basis of single-crystal X-ray diffraction data. In the solid there are trigonal planar $\left[\mathrm{GaN}_{3}\right]^{6-}$ ions and zigzag chains of vertex sharing $\mathrm{LiF}_{6}$ octahedrons surrounded by $\mathrm{Ba}^{2+}$ ions. Optical measurements and calculations of the electronic structure revealed a band gap of $\leq 1.9 \mathrm{eV}$. According to the calculations, the observed transition occurs from a nitrogen state into a hybrid $\mathrm{Ba} / \mathrm{N}$ state.


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## 1. Introduction

Ternary nitrides of $G a$ in combination with alkaline earth metals (e.g., Sr, Ba) are known since a couple of years [1-4], but only a small number of quaternary examples have been reported recently [5-7]. Most of these compounds have been synthesized in sodium melts which due to addition of alkaline earth metals exhibit an increased solubility of nitrogen [8]. Syntheses were performed in closed niobium crucibles, utilizing additional azides as nitrogen source. $\mathrm{Sr}_{3} \mathrm{GaN}_{3}, \mathrm{Sr}_{6} \mathrm{GaN}_{5}$ [2], $\mathrm{Sr}_{4} \mathrm{GaN}_{3} \mathrm{O}$ and $\mathrm{Sr}_{4} \mathrm{GaN}_{3}\left(\mathrm{CN}_{2}\right)$ [5] contain non condensed ("isolated") $\left[\mathrm{GaN}_{3}\right]^{6-}$ while $\mathrm{Ba}_{3} \mathrm{Ga}_{2} \mathrm{~N}_{4}$ and $\mathrm{Sr}_{3} \mathrm{Ga}_{2} \mathrm{~N}_{4}$ are made up of trans edge sharing $\mathrm{GaN}_{4}$ tetrahedrons building infinite chains [1,4]. Two- or threedimensional networks of vertex sharing $\mathrm{GaN}_{4}$ tetrahedrons have been found in $\mathrm{Ca}_{3} \mathrm{Ga}_{2} \mathrm{~N}_{4}, \mathrm{Sr}_{3} \mathrm{Ga}_{3} \mathrm{~N}_{5}$ [4] and LiSrGaN 2 [6], respectively. Thus, structural motifs similar to nitridosilicates (e.g., $\mathrm{BaSi}_{7} \mathrm{~N}_{10} ; \mathrm{Eu}_{2} \mathrm{SiN}_{3}$ ) $[9,10]$ occur in the nitridogallates mentioned above. Recently, we have reported about synthetic approaches to control the dimensionality of nitridosilicates employing lithium melts [11]. In comparison with nitridosilicates, the field of nitridogallates is more unexplored but utilization of the lithium

[^0]flux technique may lead to a larger structural variety analogously to our experience with nitridosilicates.

