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Electronic structure and magnetic state of InCNi₃

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ABSTRACT

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1. Introduction

The discovery of superconductivity at ~8 K in intermetallic compound MgCNi₃ [1] has stimulated considerable attention on the cubic antiperovskites $XCNi_3$ (X = Mg, Zn, Al, Ga, and In). However no superconductivity was found in ZnCNi₃ [2], GaCNi₃ [3], AlCNi₃ [4,5], and InCNi₃ [6]. For AlCNi₃ and CaCNi₃, experimental studies reported that AlCNi3 is a strongly exchange-enhanced Pauli paramagnet in the very vicinity of ferromagnetic order [5] and its transition from a paramagnetic state into weak ferromagnetic one takes place at temperature near 300 K [4], while CaCNi₃ has a large electron–electron correlation due to the proximity of ferromagnetic order from the side of exchange-enhanced Pauli paramagnet [3]. But these magnetic properties of AlCNi3 and CaCNi3 reported in experiments are contradictory to the results predicted from density functional theory (DFT) calculations [7–9] that did not support AlCNi₃ being in a magnetic ground state and CaCNi₃ being in the presence of strong electron-electron correlations. DFT calculations [8] instead suggested that the tendencies toward magnetism found in experiments for AlCNi₃ and CaCNi₃ should be explained by the deviation of the Ni/C atomic ratio from the ideal stoichiometry. For the same group compound InCNi₃, its non-stoichiometry

Using the projector augmented wave method within the local density approximation and the generalized gradient approximation, we have studied the electronic structures and magnetic state of cubic antiperovskite InCNi₃. It is found that Ni 3d states play dominant roles near the Fermi level of this compound. The analysis of bonding nature in InCNi₃ implies that Ni 3d and C 2p states exhibit strong hybridization. Both the spin-polarized calculations and the fixed-spin-moment calculations indicate that cubic antiperovskite InCNi₃ has a stable paramagnetic (non-magnetic) state.

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compound In_{0.95}CNi₃ has been synthesized recently [6]. From the measured results of electronic specific coefficient and Debye temperature, In_{0.95}CNi₃ was proposed to behave as a ferromagnetic metal below the Curie temperature (577 K) and the emergence of ferromagnetism in In_{0.95}CNi₃ is ascribed to the deviation of the Ni/In atomic ratio from the ideal stoichiometry [6], which is different to the findings from the DFT calculations [7,8] on AlCNi3 and CaCNi3. To shed more light on the understanding on the presence of ferromagnetism in In_{0.95}CNi₃, it is of great interest to study the physical properties of the ideal stoichiometry compound InCNi₃. Recently, InCNi₃ is proposed to be a paramagnetic nonsuperconductor from the calculated total energy, density of states at Fermi level, and electron-phonon coupling constant [9], however the detailed electronic structures of InCNi₃ is not presented there, especially for the bonding nature. In addition, only the generalized gradient approximation is employed in previous DFT calculations on InCNi₃ [9]. In order to completely understand the electronic and magnetic state of cubic antiperovskite InCNi₃, we conduct a comprehensive ab initio calculations on this compound by using the projector augmented wave method with plane-wave basis set within the local density approximation and the generalized gradient approximation. In this study, we predict the lattice parameter and bulk modulus of InCNi₃, discuss the bonding nature in this compound from the calculated electronic density of states and charge distribution, and also perform the fixed-spin-moment calculations to examine the stability of possible magnetic state of InCNi₃.



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2. Method of calculations

All calculations on InCNi₃ are performed by means of the projector augmented wave method [10] within the local density approximation (LDA) [11,12] and the generalized gradient approximation (GGA) [13], as implemented in VASP code [14,15]. In LDA calculations, Perdew and Zunger's parametrization [12] of Ceperley and Alder quantum Monte Carlo results [11] is used for the correlation functional. While Perdew-Wang'91 (PW91) exchange and correlation functional [13] is adopted in GGA calculations. A nickel potential with 4s and 3d electrons as valence states, an indium potential with 4d, 5s, and 5p electrons as valence states, and a carbon potential with 2s and 2p electrons as valence states have been used. Wave functions are expanded by the plane waves up to an energy cutoff of 450 eV. Brillouin-zone (BZ) integrations during the self-consistent calculations are approximated using the special k-point sampling of Monhkorst-Pack scheme [16] with a $13\times13\times13$ grid. The convergence of the total energy with respect to both k-point sampling and plane-wave energy cut-off are carefully examined, and the above calculation setup is found to provide good precision in present study. The total energies are obtained as a function of volume and then the data of total energies versus volumes are fitted to 3rd-order Birch-Murnaghan equation of state [17] to give the equilibrium lattice constant and other ground state properties of InCNi₃. In the calculation of density of states (DOS), a dense k-grid of $20 \times 20 \times 20$ is used and the angular-momentum-projected DOS is integrated within the

Table 1

Calculated lattice constants (a, in Å), bulk modulus (B, in GPa), and the first pressure derivative of B (B') for InCNi₃, with a comparison with available experimental values and other theoretical results

Property	Present		Previous [9]	Experimental [6]
	LDA	GGA-PW91	GGA-PBE	
а	3.7893	3.8805	-	3.7836
В	228.7	185.6	183.56	-
Β'	4.70	4.74	-	-

spheres of radii 1.677 Å for In atom, 0.863 Å for C atom, and 1.286 Å for Ni atom, respectively.

3. Results and discussions

Recent experimental study [6] on In_{0.9}CNi₃ reported that this compound has the typical cubic antiperovskite structure with the space group $221(Pm\bar{3}m)$ and its lattice parameter is about 3.7836 Å determined from the powder X-ray diffraction patterns. The calculated equilibrium lattice constant (*a*), bulk modulus (*B*), and first pressure derivative of bulk modulus (B') of cubic antiperovskite InCNi₃ are listed in Table 1. From the results in Table 1, one can see that the deviation of the predicted lattice constant of InCNi₃ in LDA calculation is about 0.15% and the one in GGA-PW91 calculation is about 2.5% with respect to the experimental lattice constant of In_{0.9}CNi₃, respectively. The bulk modulus of InCNi₃ is predicted to be 228.7 GPa in LDA calculation and 185.6 GPa in GGA-PW91 calculation, respectively, and especially the latter is in excellent agreement with the previous GGA-PBE calculation result (183.56 GPa) [9] from highly precise full-potential linearized augmented plane-wave (FP-LAPW) method.

In order to understand the electronic nature of this compound, its band structure and density of states (DOS) have been calculated. For InCNi3 at equilibrium, the LDA and GGA-PW91 calculations give almost identical electronic band structure around the Fermi energy. Here we show only GGA-PW91 results. Fig. 1(a) shows the calculated electronic band structure of cubic InCNi₃ along several high-symmetry lines in the first BZ. At the same time, the total DOS of cubic InCNi₃ is shown in Fig. 1(b). One band around -12 eV in Fig. 1(a) is mainly contributed by the bonding states of C 2s orbital. One band around -8 eV mostly consists of In 5s orbital and Ni 4s orbital. Complex bands extending between -7 and -4 eV are composed of C 2p and Ni 3d orbital mixed lightly with In 5s orbital. Dense bands around -2 eV are dominated by Ni 3d electrons and In 5p orbital, while the latter one contributes a little to the lower part of these dense bands. Two conduction bands crossing the Fermi energy $(E_{\rm F})$ are dominated by Ni 3d, C 2p and In 5p orbitals. The contribution of



Fig. 1. Electronic band structure and total density of states (TDOS) of cubic InCNi₃ at equilibrium, obtained from GGA calculations.

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