



Investigation of synchronous effects of multi-mesh regenerator and double-inlet on performance of a Stirling pulse tube cryocooler

M. Arablu, A. Jafarian*

Department of Mechanical Engineering, College of Engineering, Tarbiat Modares University, P.O. Box 14115-143, Tehran, Iran

ARTICLE INFO

Article history:

Received 16 July 2012

Received in revised form 24 September 2012

Accepted 28 September 2012

Available online 6 November 2012

Keywords:

Multi-mesh regenerator

Double inlet

Pulse tube cryocooler

Numerical simulation

ABSTRACT

In this paper synchronous effects of multi-mesh regenerator and double-inlet on performance of a Stirling pulse tube cryocooler (SPTC) have been considered. In this respect, a finite volume code was developed to simulate the SPTC. Set of governing equations were written in a general form such that all porous and non-porous sections of the system could be modeled. Results showed that synchronous application of double inlet and multi-mesh regenerator optimizes the phase shift between velocity and pressure at the warm end of the pulse tube, increases the regenerator's outlet pressure amplitude, decreases inertial and viscous losses in the hot end of the regenerator and consequently increases the COP of the system. Furthermore, it was observed that a minimum temperature of 60.3 K and COP of 0.03996 @ 80 K is attainable using optimum multi-mesh regenerator and double inlet; whereas, for a simple SPTC with a uniform mesh regenerator, a minimum temperature of 71.3 K and maximum COP of 0.0227 @ 80 K are concluded.

© 2012 Elsevier Ltd. All rights reserved.

1. Introduction

Nowadays, pulse tube cryocooler (PTC) is widely used to produce very low temperatures. This is due to its obvious benefits in contrast with other cryocoolers. However, its efficiency is lower than other cryocoolers and, since its invention by Gifford and Longworth in 1963 [20], many efforts have been dedicated to improve that. As a result of extensive researches in the PTC community, essential improvements achieved by two modifications: adding a reservoir via an orifice valve to the warm end of the pulse tube which led to phase shift between pressure and velocity; adding a second inlet valve, thereby reducing the mass flow through the regenerator led to a second phase shift with resulting improvement in cooling performance [19].

After four decades of experiences with PTC, it is already well-known that the regenerator as a key component plays a critical role on the cooling performance of the system. Therefore, extensive efforts have been focused on improvement in regenerator technology during the last two decades. These efforts are categorized into areas of materials and geometry, modeling, and measurement [9].

Although, there are several sources of energy losses in the regenerator; however, the most important parts of the losses occur due to the viscous and inertial losses [9]. Therefore, many efforts

have been dedicated to introduce a good relation for oscillating flow friction coefficient inside the regenerator [14].

Heat transfer between the gas and solid matrices in the regenerator micro-porous media is another important part of the losses. Therefore, exotic alloys and processes were developed so that the regenerator obtained sufficient thermal capacity in the cold region. There are many investigations published in the literature that deal with application of erbium alloys (Er_3Ni , $\text{Er}_6\text{Ni}_2\text{Sn}$, $\text{ErNi}_{0.9}\text{Co}_{0.1}$ and ErNi), lead spheres and even multi-mesh regenerators instead of simple micro-porous stainless steel regenerators to improve their performance [12,8,15,17–19].

Although increasing the interfacial heat transfer in the cold end of the regenerator has been the key reason for using multi-mesh regenerator with different materials in PTCs; however, optimum multi-mesh regenerator can reduce the inertial and viscous losses. Furthermore, the multi-mesh regenerator and double inlet can be simultaneously employed for further improvement of COP of SPTC in comparison with simple or multi-mesh regenerator SPTC. In this paper, synchronous effects of multi-mesh regenerator and double inlet on the cooling performance of a SPTC are investigated. In this respect, a 1D CFD code is developed to simulate the SPTCs.

2. Physical model and governing equations

Schematic of the double inlet SPTC is shown in Fig. 1. The physical domain includes all of the components between the piston head of the compressor and the end wall of the buffer. For simplicity three assumptions are applied as listed below:

* Corresponding author. Tel./fax: +98 21 82884909.

E-mail addresses: m.arablou@modares.ac.ir (M. Arablu), jafarian@modares.ac.ir (A. Jafarian).