## 2. Experimental

The synthesis of $\mathrm{LiBa}_{5} \mathrm{GaN}_{3} \mathrm{~F}_{5}$ was carried out in Ta crucibles ( 30 mm length, 9.5 mm diameter, 0.5 mm wall thickness). Under argon atmosphere (glove box Unilab, MBraun), 0.35 mmol ( 22.8 mg ) $\mathrm{NaN}_{3}$ (Acros, $99 \%$ ), $0.138 \mathrm{mmol}(9.6 \mathrm{mg}$ ) Ga (AluSuisse, $99.999 \%$ ), $0.549 \mathrm{mmol}(75.4 \mathrm{mg})$ Ba (Sigma Aldrich, $99.99 \%$ ) and $0.027 \mathrm{mmol}(5.7 \mathrm{mg}) \mathrm{EuF}_{3}$ (Sigma Aldrich, $99.99 \%$ ) were mixed and filled into the Ta crucible. For the flux $2.174 \mathrm{mmol}(50.0 \mathrm{mg}) \mathrm{Na}$ (Sigma Aldrich, $99.95 \%$ ) and $0.145 \mathrm{mmol}(1.0 \mathrm{mg}) \mathrm{Li}$ (Sigma Aldrich, $99.9 \%$ ) were added. The Ta crucible was sealed under argon by arc welding. To protect the Ta crucible from oxidation, it was placed into a silica tube under argon atmosphere. In a tube furnace the crucible was heated to $760{ }^{\circ} \mathrm{C}$ with a rate of $50{ }^{\circ} \mathrm{C} \mathrm{h}^{-1}$. The temperature was maintained for 48 h and then lowered with $3.7^{\circ} \mathrm{C} \mathrm{h}^{-1}$ to $200^{\circ} \mathrm{C}$. Once the temperature reached $200^{\circ} \mathrm{C}$, the furnace was turned off and cooled down to room temperature. The Ta crucible was opened and Na was separated from the reaction products by evaporation at $320^{\circ} \mathrm{C}$ under vacuum ( 0.1 Pa ) for 18 h . From the inhomogeneous gray product, red needle-shaped single crystals (200-600 $\mu \mathrm{m}$ ) were isolated (cf. Fig. 1), enclosed in glass capillaries and sealed under argon atmosphere. X-ray diffraction data were collected at room temperature with a STOE IPDS I


Fig. 1. SEM micrograph of $\mathrm{LiBa}_{5} \mathrm{GaN}_{3} \mathrm{~F}_{5}$ crystals.
diffractometer. A numerical absorption correction using the programs XRED [12] and XSHAPE [13] was applied. The crystal structure was solved by using direct methods with SHELXS [14]. The refinement of the structure was carried out by the method of least-squares using SHELXL [14]. The chemical composition was confirmed by energy dispersive X-ray spectroscopy (EDX) using a JSM-6500F scanning microscope (Jeol) provided with a Si/Li EDX detector (Oxford Instruments, model 7418). Optical spectra of $\mathrm{LiBa}_{5} \mathrm{GaN}_{3} \mathrm{~F}_{5}$ were measured with a modified microcrystal spectrophotometer CARY 17 (Spectra Services, ANU Canberra, Australia) [15-17]. Calculations of the band gap were carried out with the program package WIEN2K [18] utilizing the structural data from the single-crystal structure refinement.

Further details of the crystal structure investigations can be obtained from the Fachinformationzentrum Karlsruhe, 76344 Eggenstein-Leopoldshafen, Germany (Fax: +49-7247-808-666; E-Mail: crysdata@fiz-karlsruhe.de) on quoting the depository number CSD-421592.

## 3. Results and discussion

### 3.1. Crystal structure

The crystal structure was solved and refined in orthorhombic space group Pnma (no. 62) with $a=15.456(3), b=5.707$ (1) and $c=12.259(3) \AA$. The crystallographic data of $\mathrm{LiBa}_{5} \mathrm{GaN}_{3} \mathrm{~F}_{5}$ is summarized in Table 1, the atomic coordinates and the isotropic displacement parameters are listed in Table 2.

