

# Spontaneous strain in high-temperature superconductor $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$

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## Abstract

The lattice parameters of  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$  were measured by an x-ray Rietveld method between 15 K and room temperature. A change in a lattice parameter could be detected at superconducting phase transition. It is pointed out, for the first time, that the change can be attributed to a spontaneous strain in the superconducting phase caused by coupling between a superconducting order parameter and the strain.

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## 1. Introduction

Structural changes in high-temperature superconductors have been studied extensively since the discovery of superconducting phase transition in  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$  using high-resolution diffractometers. Nevertheless, no significant anomalies could be detected in connection with the superconducting phase transition. Recently, an accurate lattice parameter measurement has been carried out on orthorhombic  $\text{YBa}_2\text{Cu}_3\text{O}_{6.5}$  with a superconducting transition temperature  $T_c$  of 55 K by high-angle double-crystal x-ray diffractometry: The change in orthorhombicity  $2(b-a)/(a+b) = \tan^{-1}(b/a) - \pi/4$  at  $T_c$  could be clearly detected [1]. The introduction of a new critical exponent was considered necessary for explaining the difference between the observed orthorhombicity and that extrapolated from the normal state. An anomaly of the lattice parameter  $a$  of hexagonal  $\text{MgB}_2$  was clearly observed below its superconducting transition temperature  $T_c = 39$  K by high-resolution pulse neutron powder diffraction [2]. This result was analysed by introducing two Einstein temperatures, 222 K and 69 K, and two Grüneisen parameters, 1.33 and  $-0.304$ . The anomaly of the lattice parameter  $a$  of  $\text{MgB}_2$  was found to be independent of the onset of superconductivity.

Spontaneous strain is typically produced as a secondary order parameter through the coupling between the strain  $e$  and the primary order parameter  $Q$  that is an atomic shift in structural phase transitions. The free energy  $G(Q, T)$  near the transition temperature  $T_c$  can be written in its simplest form as follows using a Landau potential:

$$G(Q, T) = G_0(T) + \frac{1}{2}A(T - T_c)Q^2 + \frac{1}{4}BQ^4 + \frac{1}{2}Ce^2 - DeQ^2,$$

where  $A$ ,  $B$ ,  $C$  and  $D$  are temperature  $T$ -independent positive constants. From the equilibrium condition for the strain, we can obtain the relation  $e = (D/C)Q^2$ . The primary-order parameter  $Q$  is zero above  $T_c$ , whereas  $Q$  is nonzero below  $T_c$ . The spontaneous strain is defined as follows by the low-temperature phase lattice parameter  $a_{LT}$  and high-temperature phase lattice parameter  $a_{HT}$  when extrapolated to the same temperature [3]:  $e = (a_{LT} - a_{HT})/a_{HT}$ . In the case of tetragonal–orthorhombic structural phase transition, the orthorhombicity  $2(b-a)/(a+b)$  is twice the spontaneous strain in the  $ab$ -plane, because the lattice parameter of an average structure (i.e. a hypothetical tetragonal structure at a low temperature) is  $(a+b)/2 \approx (ab)^{1/2}$  as a first approximation [3]. We have been studying the phenomenological relations of order parameters in the antiferroelectric phase transitions of the perovskites  $\text{PbZrO}_3$  and  $\text{PbHfO}_3$  by

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X-ray and neutron diffraction analyses [4,5]: The spontaneous strain perpendicular to the plane in which an antiferroelectric order parameter appears could be explained well by conventional coupling. The strain in the plane, however, showed a complex dependence on temperature, which is difficult to explain in terms of the coupling.

If the spontaneous strain in high-temperature superconductors is caused by the same mechanism as that in structural phase transitions, the strain perpendicular to the  $\text{CuO}_2$  plane is expected to show a temperature dependence proportional to the square of a primary order parameter, which is a superconducting gap, because a gap appears in the  $\text{CuO}_2$  plane. The lattice parameter  $c$  of YBCO changes considerably with decreasing temperature above  $T_c$ , which prevents drawing a definite conclusion. The saturation of the lattice parameter above  $T_c$  is expected for superconductors with low  $T_c$ 's. Thus, we chose an optimally doped LSCO sample,  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ . We carried out precise lattice parameter measurements between 15 K and room temperature using a conventional x-ray powder diffractometer. We detected a change in strain perpendicular to the  $\text{CuO}_2$  plane in the superconducting phase, which could be explained by the coupling between the strain and the superconducting gap. The strain of  $\text{MgB}_2$  was also analysed by considering the coupling. To the best of our knowledge, this is the first indication of the coupling between the superconducting order parameter and the strain.

## 2. Experimental details

$\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$  samples were prepared by a ceramic solid-state reaction. Appropriate amounts of  $\text{La}_2\text{O}_3$ ,  $\text{SrCO}_3$  and  $\text{CuO}$  were mixed and calcined in air at 950 °C for 20 h, then pulverized, pressed into pellets and sintered in air at 1000 °C for 40 h. The obtained products were then ground, pressed into pellets again, and sintered in air at 1050 °C for 40 h. The samples were cooled to room temperature in a furnace.

