



# Finite-time thermodynamics optimization of absorption refrigeration systems: A review

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## ABSTRACT

This paper presents a literature review of the optimization of absorption refrigeration systems based on finite-time thermodynamics. An overview of the various objective functions is presented.

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## 1. Introduction

The absorption refrigeration systems are thermodynamic processes which produce cold thanks to thermal energy. Then, they exchange heat with at least three sources at different temperatures without receiving work. A three-heat-source reversible refrigerator operates between heat hot reservoir, heat cold reservoir and heat sink. When  $T_H$ ,  $T_L$  and  $T_O$  denote the temperatures of heat hot reservoir, heat cold reservoir and heat sink respectively, the coefficient of performance for three-heat-source reversible refrigerators is

expressed as:  $\varepsilon_r = [(T_H - T_O)/T_H][T_L/(T_O - T_L)]$  [1]. This expression reveals the product of thermal efficiency of Carnot cycle for heat engines working between  $T_H$  and  $T_O$  and coefficient of performance of reversible Carnot refrigerator producing cold at  $T_L$  and rejecting heat at  $T_O$ :  $\varepsilon_r = \eta_C \times \varepsilon_C$  with  $\eta_C = (T_H - T_O)/T_H$  and  $\varepsilon_C = T_L/(T_O - T_L)$ . In classical thermodynamics, the efficiency of a cycle operating on reversibility principles proposed by Carnot [2] became the upper bound of thermal efficiency for heat engines that work between the same temperature limits. This equally applies to the coefficient of performance of refrigeration cycles that execute a reversed Carnot cycle (Carnot refrigerator). This implies that the coefficient of performance defined above is the maximum coefficient of performance for three-heat-source refrigerators from the point of view of classical thermodynamics.

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**Nomenclature**

$A$	total heat-transfer area ( $\text{m}^2$ )	$T_A$	temperature of the absorber-side heat sink (K)
$A_A$	heat-transfer area of absorber ( $\text{m}^2$ )	$T_C$	temperature of the condenser-side heat sink (K)
$A_C$	heat-transfer area of condenser ( $\text{m}^2$ )	$T_{env}$	temperature in environmental conditions
$A_L$	heat-transfer area of evaporator ( $\text{m}^2$ )	$T_H$	temperature of the heat source (K)
$A_H$	heat-transfer area of generator ( $\text{m}^2$ )	$T_L$	temperature of the cooled space (K)
$A_O$	$A_A + A_C$	$T_O$	$T_A = T_C$
$ECOP$	ecological coefficient of performance	$U_H$	overall heat-transfer coefficient of generator ( $\text{W K/m}^2$ )
$I$	internal irreversibility parameter	$U_L$	overall heat-transfer coefficient of evaporator ( $\text{W K/m}^2$ )
$K_{LC}$	heat leak coefficient (W K)	$U_O$	overall heat-transfer coefficient of absorber and condenser ( $\text{W K/m}^2$ )
$K_H$	thermal conductance of heat source ( $\text{W K}^{-1}$ )	$\dot{W}$	power output (W)
$K_L$	thermal conductance of cooled space ( $\text{W K}^{-1}$ )		
$K_O$	thermal conductance of heat sink ( $\text{W K}^{-1}$ )	<i>Symbol</i>	
$ncu$	national currency unit	$\varepsilon$	coefficient of performance for absorption refrigerators
$\dot{Q}_A$	heat reject load from absorber to heat sink (W)	$\varepsilon_C$	coefficient of performance of reversible Carnot refrigerator
$\dot{Q}_C$	heat reject load from condenser to heat sink (W)	$\eta_C$	thermal efficiency of Carnot cycle
$\dot{Q}_L$	heat input load from cooled space to evaporator (W)	$\lambda$	Dissipation coefficient of cooling rate
$\dot{Q}_H$	heat input load from heat source to generator (W)	$\sigma$	Entropy generation rate (W/K)
$\dot{Q}_O$	$\dot{Q}_C + \dot{Q}_A$	$\varepsilon_l$	coefficient of performance for three-heat-source refrigerator affected only by internal irreversibility
$R$	cooling load (W)	$\varepsilon_r$	coefficient of performance for reversible three-heat-source refrigerator
$R_m$	cooling rate at maximum coefficient of performance (W)	$\varepsilon_m$	coefficient of performance at maximum cooling rate
$r$	specific cooling load ( $\text{W m}^{-2}$ )	<i>Subscripts</i>	
$r_m$	specific cooling rate at maximum coefficient of performance ( $\text{W m}^{-2}$ )	max	maximum
$T_1$	temperature of working fluid in generator (K)		
$T_2$	temperature of working fluid in evaporator (K)		
$T_3$	temperature of working fluid in absorber and condenser (K)		
$T_3$	temperature of working fluid in absorber and condenser (K)		

