



High energy 523 nm Nd:YLF pulsed slab laser with novel pump beam waveguide design



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ABSTRACT

A laser diode pumped Nd:YLF master oscillator power amplifier (MOPA) green laser system with high pulse energy and high stable output is demonstrated. At a repetition rate of 50 Hz, 840 mJ pulse energy, 9.1 ns pulse width of 1047 nm infrared laser emitting is obtained from the MOPA system. The corresponding peak power is 93 MW. Extra-cavity frequency doubling with a LiB₃O₅ crystal, pulse energy of 520 mJ at 523 nm wavelength is achieved. The frequency conversion efficiency reaches up to 62%. The output pulse energy instability of the laser system is less than 0.6% for one hour.

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1. Introduction

Solid-state green lasers with high pulse energy and high beam quality have been rapidly developed recently. Green lasers perform better than IR lasers in many applications, such as large-area material treatment, precision micro-fabrication and underwater communication. Even more, high energy green lasers can be used to pump Ti:sapphire lasers and generate UV lasers. Compared with lamp-pumped solid-state green lasers, the diode-pumped system is reliable, efficient and stable. Traditionally, extra-cavity frequency doubling is a widely used method to obtain high energy green pulse for a master oscillator power amplifier (MOPA) system [1–4]. In 2000, Hirano et al. designed an external two-stage KTP crystal architecture. The system produced 131 W average power in green with a frequency conversion efficiency as high as 65.2%, and the beam quality factor M^2 was around 5.2 [1]. Using an external KTP crystal, Kiriya demonstrated 132 W of green pulse train output at 1 kHz, and the frequency conversion efficiency was 60% [2]. In a high power Q-switched Nd:YVO₄ MOPA system, Liu demonstrated 103.5 W of 532 nm wavelength laser output [3]. In 2013, Li reported a single-frequency Nd:YAG MOPA system with extra-cavity frequency doubling; at a repetition rate of 250 Hz, 532 nm green laser with 12.6 ns pulse width and 400 mJ pulse energy was obtained [4].

Nd:YLF, as a high quality fluoride, has been proven a very good

laser crystal to generate high energy laser. The laser transition of Nd:YLF is at 1047 nm (extraordinary) or at 1053 nm (ordinary); therefore the wavelength of second harmonic generation (SHG) is closer to the spectral window of sea water than 532 nm obtained from Nd:YAG lasers [5]. Nd:YLF material has other advantages, such as long fluorescence lifetime (480 μ s), natural birefringence and low thermal effect. The main challenge of Nd:YLF crystal for high energy laser application is its small tensile strength (33 MPa), which critically limits the maximum average pump power density. In 1993, Beach demonstrated a diode-end-pumped Nd:YLF laser emitting at 1047 nm wavelength, 100 mJ of pulse energy with 4 ns pulse width was obtained, and the repetition rate was 30 Hz [6]. In 1998, Clarkson reported a Nd:YLF laser at 1053 nm wavelength, pumped by two beam-shaped 20 W diode bars. The pulse energy was around 2.6 mJ at a repetition frequency of 1 kHz [7]. In 2004, Q-Peak company demonstrated a high-repetition rate Nd:YLF MOPA system; 45 W green light was achieved, but the pulse energy was low [8]. Similarly, Li reported a stable-unstable hybrid resonator Nd:YLF laser in 2008, and 15.1 mJ pulse energy with 7.1 ns pulse width at 523 nm wavelength was achieved [9]. Lu also designed a diode-end-pumped, conductively cooled intra-cavity frequency doubling Nd:YLF lasers for 527 nm or 523 nm output [10,11]. But until now, high energy diode-pumped 523 nm Nd:YLF green lasers have rarely been reported.

In this paper, we demonstrate a conductively cooled, high pulse energy Nd:YLF green laser at 523 nm. A MOPA system was used to generate 1047 nm fundamental wavelength, and a LiB₃O₅ (LBO) crystal was used for extra-cavity frequency doubling. Two

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cylindrical lenses were employed to compensate the thermal lens effect in two different directions. The MOPA system provided output energy of 840 mJ with pulse duration of 9.1 ns. Using extra-cavity frequency doubling method, 520 mJ pulse energy at 523 nm wavelength was achieved with 62% frequency conversion efficiency.

2. Experimental setup

The schematic of the experimental setup is shown in Fig. 1. The system consists of three parts, an EO Q-switched oscillator, Nd:YLF slab amplifiers and a frequency conversion stage. A U-type arrangement of the oscillator [11] could make the construction compact. The gain material was a pair of a-cut Nd:YLF crystal slabs (1.0 at.% Nd³⁺-doped), the dimension of each crystal was $4 \times 4 \times 12 \text{ mm}^3$. Total 24 mm slab in length was designed to absorb the pump energy completely. The crystals were end-pumped by two fiber coupled QCW 806 nm laser diodes operated at 50 Hz repetition rate with 480 μs pulse duration. The folded resonator consisted of a high-reflectivity mirror M1 with curvature radius of 2000 mm; two flat, high-reflectivity mirrors M2 and M3 ($R > 99.8\%$ at 1047 nm, high transmission at 806 nm), and a flat output coupler M4 with a transmission of 60% at 1047 nm. The total cavity length was 620 mm. A KD*P Pockel cell, a polarizer, and a quarter wave plate were used as the electro-optic Q-switcher.

