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## Exciton spectra and energy band structure of CuAlS<sub>2</sub> crystals

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### ARTICLE INFO

#### Article history:

Received 21 March 2009

Received in revised form

27 February 2010

Accepted 22 April 2010

#### Keywords:

Chalcopyrites

Optical reflectivity

Photoluminescence

Exciton polaritons

### ABSTRACT

Three exciton series are investigated in the reflectivity spectra of CuAlS<sub>2</sub> crystals at the temperature of 10 K. The  $n^A=1$  ( $\omega_T=3.543$  eV,  $\omega_L=3.546$  eV) and  $n^A=2$  (3.565 eV) lines of the  $\Gamma_4$  (A-series) excitons are observed in the  $E\parallel c$  polarization. In the  $E\perp c$  polarization, the  $\Gamma_5$  excitons ( $n^B=1$  at 3.668 eV and  $n^B=2$  at 3.686 eV) of the B-series, and  $n^C=1$  at 3.813 eV of the C-series are observed. The parameters of the excitons and the more exact values of the  $\Gamma_7 - \Gamma_6$ ,  $\Gamma_6 - \Gamma_6$ , and  $\Gamma_7 - \Gamma_6$  energy gaps are determined. The crystal field and spin-orbit splitting of the valence band is calculated. The electron ( $m_{c1}^*$ ) and hole ( $m_{v1}^*$ ,  $m_{v2}^*$ , and  $m_{v3}^*$ ) effective masses have been estimated.

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### 1. Introduction

The CuAlS<sub>2</sub> compound belonging to the I-III-VI<sub>2</sub> materials crystallizes in the chalcopyrite structure with the  $I\bar{4}2d - D_{2d}^{12}$  space group and it is the most wide bandgap material from this group. The I-III-VI<sub>2</sub> crystals are of interest from the point of view of their implementation in optoelectronic devices [1,2]. Solar cells are successfully developed on the basis of materials from this group [3–9]. Thin film device structures were obtained on some of these materials. The photoelectrical properties of surface barrier structures have been investigated [4,5]. The energy band structure of I-III-VI<sub>2</sub> chalcopyrites is calculated as for nearest analogs of binary compounds with zincblende structure (ZnSe, ZnS) [10–12]. All the crystals from the I-III-VI<sub>2</sub> group are characterized by strong anisotropy of optical properties in the region of fundamental absorption [13–29]. Many of chalcopyrite compounds are obtained by molecular beam epitaxy in the form of epitaxial layers, including the CuAlS<sub>2</sub> compound [13]. The CuGaS<sub>2</sub>–CuAlS<sub>2</sub> structures are interesting for the development of new optical devices for the green and blue spectral ranges. There is a relatively small mismatch of crystal lattices (0.4%) and a relatively big difference in the bandgaps of CuGaS<sub>2</sub> and CuAlS<sub>2</sub> crystals in this system.

In this paper we investigate the optical spectra of CuAlS<sub>2</sub> crystals in the exciton region. The energy positions of  $n=1$  and  $n=2$  lines of the three A-, B-, and C-exciton series are determined.

The exciton binding energy and the more precise values of the  $\Gamma_7 - \Gamma_6$ ,  $\Gamma_6 - \Gamma_6$ , and  $\Gamma_7 - \Gamma_6$  energy gaps are determined. The crystal field and spin-orbit splitting of the valence band are calculated. The electron ( $m_{c1}^*$ ) and hole ( $m_{v1}^*$ ,  $m_{v2}^*$  and  $m_{v3}^*$ ) effective masses are estimated. The contours of the exciton reflectivity spectra are calculated according to the one-oscillator model of the dispersion relations.

### 2. Experimental

CuAlS<sub>2</sub> crystals in the form of plates with  $2.5 \times 8$  mm<sup>2</sup> mirror-like surfaces and thickness of 1–2 mm were grown by chemical vapor transport. The surfaces of some plates are parallel to the C-axis. The orientation of crystals was determined according to X-ray diffraction data. The optical reflectivity and luminescence spectra were measured with a double SDL-1 spectrometer. For the low temperature measurements the samples were mounted on the cold station of an LTS-22 C330 optical cryogenic system. The luminescence spectra were excited by the 325 nm line of a Melles Griot He–Cd laser.

### 3. Results and discussions

According to theoretical calculations of energy band structure [10–12], the minimum inter-band gap is formed by direct electron transition in the center of the Brillouin zone. The lower conduction band is of  $\Gamma_6$  symmetry, while the upper  $V_1$ ,  $V_2$ , and  $V_3$  valence bands are of  $\Gamma_7$ ,  $\Gamma_6$ , and  $\Gamma_7$  symmetries, respectively.