Nomenclature

A	cross sectional area of gas flow (m^2)	x_{amp}	piston motion amplitude
A_L	interfacial heat transfer area (m^2)	x_d	length of compressor dead volume
A_1	small area of flow passing section through a joint	x_{comp}	compressor length
A_2	large area of flow passing section through a joint	1D	one dimensional
C_F	Forchheimer's inertial coefficient	$[A]_b^a$	$(A)_a - (A)_b$
C_p	constant pressure specific heat ($\text{J kg}^{-1} \text{K}^{-1}$)	Greek letters	
C_s	specific heat of solid ($\text{J kg}^{-1} \text{K}^{-1}$)	α	heat transfer coefficient
C_v	constant volume specific heat ($\text{J kg}^{-1} \text{K}^{-1}$)	γ	expansion angle of conical sections
E_{or}	energy flow through orifice	Δx	control volume length
J_{or}	correction factor of orifice	δt	time step
K	permeability (m^2)	δx	node distance
K_l	local energy loss coefficient	ε	porosity
L	conical section length or distance between adjacent nodes of an abrupt joint	ϕ	opening area ratio of screen
Nu	Nusselt number	θ	system angle with respect to gravity
P	pressure (Pa)	λ	matrix conductivity factor
Pr	Prandtl number	μ	viscosity (Pa s)
$\dot{Q}_c$	cooling power (W)	μ_w	viscosity at wall temperature (Pa s)
R	gas constant	ρ	density (kg m^{-3})
Re	Reynolds number	ω	frequency of piston motion (rad s^{-1})
S'_{gen}	entropy generation rate ($\text{W m}^{-3} \text{K}^{-1}$)	ζ	number of packed screens per length
T	temperature (K)	Subscripts	
V	volume (m^3)	cv	center of control volume
$\dot{W}$	P – V power of the compressor (W)	f	face of control volume
X_t	screen transverse pitch	i	control volume numbered “i”
X_u	amplitude of velocity oscillation	in	inlet section
d_h	hydraulic diameter	or	double inlet orifice
f	friction factor	$oscm$	oscillating mean value
g	gravity acceleration (m s^{-2})	out	outlet section
k	thermal conductivity ($\text{W m}^{-1} \text{K}^{-1}$)	s	solid matrix
l	mesh distance	$u1$	oscillating velocity at Reg. inlet
m	mass (kg)	$u2$	oscillating velocity at Reg. outlet
$\dot{m}$	mass flow rate (kg s^{-1})	Superscripts	
n	number of control volumes	j	current time
$\dot{q}_c$	cooling power/CHX solid matrix volume		
u	velocity (m s^{-1})		
x	longitudinal coordinate		

- 1D compressible flow of ideal gas;
- Specified heat transfer rate in the cold heat exchanger as cooling power;
- Sink of energy in compressor, transition line, after cooler and hot heat exchanger walls equal to radiative and convective heat transfer with ambient.

Taking an Eulerian point of view, the flow field unknowns are described as $\dot{m}(x, t)$, $P(x, t)$, $\rho(x, t)$ and $T(x, t)$. Finite volumes are defined to construct piecewise uniform approximations to these functions. Control volumes and the corresponding nodal points in their center are shown in Fig. 1. Solid circles ($1 \leq i \leq n$) store the nodal values of the temperature, pressure, density, mass, and gas properties. The dashed lines ($1 \leq i \leq n + 1$) store values of the mass flow rate, velocity, energy flow, and entropy flow. This staggered arrangement of finite volumes helps to prevent occurring non-physical numerical solutions, Checkerboard, in the Finite Volume Method. With aforementioned assumptions the continuity, momentum, energy of gas, energy of solid and the state equations can be respectively represented as follows:

$$\partial(m)_{cv,i}/\partial t = [\dot{m}]_{f,i+1}^{f,i} \quad (1)$$

$$\partial(mu)_{f,i}/\partial t = [\dot{m}u]_{cv,i}^{cv,i-1} + A_i[P]_{cv,i}^{cv,i-1} - [\varepsilon\mu/(K\rho) + c_F\bar{e}^2|u|/K^{1/2}]_{f,i}\dot{m}_i\delta x_i + m_{f,i}g \cos \theta. \quad (2)$$

$$\begin{aligned} \frac{\partial}{\partial t}(m \frac{u^2}{2} + mC_v T)_{cv,i} = & \left[\dot{m} \frac{u^2}{2} + \dot{m}C_p T - kA \frac{\partial T}{\partial x} \right]_{f,i+1}^{f,i} \\ & + \alpha_i A_{L,i} (T_{s,i}^j - T_i^j) + P_i \frac{dV_i}{dt}. \end{aligned} \quad (3)$$

$$\frac{\partial}{\partial t}(m_s C_s T_s)_{cv,i} = \left[\dot{m}_s k_s \frac{\partial T_s}{\partial x} \right]_{f,i+1}^{f,i} + \alpha_i A_{L,i} (T_i^j - T_{s,i}^j) + \dot{q}_c V_{s,cv,i}. \quad (4)$$

$$P_i^j V_i = m_i^j R T_i^j. \quad (5)$$

3. Numerical procedure

The implicit control volume method is used to solve Eqs. (1)–(5) which are changed to algebraic equations. The first order method is used for time-direction in first step and the second order for the second step. In x -direction, for diffusive and convective terms, central and second order upwind discretization methods are employed, respectively; except near the boundaries which the first order upwind is used. Herein, only second order method of time-direction discretization is presented.

3.1. Continuity and momentum equations

To explicit continuity equation for the pressure as a function of mass flow rate, the mass of gas is replaced from the equation of

Download English Version:

<https://daneshyari.com/en/article/1507701>

Download Persian Version:

<https://daneshyari.com/article/1507701>

[Daneshyari.com](https://daneshyari.com)