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X-ray emission as a diagnostic from pseudospark-sourced electron beams

ABSTRACT

imaging of materials.

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1. Introduction

Since its conception as an electron beam (e-beam) source at the university of Erlangen in 1978 [1], the pseudospark discharge has been proven to be of note both as an e-beam source and for its use in high-power switching applications [2,3]. The discharge operates within a hollow cathode/planar anode configuration on the left-hand side of the hollow cathode analogue to the Paschen minimum, typically at pressures of 6.5–65 Pa. It is notable for its fast rise time and high discharge current, as well as for the generation of a high quality electron beam in the conductive phase of the discharge. Such beams possess high current density of up to 10^4 A cm^{-2} , high brightness up to $10^{12} \text{ A m}^{-2} \text{rad}^{-2}$ [4] and require no guiding magnetic field. This is due to the presence of an ion channel which is formed with the generation of the electron beam [5–9].

Because of its special discharge characteristics, the pseudospark discharge has gained considerable attention during the last 30 years, particularly with regard to its breakdown characteristics and the underlying plasma physics responsible for high current emission, and for its potential application in various fields [10,11]. For example, the pseudospark has been shown to act as a source of soft X-rays by means of X-ray fluorescence [12–14]

and X-ray bremsstrahlung [15]. In this paper, measurements of electron beams of 3 mm diameter from a four-gap pseudospark, impacting on a molybdenum plate (which also served as a witness plate of the electron beam) for the generation of X-rays and X-ray imaging of a micro-sized object are presented. In addition to their role in imaging, it has been found that these X-rays may function as a diagnostic tool for indicating the presence of the generated

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2. Materials and methods

electron beam.

X-ray emission has been achieved using an electron beam generated by a pseudospark low-pressure dis-

charge and utilised as a diagnostic for beam detection. A 300 A, 34 kV PS-sourced electron beam pulse of

3 mm diameter impacting on a 0.1 mm-thick molybdenum target generated X-rays which were detected

via the use of a small, portable X-ray detector. Clear X-ray images of a micro-sized object were captured

using an X-ray photodetector. This demonstrates the inducement of proton induced X-ray emission (PIXE) not only as an indicator of beam presence but also as a future X-ray source for small-spot X-ray

The X-ray experimental setup with the PS discharge is shown in Fig. 1. The discharge chamber itself is a four-gap PS discharge structure consisting of a planar anode, a planar cathode with a cylindrical hollow cavity, three stainless steel inter-electrodes of 3 mm thickness and four Perspex insulation discs of 4 mm thickness. Both the anode and cathode have an on-axis hole of 3 mm diameter. The hollow cathode cavity was made of stainless steel, having a length of 50 mm and outer and inner diameters of 63 and 50 mm respectively. These dimensions were based on the requirements of the hollow cathode effect [4]. To reduce the inductance in the discharging circuit and to act as an energy source to sustain the discharge, an external energy storage capacitor of 428 pF was placed across the cathode and anode, consisting of three chains of capacitors in parallel. Each chain consisted of seven 15 kV, 1000 pF resin dipped ceramic capacitors connected in series.







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Fig. 1. Schematic experimental setup.

A mechanical rotary pump evacuated the experimental system from the anode end through a vacuum valve. The working gas was air and entered the chamber through a very fine controlled needle valve at the anode side and its pressure was measured by a Baratron-type vacuum gauge located at the anode together with a display meter. The hollow cathode was connected through a 30 M Ω charging resistor to a negative voltage source (-100 kV, 40 mA dc power supply). The anode was grounded. A capacitive voltage probe of sub-ns rise-time was connected to the cathode to measure the applied voltage as well as the discharge voltage while a Rogowski coil was connected next to the external storage capacitor in the discharge circuit to measure the discharge current. The measurement of the beam current was achieved using another Rogowski belt between the anode and the drift tube, as shown in Fig. 1.