However, since the absorption refrigeration cycles are in direct contact with reservoirs and sink, the heat transfers during the isothermal processes are supposed to be carried out infinitely slowly. Therefore, duration of the processes will be infinitely long and hence it is not possible to obtain a certain amount of cooling load  $\dot{Q}_L$  with heat exchangers having finite heat-transfer areas, i.e.  $\dot{Q} = 0$  for  $0 < A < \infty$ . If we require certain amount of cooling load in an absorption refrigerator executing a reversible cycle, the necessary heat exchanger area would be infinitely large, i.e.  $A \rightarrow \infty$  for  $\dot{Q} > 0$ .

Thus in classical thermodynamics the real absorption refrigerators producing cold with a certain amount of cooling load are compared with the ideal absorption refrigerators developing no cooling load. In other words the performance of an absorption refrigerator of given size (in term of total heat-transfer area) is measured with an ideal absorption refrigerator which would require an infinite total heat-transfer area to produce the same amounts of cooling load. In practice, all absorption refrigeration processes take place in finite-size devices in finite-time; therefore, it is impossible to meet reversibility conditions between the absorption refrigeration system and the surroundings. For this reason, the reversible absorption cycle cannot be considered as a comparison standard for practical absorption refrigeration systems from the view of cooling load on size perspective, although it gives an upper bound for coefficient of performance. The performance bound of classical thermodynamics [3–6] is highly important in theory, but it is usually too rough to predict the coefficient of performance of practical absorption refrigerators. Therefore, it is necessary to establish the bound of finite-time thermodynamics [7].

The finite-time thermodynamics has been first proposed by Chambadal [8] and Novikov [9] independently on 1957, then popularized in many works including Curzon and Ahlborn [10], De Vos [11], Sieniutycz et al. [12], Bejan [13–18], Wu [19], Chen [20], Stitou [21,22], Feidt [23,24], Leff and Teeters [25], Blanchard [26], Stitou and Feidt [27], Andresen [28], Sieniutycz and Salamon [29], De Vos [30], Bejan et al. [31], Bejan and Mamut [32], Berry et al. [33], Radcenco [34] and in many review articles including Sieniutycz and Shiner [35], Chen et al. [36], Hoffmann et al. [37] and Durmayaz et al. [38].

The finite-time thermodynamics tends to model the real systems in a way closer to reality and enable to distinguish the irreversibilities due to internal dissipation of the working fluid from those due to finite-rate heat transfer between the system and the external heat reservoirs and heat sink.

The objective of this paper is to review the present state of optimization of absorption refrigeration processes based on finite-time thermodynamics. The different performance optimization criteria are provided and discussed.

## 2. Optimization based on the coefficient of performance and cooling load criteria

### 2.1. Three-heat-source absorption refrigerator

An absorption refrigeration system (equivalent to three-heat-reservoir refrigeration system) affected by the irreversibility of finite rate heat transfer may be modeled as a combined cycle which consists of an endoreversible heat engine and an endoreversible

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