In the crystal $\mathrm{LiBa}_{5} \mathrm{GaN}_{3} \mathrm{~F}_{5}$ zigzag chains of vertex sharing $\mathrm{LiF}_{6}$ octahedrons running along [010] (cf. Fig. 2). Perpendicular to these chains, Ba atoms are arranged in layers. Likewise perpendicular to [010] "isolated" trigonal planar $\left[\mathrm{GaN}_{3}\right]^{6-}$ ions are found. The $\mathrm{Ga}-\mathrm{N}$ bond-lengths range from 1.90 to $1.95 \AA$ (Fig. 3) and agree well with the sum of the ionic radii [19-21] as well as with typical Ga-N distances (e.g., $\mathrm{Sr}_{4} \mathrm{GaN}_{3} \mathrm{O}, \mathrm{Ga}-\mathrm{N}: 1.88-1.92 \AA$ Å) [5]. Similar [GaN $]_{3}{ }^{6-}$ ions have been found in $\mathrm{Sr}_{3} \mathrm{GaN}_{3}, \mathrm{Sr}_{6} \mathrm{GaN}_{5}$ [2], $\mathrm{Sr}_{4} \mathrm{GaN}_{3} \mathrm{O}$ and $\mathrm{Sr}_{4} \mathrm{GaN}_{3}\left(\mathrm{CN}_{2}\right)$ [5]. The coordination sphere of the trigonal planar $\left[\mathrm{GaN}_{3}\right]^{6-}$ ions can be described as three-capped trigonal prisms of $\mathrm{Ba}^{2+}$ atoms. In $\mathrm{Sr}_{3} \mathrm{GaN}_{3}$ [2] similar trigonal prism of $\mathrm{Sr}^{2+}$ atoms have been observed. The $\mathrm{N}^{3-}$ atoms are likewise surrounded in distorted octahedrons of five $\mathrm{Ba}^{2+}$ and one $\mathrm{Ga}^{3+}$ atom.

Table 1
Crystallographic data of $\mathrm{LiBa}_{5} \mathrm{GaN}_{3} \mathrm{~F}_{5}$.

| Formula | $\mathrm{LiBa}_{5} \mathrm{GaN}_{3} \mathrm{~F}_{5}$ |
| :--- | :--- |
| Crystal system | orthorhombic |
| Space group | Pnma (no. 62$)$ |
| Lattice parameters $(\AA)$ ) | $a=15.456(3), b=5.707(1), c=12.259(3)$ |
| Cell volume $\left(\AA^{3}\right)$ | $1081.3(4)$ |
| Formula units per unit cell | 4 |
| Density $\left(\mathrm{g} \mathrm{cm}^{-3}\right)$ | 5.531 |
| $\mu\left(\mathrm{~mm}^{-1}\right)$ | 20.41 |
| $\mathrm{~T}(\mathrm{~K})$ | $293(2)$ |
| $\mathrm{F}(000)$ | 1520 |
| Profile range | $5.8 \leq 2 \theta \leq 63.2$ |
| Index ranges | $-20 \leq h \leq 18$ |
|  | $-6 \leq k \leq 6$ |
| Independent reflections | $-15 \leq l \leq 15$ |
| Refined parameters | $1286[R($ int $)=0.072]$ |
| Goodness of fit | 89 |
| $R_{1}$ (all data); $R_{1}\left(F^{2}>2 \sigma\left(F^{2}\right)\right)$ | 1.057 |
| $w R_{2}($ all data $) ; w R_{2}\left(F^{2}>2 \sigma\left(F^{2}\right)\right)$ | $0.0230,0.0212$ |
| $\Delta \rho_{\text {max }}, \Delta \rho_{\text {min }}\left(\mathrm{e} \AA^{-3}\right)$ | $0.0512,0.0505$ |
|  | $1.56,-2.13$ |

