

# Tunable periodic graphene antidot lattices fabricated by e-beam lithography and oxygen ion etching



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## ABSTRACT

Intrinsic monolayer graphene has no band gap between its conduction and valence bands, which limits its application in many aspects as a semiconductor. Antidot lattices by constructing periodic holes on graphene have been proved to be an intriguing strategy to introduce a band gap into graphene. Here we used the e-beam lithography (EBL) combined with the oxygen reactive ion etching (RIE) to fabricate tunable antidot lattices with different and uniform regularly spaced holes on graphene. In this way, tunable periodic graphene nanostructures with the dimensions ranging from  $\sim 20$  nm to several hundreds of nanometers can be fabricated by controlling exposure dose and etching time.

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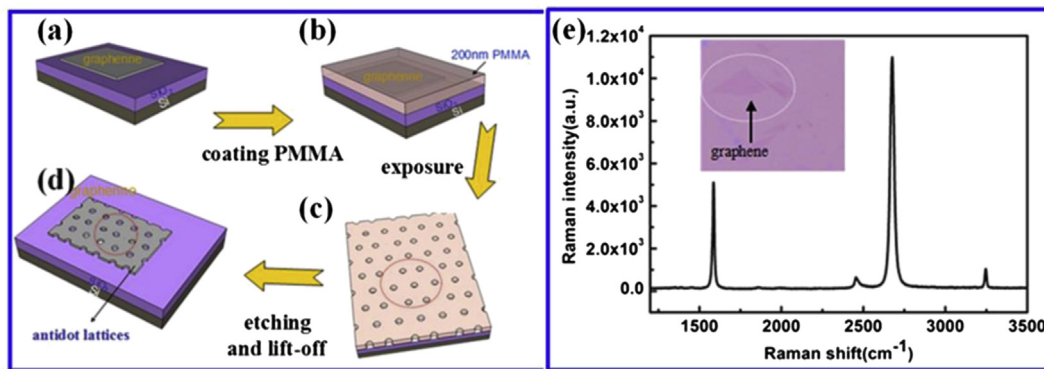
## 1. Introduction

Since its experimental discovery in 2004 [1], graphene has attracted extraordinary attention owing to its many outstanding physical properties, such as high mobility [2], quantum electronic transport [3], room-temperature quantum Hall effect [4], high elasticity [5] and good optical transparency [6]. Many potential electronic applications of graphene have also been demonstrated or proposed as transparent conductors [7], spintronic devices [8], sensors [9,10] and flexible electronics devices [11,12], but the deficiency of band gap around the Fermi level limits graphene application as a semiconductor material. However, graphene can show tunable band gap characteristics by adjusting size, shape, or surface functionalization [13]. Several kinds of structures have been proposed to introduce a band gap into graphene, such as nanoribbon, nanomesh and quantum dots [14–16]. Among these structures, graphene nanoribbon devices are the most commonly used, but they have relatively low driving currents or conductance for graphene-based semiconductor devices [17]. Therefore, in order to improve the performances of the graphene-based devices, the dense arrays of ordered nanoribbons will be required, which is still a great challenge.

Recently, a strategy of constructing periodic holes on graphene to form graphene antidot lattices has been extensively proposed. Graphene antidot lattices consist of the holes of a periodic arrangement in a graphene sheet [18]. Theoretical calculations have predicted that antidot lattices can introduce a band gap in graphene, and the width of opened gap can be tuned by the size, shape, and symmetry of both the hole and the lattice cell [19,20]. Some experimental studies have demonstrated that graphene antidot lattice have an effective energy gap of 100 meV and field effect transistors have an ON – OFF current ratio of up to 10, which demonstrate the utility of these structures for applications [21,22]. In comparison with a single nanoribbon, graphene antidot lattices can ensure a higher drive current when used as a conduction channel. More importantly, it is also an important foundation for constructing graphene superlattices with various nanoscopic functional units. To widening their application range in electronic devices, various methods such as lithography, nanoimprinting, and chemical techniques have been used to fabricate patterning graphene. Comparing other approaches of preparation of graphene antidot lattices, EBL combined with RIE as the most commonly used method have a lot of advantages, such as accurate alignment accuracy, high pattern resolution, large scale area, good uniformity and arbitrary control the size of antidot lattice. However, how to fabricate large-scale, uniform and high-density antidot lattice by using EBL and RIE is still a great challenge due to the well-known electron scattering effects in common electron beam photoresists and the uncertainty during reactive ions etching. In this work, we carried out EBL and RIE technique to fabricate tunable antidot

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**Fig. 1.** A schematic illustration of the simple and scalable graphene patterning processes. (a) Natural graphene deposited on SiO<sub>2</sub> (300 nm)/Si substrate. (b) A layer of PMMA was coated on the top of graphene. (c) EBL was used to pattern the PMMA into desired shapes at desired locations. (d) RIE was used to etch graphene/PMMA film. (e) Raman spectrum of monolayer graphene on SiO<sub>2</sub>/Si substrate and the inset shows the optical microscope photograph of the graphene.