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From group theoretical arguments and direct product of the group representations, the interaction of a hole from the  $\Gamma_7$  band with an electron from the  $\Gamma_6$  band leads to the formation of intrinsic excitons with ground states of  $\Gamma_3, \Gamma_4$ , and  $\Gamma_5$  symmetries ( $\Gamma_6 \times \Gamma_7 = \Gamma_3 + \Gamma_4 + \Gamma_5$ ). The  $\Gamma_4$  exciton transition is allowed in  $E\parallel c$  polarization, the  $\Gamma_5$  exciton transition is allowed in  $E\perp c$  polarization, whereas the  $\Gamma_3$  exciton transition is forbidden in both polarizations. The interaction of a hole from the  $\Gamma_6$  band with an electron from the  $\Gamma_6$  band leads to the formation of three exciton series with  $\Gamma_1, \Gamma_2$  and  $\Gamma_5$  symmetries. The  $\Gamma_5$  excitons are allowed, while  $\Gamma_1$  and  $\Gamma_2$  excitons are forbidden in  $E\perp c$  polarization according to the selection rules.

The  $n^A=1$  ( $\omega_T=3.543$  eV,  $\omega_L=3.546$  eV) and  $n^A=2$  (3.565 eV) lines of the  $\Gamma_4$  exciton hydrogen-like series are observed in the reflectivity spectra of CuAlS<sub>2</sub> crystals measured at 10 K in the  $E\parallel c$  polarization (Fig. 1). The reflectivity spectra in the region of the  $n^A=1$  line are of a usual excitonic shape with a maximum at 3.543 eV and a minimum at 3.546 eV. These peculiarities are due to the presence of the transversal and longitudinal excitons. A longitudinal-transversal exciton splitting of 3 meV is estimated for the  $\Gamma_4$  excitons from these data. A Rydberg constant of 32 meV is determined for the  $\Gamma_4$  exciton series from the position of  $n^A=1$  and  $n^A=2$  lines (Fig. 1). The energy of the continuum ( $E_g, n = \infty$ ) is equal to 3.575 eV. These energy values of the ground ( $n^A=1$ ) and excited ( $n^A=2$ ) states of excitons are in accordance with previously reported values of 3.534 and 3.665 eV, measured at 77 K [14–16].

The background dielectric constant has been estimated from the measurements reflectivity in the IR ( $400\text{ cm}^{-1}$ ) and near-IR ( $12,000\text{ cm}^{-1}$ ) regions.

$$\varepsilon_b = \left| \frac{1 + \sqrt{R}}{1 - \sqrt{R}} \right|^2$$

The reported value of  $\varepsilon_b$  in CuAlS<sub>2</sub> crystals equals 7.05 in the ( $E\parallel c$ ) polarization and 8.14 in the ( $E\perp c$ ) polarization far from the exciton resonances ( $\nu=3000\text{--}4000\text{ cm}^{-1}$ ) [26]. According to our data, the reflection coefficient at quantum energies lower than the energy of exciton resonances is  $R=0.21$  for ( $E\parallel c$ ) polarization. The calculated value of  $\varepsilon_b(E\parallel c)$  is 7.26. With this value of  $\varepsilon_b$  and Rydberg constant  $R=0.032$  eV, the  $\Gamma_4$ -excitons reduced mass equals to  $\mu = (\varepsilon_b^2 R / R_H) = 0.11 m_0$ , where  $R_H$  is the Rydberg energy of the hydrogen atom (13.6 eV). The Bohr radius ( $a_B$ ) of the S-state of the  $\Gamma_4$ -exciton equals  $0.3 \times 10^{-6}$  cm.

A maximum at 3.668 eV (transversal exciton) and a minimum at 3.670 eV (longitudinal exciton) are observed in the  $E\perp c$  polarization for the  $n^B=1$   $\Gamma_5$  exciton series (Fig. 1). The longitudinal-transversal splitting of the  $\Gamma_5$  exciton equals

