



First-principles investigation of structural, mechanical and electronic properties for Cu–Ti intermetallics



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ABSTRACT

To design a high-performance Cu–Ti intermetallic coating, first-principles calculation is conducted to systematically investigate the phase stability, mechanical and electronic properties of the Cu–Ti binary intermetallics. The phase stability and mechanical properties are studied by calculating the formation enthalpy and independent elastic constants, respectively. The calculated formation enthalpy indicates that β -Cu₄Ti, CuTi and CuTi₂ are stable phases at 0 K, while α -Cu₄Ti, Cu₃Ti, Cu₃Ti₂, and Cu₄Ti₃ are metastable phases. In addition, the mechanical properties of Cu–Ti intermetallics present a positive correlation with their formation enthalpy. The electronic structures of the Cu–Ti intermetallics are evaluated by analyzing the bonds character to reveal the bonding characteristics, which is crucial to the phase stability and mechanical properties. Among the Cu–Ti intermetallics studied, CuTi phase exhibits the highest stability, hardness and a higher brittleness among all Cu–Ti intermetallics, while other Cu–Ti intermetallics show good toughness. Based on the calculated results, a high strength Cu–Ti intermetallic coating consisting of hard CuTi particles on a ductile Cu₄Ti₃ matrix is proposed.

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1. Introduction

Copper (Cu) and Cu-based alloys are widely used for various functional components such as conductive springs and interconnections due to their excellent electrical and thermal conductivity [1]. Meanwhile, the strength and wear resistance of Cu-based alloys are critical factors in some specific applications, such as spot welding electrodes and electric railway wires [2], where Cu–Be alloy is now widely used [3]. However, the expensive cost and toxicity prevents its further utilization [4]. Fortunately, the Cu–Ti binary alloys containing 1–6 mol% Ti could work as a good candidate as their electronic and physical properties are close to Cu–Be alloys after ageing [5,6]. Moreover, both of the strength and hardness of Cu–Ti alloys could be improved due to the formation of Cu–Ti intermetallic compounds, while maintaining the high electrical conductivity of bulk alloy [7–9]. Therefore, surface modification by Cu–Ti intermetallic coatings could be a good way to improve the wear resistance of Cu-based alloys. In our previous work, a Cu–Ti intermetallic coating on C17000 Cu–Be alloy by thermally diffusion method have been successfully prepared [10], but the Cu–Ti intermetallic coating always contains various phases such as Cu₃Ti, Cu₃Ti₂, CuTi and CuTi₂, whose mechanical and electrical

properties are still unclear. Ghosh [11] investigated the structural stability of Cu–Ti intermetallics by first principles calculations, but the mechanical properties of Cu–Ti intermetallics were not revealed. In addition, there is still a lack of comparison between the crystalline and amorphous Cu–Ti intermetallics [12]. Therefore, to optimize the composition of Cu–Ti coatings and prevent the negative effects caused by the weak Cu–Ti intermetallics, a further systematic study from theoretical perspectives on the Cu–Ti intermetallics is needed.

In the present work, first-principles calculations are conducted to systematically study the phase stability, mechanical properties and electronic structure of Cu–Ti intermetallics. Specifically, the formation enthalpy, elastic constants, elastic modulus, microhardness, Debye temperature, density of state (DOS) and the difference of charge density are calculated. The results could provide suggestions for rational design of Cu–Ti intermetallic coatings on the surface of Cu alloys.

2. Computational method

The Cu–Ti binary phase diagram shows that there are six stable intermetallic compounds in the Cu–Ti binary system, i.e. CuTi₂, CuTi, Cu₄Ti₃, Cu₃Ti₂, Cu₂Ti and Cu₄Ti [13]. In addition, some other possible intermetallic compounds, such as CuTi₃, Cu₃Ti, were experimentally obtained [14,15]. Table 1 shows the built initial

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crystal structures of the above eight Cu–Ti intermetallics, as well as Cu and α -Ti with the crystallographic data from Refs. [11,13,22]. All calculations were performed with CASTEP code based on the density functional theory (DFT) [16]. The generalized gradient approximation (GGA) in Perdew–Burke–Eruzerh (PBE) formula was applied to describe the exchange–correlation potential [17]. The ion–electron interaction was evaluated by ultrasoft pseudopotentials of Vanderbilt type [18], and the Broyden–Fletcher–Goldfarb–Shanno (BFGS) algorithm was chosen to perform Geometry optimizations of these structures [19]. The cutoff energy of 600 eV was used for plane-wave expansions. The k-points meshes within Monkhorst–Pack framework were set as $10 \times 10 \times 10$, $6 \times 6 \times 9$, $6 \times 6 \times 2$, $6 \times 8 \times 8$, $10 \times 5 \times 6$, $11 \times 11 \times 2$, $8 \times 8 \times 2$, $11 \times 11 \times 2$, $1 \times 1 \times 1$, $5 \times 5 \times 3$, $9 \times 9 \times 2$, $8 \times 8 \times 8$, $12 \times 12 \times 12$ for Cu, α -Cu₄Ti, β -Cu₄Ti, Cu₃Ti, Cu₂Ti, Cu₃Ti₂, Cu₄Ti₃, CuTi, CuTi₂, CuTi₃ and α -Ti, respectively [20]. The geometry optimizations were considered to be converged when the energy change per atom, maximum residual force, maximum atomic displacement, maximum stress were less than 5×10^{-6} eV, 0.01 eV/Å, 5×10^{-4} Å and 0.02 GPa, respectively. The crystal structure of these compounds is shown in Fig. 1. Using the equilibrium structure, the elastic constants and modulus were calculated by the stress–strain method based on Hooke’s Law [21].