The electric resistance of a sample was measured by a conventional four-probe method with a current of 1 mA. Au wire probes were attached to bar-shaped ceramic samples using Ag paste. The resistance slightly increased at around 70 K and disappeared at about 35 K.

X-ray diffraction patterns were obtained using a Rigaku X-ray diffractometer, RINT2500, with a graphite counter monochromator and an X-ray generator with a rotating Cu anode. The generator was operated at 50 kV and 300 mA. A powder sample was obtained by grinding the ceramics. A platelike powder sample was mounted on a sample holder made of copper. The sample was fixed in a closed-cycle He gas refrigerator mounted on the diffractometer. The sample was cooled from room temperature to 15 K. The diffraction patterns between 20° and 140° were measured at a scanning speed of  $2\theta = 0.4^\circ/\text{min}$ . Data were collected at every  $2\theta = 0.02^\circ$ . The diffraction patterns were analysed by the Rietveld method to obtain accurate lattice parameters using the computer program RIETAN-2000 [6]; calculations were carried out by a conjugate direction method.  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$  shows tetragonal–orthorhombic structural phase transition at about 180 K [7]. We used the space groups  $I4/mmm$  (No. 139) [8]

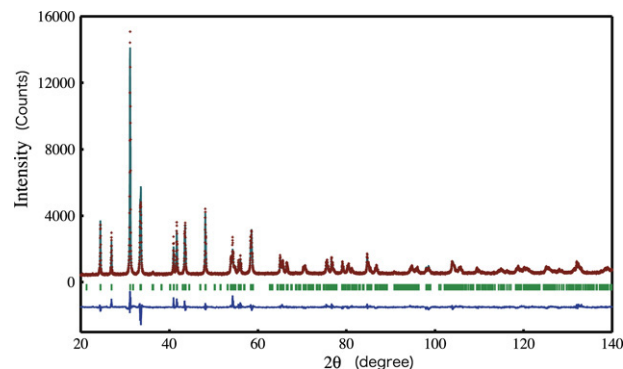


Fig. 1. Observed x-ray powder diffraction pattern and best-fit Rietveld refinement profile for orthorhombic  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$  at 15 K.

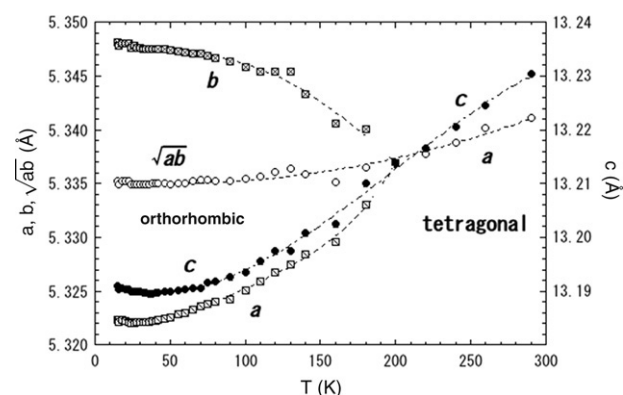


Fig. 2. Lattice parameters  $a$ ,  $b$ ,  $(ab)^{1/2}$  and  $c$  of  $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$  vs temperature based on X-ray powder diffraction measurements. The orthorhombic phase lattice parameter  $c$  is determined along the direction of  $c$  in the tetragonal phase to enable an easy comparison. The parameter  $a$  in the tetragonal phase is increased by a factor of  $\sqrt{2}$  for easy comparison with  $a$ ,  $b$  and  $(ab)^{1/2}$  of the orthorhombic phase. Standard deviations are smaller than the symbols. Lines serve as visual guides.

and  $\text{Cmca}$  (No. 64) [9] for the tetragonal and orthorhombic phases, respectively.

## 3. Results and discussion

The result of the profile fit at 15 K for  $x = 0.15$  is shown in Fig. 1, as an example. There are no clear impurity peaks. The obtained lattice parameters are delineated in Fig. 2 as a function of temperature. Tetragonal–orthorhombic structural phase transition, which produces a superstructure, occurs at around 180 K. The orthorhombic lattice parameter  $c$  is determined along the direction of the tetragonal lattice parameter  $c$  in Fig. 2 to enable an easy comparison of the lattice parameters. The tetragonal lattice parameter  $a$  shown in Fig. 2 is increased by a factor of  $\sqrt{2}$  for easy comparison with  $a$ ,  $b$  and  $(ab)^{1/2}$  of the orthorhombic phase. Standard deviations are smaller than the symbols in Fig. 2. A distinct change in the lattice parameter  $c$  can be observed near the superconducting transition temperature.

The lattice parameters  $c$  of the tetragonal and orthorhombic phases are shown in Fig. 3 as a function of temperature along with the electric resistance for comparison. The lattice parameter is expected to show saturation at low temperatures if superconducting phase transition does not occur as the

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