The amplification stage involved two pre-amplifiers and two power-amplifiers. For every amplifier head, LD pumping and cooling architecture is shown in Fig. 2. The pump sources were quasi-cw diode stacks; each stack consisted of four 10 mm long diode bars (pre-amplifier) or six 10 mm long diode bars (power-amplifier). The maximum peak power of each LD bar was 150 W at the central wavelength of 806 nm. The full divergence angles in the fast and slow axes were approximately 40° and 10° . The pumping LD operated at a repetition rate of 50 Hz, with a pulse width of 400 μs . Six LD stacks were arranged in each pre-amplifier, 10 LD stacks were arranged in power-amplifier-1, and 11 LD stacks were adopted in power-amplifier-2. The LD stacks were staggered

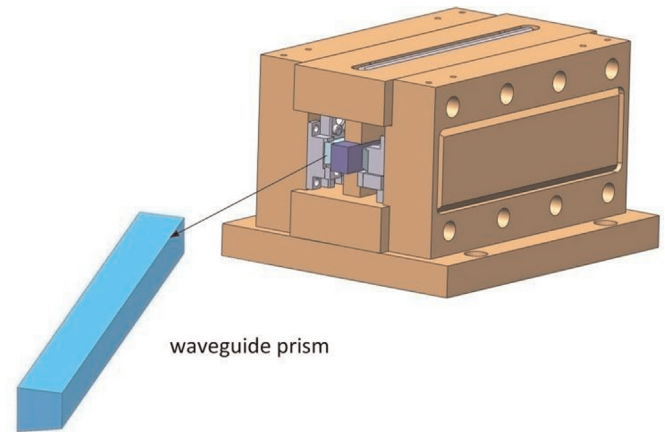


Fig. 2. Schematic of the amplifier head.

at two sides of each slab. In order to suppress parasitic oscillation, the Nd:YLF slabs were cut to 3° angle. For the sake of increasing the absorbing efficiency, 1 at.% concentration of Nd³⁺ was chosen. Considering the damage threshold of the crystal and the fluence of the signal, the beam sizes of the signal were controlled at $4.5 \times 4.5 \text{ mm}^2$, $6 \times 6 \text{ mm}^2$ and $8 \times 8 \text{ mm}^2$ respectively by expansion telescope. In order to increase the overlap efficiency and decrease the diffraction effect, slabs with size of $6 \times 6 \times 72 \text{ mm}^3$ were used in the pre-amplifiers, one $8 \times 8 \times 110 \text{ mm}^3$ slab was used in the power-amplifier-1, and a $10 \times 10 \times 120 \text{ mm}^3$ slab was used in the power-amplifier-2.

As we know, the main challenges involved in an LD array are its highly divergent output and large beam divergence. These drawbacks lead to the irregular and non-uniform output radiation of LD array. To overcome these problems, many researchers have applied lens ducts to shape the pumping beam [12,13], but the coupling efficiency of lens ducts is low, and the size is large. We therefore designed a trapezoid waveguide prism to couple the pump laser. The structure configuration of the waveguide prism is also shown in Fig. 2. The dimension of the waveguide prism was accurately calculated by ray tracing method. The height of the

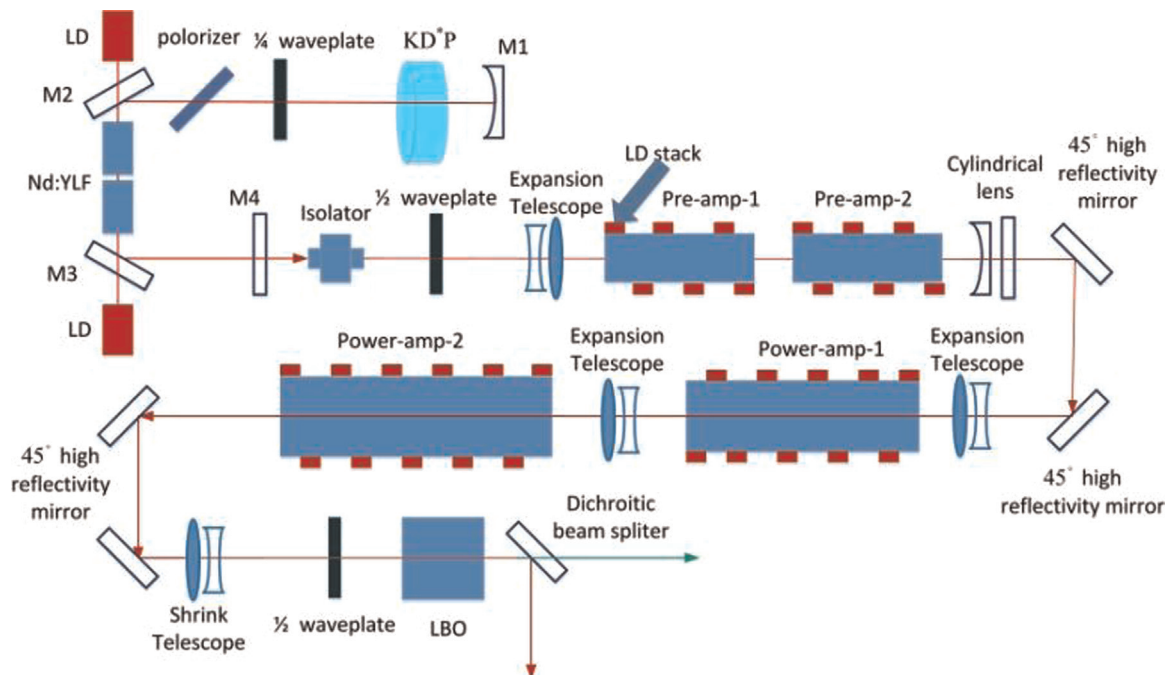


Fig. 1. Schematic of the experimental setup.

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