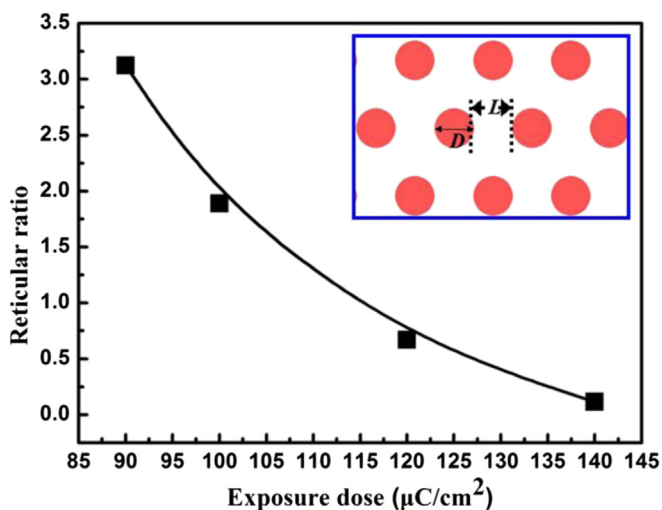
lattices, and different and uniform regularly spaced holes on graphene were achieved by controlling exposure dose and etching time. Finally, large-scale, uniform and dense grid of antidot lattices on graphene were fabricated successfully. Our results indicate that accurate control in both exposure dose and etching time is very desired approach to fabricate tunable antidot lattices nanostructures on graphene.

## 2. Experiment

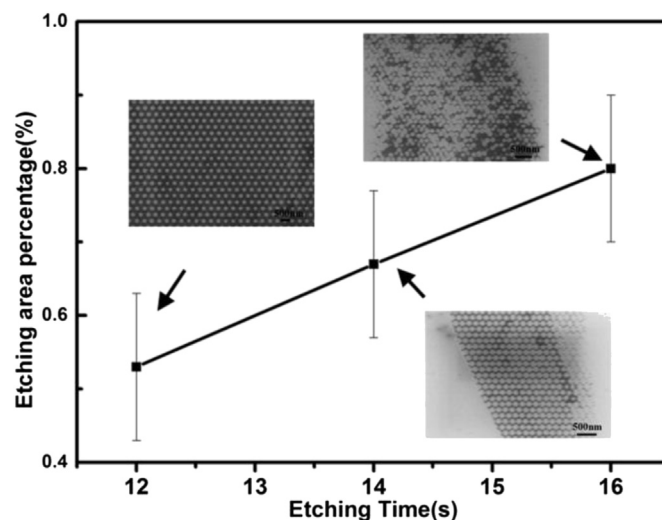
**Fig. 1** shows the schematic illustration of our fabricating procedure. Initially, the monolayer graphene was deposited on SiO<sub>2</sub>(300 nm)/Si substrate by mechanical exfoliation [**Fig. 1(a)**]. Then a layer of PMMA (about 200 nm) was coated on the top of graphene [**Fig. 1(b)**]. By changing exposure dose of EBL system (Raith 150), antidot lattices with different periods were formed on the PMMA layer [**Fig. 1(c)**]. After exposure, the sample was developed in methyl isobutyl ketone: isopropyl alcohol (IPA) (1:3) for 40 s and IPA for 30 s, and then was dried using pure nitrogen. Finally, short time (around 10–16 s) based on oxygen etching (O<sub>2</sub> flow: 100 sccm; Pressure: 100 mTorr; Power: 100 W) didn't influence the overall quality of antidot lattices on graphene flake. After lift-off was done in hot acetone ( $T = 80\text{ }^{\circ}\text{C}$ ) about 5 min in order to remove photoresist on the graphene [**Fig. 1(d)**].

## 3. Results and discussion

**Fig. 1(e)** shows a typical Raman spectrum of graphene on SiO<sub>2</sub>/Si substrate. There are two prominent peaks (G and 2D) and two small peaks (G' and 2D'). Two intense peaks are the G band at  $1585\text{ cm}^{-1}$ , and the 2D band at  $2676\text{ cm}^{-1}$ , respectively. The G peak corresponds to the stretching vibration mode, E<sub>2g</sub> phonon at the Brillouin zone center. The 2D peak originates from a two phonon double resonance Raman process, and it is always present for crystalline carbon materials [23]. The intensity of 2D peak is more than twice that those of G peak and the full width at half maximum (FWHM) of 2D band is  $29\text{ cm}^{-1}$ , suggesting that the sample is monolayer graphene. The Raman G' band ( $\sim 2450\text{ cm}^{-1}$ ) has a small intensity, which originates from a combination of the zone boundary inplane longitudinal acoustic phonon and the inplane transverse optical phonon modes [24]. The 2D' peak ( $\sim 3250\text{ cm}^{-1}$ ) can be assigned to the second harmonic of D peak ( $\sim 1620\text{ cm}^{-1}$ ) [25], which does not require the existence of defects for its activation and is always present in the Raman spectrum of high-quality monolayer graphene. Moreover, it is noted that the D band in the Raman spectrum is not obvious, indicating the good crystalline structure of the graphene sample with few defects and impurities. The inset of **Fig. 1(e)** shows the optical microscope photograph of



**Fig. 2.** The change tendency of reticular ratio with different exposure dose, and the reticular ratio ( $L/D$ ) represents the ratio of hole space ( $L$ ) to hole diameter ( $D$ ).



**Fig. 3.** The measured average etching area percentage of graphene antidot lattices change with the duration of plasma etching under a given O<sub>2</sub>-based RIE condition.

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