3. Results and discussion

3.1. Structural properties

The calculated lattice parameters of Cu–Ti intermetallics are listed in Table 1. Comparing with the literatures [11,13,22], it can

be seen that the calculated lattice parameters are close to the reported values with a deviation of less than 1%, indicating the good accuracy of the present work. The deviation between the calculated values and experimental ones are mainly caused by the GGA process and the temperature-dependence of lattice parameters. Especially for CuTi₃, the calculated lattice parameter shows an obvious difference from the experimental result for its metastable nature. In addition, Cu₃Ti₂, Cu₄Ti₃ and CuTi₂ have similar lattice parameters and structure of body-centered cubic with Cu and Ti atoms stacked along (100) plane alternately. Moreover, it can be readily seen that the Cu–Ti structure will turn into other low-energy structures of lower symmetry, as Ti content is less than 40 at.%.

3.2. Energies and phase stability

To evaluate the thermodynamic stability of the Cu–Ti intermetallics, the formation enthalpies (E_f) of these phases are calculated by the following equation:

$$E_f^{Cu_xTi_y} = \left[E_t^{Cu_xTi_y} - \left(x \cdot E_{solid}^{Cu} + y \cdot E_{solid}^{Ti} \right) \right] / (x + y) \quad (1)$$

where E_t and E_{solid} is the total energy of the calculated cell and the energy per atom of the pure constituent in solid state, respectively. x and y is the number of Cu and Ti atoms in the calculated cell, respectively. The calculated formation enthalpies of the Cu–Ti intermetallics are summarized in Table 2, compared with the available experimental results [23] and other theoretical values [11,24]. From Table 2, it can be seen that the calculated formation enthalpies are higher and closer to the experimental values than that in Ref. [11]. Based on the calculated formation enthalpies, the thermodynamic

Table 1
The calculated lattice parameters (a , b and c in nm) of Cu–Ti intermetallics compared with the available experimental and other theoretical results.

Phase	Space group	Lattice parameters			Refs.
Cu	Fm3-m	$a = 0.3599$			Experiment [11]
		$a = 0.3631$			This work
		$a = 0.3636$			Other work [22]
α -Cu ₄ Ti	I4/m	$a = 0.5840$		$c = 0.3620$	Experiment [13]
		$a = 0.58853$		$c = 0.36913$	This work
		$a = 0.58806$		$c = 0.37183$	Other work [11]
β -Cu ₄ Ti	Pnma	$a = 0.4522$	$b = 0.4344$	$c = 1.2897$	Experiment [13]
		$a = 0.45382$	$b = 0.4354$	$c = 1.29627$	This work
		$a = 0.45401$	$b = 0.43663$	$c = 1.29919$	Other work [11]
Cu ₃ Ti	Pmmn	$a = 0.5162$	$b = 0.4347$	$c = 0.4531$	Experiment [13]
		$a = 0.5437$	$b = 0.4322$	$c = 0.4446$	This work
		$a = 0.54658$	$b = 0.43299$	$c = 0.44342$	Other work [11]
Cu ₂ Ti	Amm2	$a = 0.4363$	$b = 0.7997$	$c = 0.4478$	Experiment [13]
		$a = 0.44192$	$b = 0.79922$	$c = 0.4579$	This work
		$a = 0.44285$	$b = 0.79723$	$c = 0.45934$	Other work [11]
Cu ₃ Ti ₂	P4/nmm	$a = 0.3130$		$c = 1.3950$	Experiment [13]
		$a = 0.3155$		$c = 1.3992$	This work
		$a = 0.31626$		$c = 1.39514$	Other work [11]
Cu ₄ Ti ₃	I4/mmm	$a = 0.3120$		$c = 1.9940$	Experiment [13]
		$a = 0.3146$		$c = 1.9920$	This work
		$a = 0.31525$		$c = 1.98531$	Other work [11]
CuTi	P4/nmm	$a = 0.3108$		$c = 0.5887$	Experiment [13]
		$a = 0.3124$		$c = 0.5921$	This work
		$a = 0.31256$		$c = 0.59152$	Other work [11]
CuTi ₂	I4/mmm	$a = 0.2944$		$c = 1.0786$	Experiment [13]
		$a = 0.29375$		$c = 1.07791$	This work
		$a = 0.29407$		$c = 1.06944$	Other work [11]
CuTi ₃	P4/mmm	$a = 0.4158$		$c = 0.3594$	Experiment [11]
		$a = 0.3979$		$c = 0.3978$	This work
		$a = 0.39690$		$c = 0.39691$	Other work [11]
α -Ti	P6 ₃ /mmc	$a = 0.2944$		$c = 0.4668$	Experiment [11]
		$a = 0.2945$		$c = 0.4627$	This work
		$a = 0.29229$		$c = 0.46271$	Other work [11]

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