$\mathrm{Li}^{+}$does not directly coordinate to the $\left[\mathrm{GaN}_{3}\right]^{6-}$ ions but is surrounded by six $\mathrm{F}^{-}$in a distorted octahedron. The $\mathrm{Li}-\mathrm{F}$ distances are ranging between 1.97 and $1.98 \AA$ for the equatorial F2 and F1 and $2.33-2.39 \AA$ for the axial F3 and F4 (cf. Fig. 3). The latter ones exceed significantly the sum of the ionic radii ( $2.06 \AA$ ) [19] while the equatorial $\mathrm{Li}-\mathrm{F}$ distances are slightly shorter than reported distances in other $\mathrm{LiF}_{6}$ octahedrons ( $\mathrm{K}_{2} \mathrm{LiAlF}_{6}, \mathrm{Li}-\mathrm{F}: 2.109 \AA$ ) [22,23]. In the second coordination sphere, the octahedron $\operatorname{LiF}_{6}$ is coordinated by eight $\mathrm{Ba}^{2+}$ in a cubic way, comparable with the cubic $\mathrm{Ca}^{2+}$-coordination of $\mathrm{Ti}^{4+}$ in $\mathrm{CaTiO}_{3}$. The coordination of $\mathrm{F}^{-}$ occurs in distorted octahedrons of five $\mathrm{Ba}^{2+}$ and one $\mathrm{Li}^{+}$. Only F2 is coordinated by two $\mathrm{Li}^{+}$and four $\mathrm{Ba}^{2+}$. The atomic distances $\mathrm{Ba}-\mathrm{F}$ and $\mathrm{Li}-\mathrm{F}$ are mentioned above. The thermal displacement parameter $U_{\text {iso }}{ }^{*} / U_{\text {eq }}$ of F4 is considerably higher (cf. Table 3) in comparison with the other values for fluorine atoms. One reason for this observation may be the fact, that - considering the octahedron $\mathrm{LiF}_{6}$ - the F4 atom is a "free" one and does not connect to the next octahedrons. Additionally, the $U_{11}$ value for F 4 is almost three times higher than the values for the other $\mathrm{F}^{-}$. The short distance between Ba4-F4 ( 2.83 Å) may be responsible for this observation. The F3 atom has an even shorter distance to Ba4 ( $2.76 \AA \AA$ ) and is also a non-bridging one. Here we can also observe a little higher $U_{\text {iso }}{ }^{*} / U_{\text {eq }}$ value.

Table 2
Atomic coordinates and isotropic displacement parameters $\left(\AA^{2}\right)$ of $\mathrm{LiBa}_{5} \mathrm{GaN}_{3} \mathrm{~F}_{5}$, standard deviations in parentheses.

| Atom | Wyckoff <br> position | $x$ | $y$ | $z$ | $U_{\text {iso }}{ }^{*} / U_{\text {eq }}$ |
| :--- | :--- | :--- | :--- | :--- | :--- |
| Ba1 | $4 c$ | $0.43534(2)$ | $1 / 4$ | $0.68180(2)$ | $0.0092(1)$ |
| Ba2 | $4 c$ | $0.51012(2)$ | $-1 / 4$ | $0.87769(3)$ | $0.0090(1)$ |
| Ba3 | $4 c$ | $0.34365(2)$ | $-1 / 4$ | $0.47028(3)$ | $0.0089(1)$ |
| Ba4 | $4 c$ | $0.24513(2)$ | $-1 / 4$ | $0.77891(3)$ | $0.0100(1)$ |
| Ba5 | $4 c$ | $0.16417(2)$ | $1 / 4$ | $0.63252(2)$ | $0.0109(1)$ |
| Ga1 | $4 c$ | $0.35910(4)$ | $1 / 4$ | $0.92467(4)$ | $0.0076(2)$ |
| F1 | $8 d$ | $0.4111(2)$ | $0.0121(5)$ | $0.3013(2)$ | $0.0151(6)$ |
| F2 | $4 a$ | $1 / 2$ | 0 | $1 / 2$ | $0.0131(7)$ |
| F3 | $4 c$ | $0.4075(3)$ | $-1 / 4$ | $0.6854(3)$ | $0.0171(8)$ |
| F4 | $4 c$ | $0.1779(4)$ | $-1 / 4$ | $0.9936(3)$ | $0.027(1)$ |
| N1 | $4 c$ | $0.4799(4)$ | $1 / 4$ | $0.8980(4)$ | $0.014(2)$ |
| N2 | $4 c$ | $0.2813(4)$ | $1 / 4$ | $0.8020(4)$ | $0.013(2)$ |
| N3 | $4 c$ | $0.3152(4)$ | $1 / 4$ | $1.0734(5)$ | $0.012(2)$ |
| Li1 | $4 c$ | $0.459(2)$ | $1 / 4$ | $0.4020(8)$ | $0.026(3)$ |

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