The capacitive voltage probe with fast response and high load impedance was used to measure the fast risetime (\sim 20 ns), short duration (20–50 ns) discharge voltage of magnitude up to 46 kV. It is based on the derivation and then integration of the signal to be measured, which in this case was the applied as well as the discharge voltage. A pure resistive or capacitive voltage divider [16] was not suitable as it would have affected the charging circuit. In the pseudospark experiments, the voltage was measured using a 100 pF capacitive probe with a resistance of 50 Ω and a corresponding sensitivity and risetime of 1.38 kV/V and 15 ns, respectively.

Both the discharge current and the beam current were measured by Rogowski coils as shown in Fig. 1. In principle a Rogowski coil is a current transformer suitable for high frequency current measurement. Its primary is the current to be measured and the secondary often consists of many turns depending on the attenuating ratio that is required. It can operate in two modes, the self-integration or differential mode of operation with the self-integration operation having the larger sensitivity. In the pseudospark experiments, the Rogowski coils were made by winding ten turns of coil around a ferrite core. The coil was made of 0.7 mm diameter wire and operated in a self-integration mode with a 10Ω external resistor connected in series [17]. The calibrated sensitivity of the Rogowski coils was 1.37 A/V with a risetime of less than 10 ns [18].

A 0.1 mm-thick molybdenum target for effective X-ray production was placed immediately after the PS anode. Typically, target materials must be made of metals with high melting points, and possess both good thermal conductivities and low vapour pressures [19]. Additionally, the higher the atomic number of the metal in the target, the higher the efficiency of X-ray production [20]. The most commonly used target metal is tungsten, which has a high atomic number, a high melting point and the lowest vapour pressure of all metals. For applications in which the X-rays required are of much lower energy, the anode usually consists of molybdenum rather than tungsten. In this experiment, the molybdenum target also served as a witness plate for the PS electron beam. The cross-sectional shape and size of the electron beam from the PS could be recorded after a number of shots by repeating the PS discharge. A small, portal X-ray detector was situated within the experimental area to allow for instantaneous confirmation of the discharge while an X-ray photodetector was located opposite the anode with a micro-sized object made of two crossed metal wires of 100 um in diameter attached for X-ray photography.

The X-ray detector used was made by Photonic Science Ltd. model CoolView FDI 1:1. The available input area of the detector was 9.0×6.7 mm with 1392×1040 pixels and a spatial resolution of 6.45 µm. The detector achieves digital X-ray imaging in an indirect way. An X-ray source sends a beam of X-ray photons through an object and any X-ray photons not absorbed by the object strike a layer of scintillating material that converts them into visible light photons. These visible light photons then form an image on a charge-coupled device (CCD) camera through a fibre-optic coupler with unit magnification. The scintillator of the detector is a polycrystalline layer of Gadolinium Oxysulphide with the density of 10 mg/cm^2 and $30 \mu \text{m}$ thickness optimised for resolution with X-ray energies of 5–60 keV. The integral electronic shutter can be selected from 1 ms to 10 min with automatic dark subtraction. The analogue-to-digital signal converter has 12 bits digitization. An aluminium foil light exclusion membrane if required could be placed in front of the scintillator layer.

The pseudospark was configured for operation in a free-running, as opposed to triggered, mode. At the beginning of a single shot experiment, the pseudospark chamber was evacuated through a vacuum port at the anode by a mechanical pump to ~0.4 Pa. After closing the vacuum pump by an electromagnetic valve, the cathode was charged to a given voltage and gas slowly filled the chamber by adjustment of the fine control needle valve (at a very slow rate of $dp/dt \sim 0.65$ Pa/s in order to obtain a uniform gas pressure) resulting in an increase in gas pressure until breakdown occurred. Critical breakdown pressure was determined by reading the display metre of the Baratron-type pressure gauge. Applied voltage, discharge



Fig. 2. Typical discharge voltage, discharge current and beam current.

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