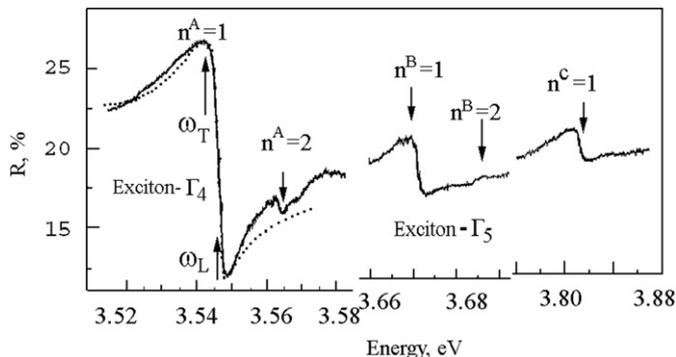


Fig. 1. Optical reflectivity spectra of CuAlS<sub>2</sub> crystals. The dot-line is the result of calculation of the reflectivity contour of the  $n=1$  line of the  $\Gamma_4$  exciton taking into account the spatial dispersion, i.e. limited value of the exciton mass  $M$ , and the boundary conditions.

2.0 meV. The  $n^B=2$  excited exciton state is observed at 3.687 eV. The binding energy of the  $\Gamma_5$  exciton equals 25 meV, and the energy of the continuum equals 3.693 eV. The C-exciton is observed at 3.813 eV ( $n^C=1$ ) in the same polarization. A band at 4.39 eV was observed at 77 K in the energy interval of 3.6–4.98 eV. Since no other lines were observed between the 3.665 and 4.39 eV at 77 K, the 4.39 eV band has been associated with C-exciton series [16]. As mentioned above, a line at 3.813 eV is observed in the reflectivity spectrum (Fig. 1). Taking into account these data and the exciton spectra of CuGaS<sub>2</sub> crystals [23–29], we assume that the 3.813 eV line is related to the ground state of the C-exciton.

For the B-exciton series, the reflection coefficient measured in the ( $E\perp c$ ) polarization equals 21% at 3.6 eV, and the dielectric constant  $\varepsilon_b(E\perp c)=7.2$ . The calculated  $\Gamma_5$ -exciton reduced mass equals  $0.09 m_0$  with the exciton binding energy of 25 meV. One should note that the values of  $\varepsilon_b(E\parallel c)$  and  $\varepsilon_b(E\perp c)$  are very close to each other in the region of exciton resonances, in contrast to a bigger difference observed far from the exciton resonances ( $\nu=3000\text{--}4000\text{ cm}^{-1}$ ) [26]. This is because the exciton resonances are close to the optical isotropic wavelength [30].

A weak luminescence is observed at 200 K in CuAlS<sub>2</sub> crystals under the excitation by the 325 nm line of a He–Cd laser. The luminescence intensifies with decreasing temperature to 10 K, and further to 10 K (Fig. 2). A broad photoluminescence band is observed at 10 K around 3 eV, along with two narrow bands around 3.3 eV and a band at 3.54 eV. The band at 3.54 eV is assumed to be related to the recombination of  $\Gamma_4$  exciton polaritons, while the narrow lines are due to the recombination of donor–acceptor pairs (DAP) (Fig. 2).

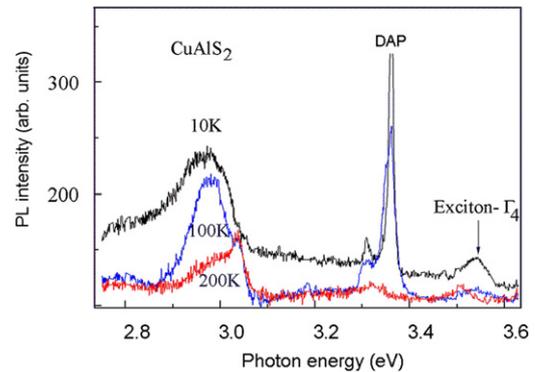


Fig. 2. The luminescence spectra of CuAlS<sub>2</sub> crystals excited by the 325 nm line of a He–Cd laser measured at different temperatures.

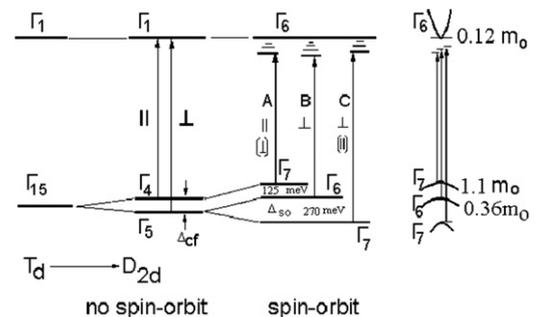


Fig. 3. Band structure at the  $\Gamma$  point showing the transition from zincblende ( $T_d$ ) to chalcopyrite ( $D_{2d}$ ) structure for the CuAlS<sub>2</sub>.

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