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Monte Carlo and least-squares methods applied in unfolding of X-ray spectra measured with cadmium telluride detectors

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Abstract

Spectra of calibration sources and X-ray beams were measured with a cadmium telluride (CdTe) detector. The response function of the detector was simulated using the GEANT4 Monte Carlo toolkit. Trapping of charge carriers were taken into account using the Hecht equation in the active zone of the CdTe crystal associated with a continuous function to produce drop of charge collection efficiency near the metallic contacts and borders. The rise time discrimination is approximated by a cut in the depth of the interaction relative to cathode and corrections that depend on the pulse amplitude. The least-squares method with truncation was employed to unfold X-ray spectra typically used in medical diagnostics and the results were compared with reference data.

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1. Introduction

The employment of cadmium telluride (CdTe) detectors is increasing in applications where portability is required. Compactness and portability are very attractive properties for using CdTe in characterization of X-ray beams commonly employed in medical applications, for example.

Characteristics of X-ray beams are usually obtained by unfolding the measured spectra, which requires the knowledge of the detector's response function. If the response function is well known, then the main source of errors in the reconstruction process would be the statistical errors of the measured spectrum. The stripping method [1,2] is commonly employed to reconstruct the original spectrum, although the propagation of errors for this procedure is not found in the literature. In this work, the least-squares method was employed in the unfolding of spectra measured with a CdTe detector and in the calculation of the associated covariance matrices of the reconstructed spectra. The response of the CdTe detector was simulated with the GEANT4 Monte Carlo toolkit [3], including a model for the charge collection efficiency. Simulated results were compared with experimental data obtained for standard calibration sources and X-ray beams.

2. The problem of spectrum unfolding

The detected spectrum, represented by a vector $\mathbf{d} = d_1, d_2, \ldots, d_n$, is related to the discrete spectrum of the sources of radiation $\mathbf{s} = s_1, s_2, \ldots, s_m$ through the expression

$$\mathbf{d} = \mathbf{R}\mathbf{s} \tag{1}$$

where $\mathbf{R}$ is a matrix that represents the response function of the instrument [4]. In the commonly employed method to obtain $\mathbf{s}$, called stripping method [1], the original spectrum is reconstructed from the largest to the smallest energies by applying the algorithm

$$s_i = \left[d_i - \sum_{j=i+1}^{m} R_{i,j} s_j \right] R_{i,i}^{-1}. \tag{2}$$

The propagation of experimental errors using this procedure is not found in the literature.

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Instead of the stripping method, in this study the least squares method was employed to obtain an estimate $\mathbf{s}'$ of the original spectrum

$$\mathbf{s}' = (\mathbf{R}^t\mathbf{V}_\mathbf{d}^{-1}\mathbf{R})^{-1}\mathbf{R}^t\mathbf{V}_\mathbf{d}^{-1}\mathbf{d} \quad (3)$$

and its covariance matrix, $\mathbf{V}_{\mathbf{s}'}$, given by

$$\mathbf{V}_{\mathbf{s}'} = (\mathbf{R}^t\mathbf{V}_\mathbf{d}^{-1}\mathbf{R})^{-1} \quad (4)$$

where the covariance matrix of $\mathbf{d}$, $\mathbf{V}_\mathbf{d}$, is diagonal with elements $\mathbf{V}_{\mathbf{d}i,i} = \sigma_i^2 = d_i$.

It is known from the theory of inverse problems that the solution of Eq. (3) presents huge oscillations and, in several cases, the inversion procedure is impossible due to singular matrix. For this class of ill-posed problem, regularization or truncation ($m<n$) methods [4,5] are applied to obtain a solution and eventually attenuate the oscillations.

The matrix $\mathbf{V}_{\mathbf{s}'}$ is necessary for the propagation of errors in any operation that involves $\mathbf{s}'$, as for example in the calculation of the mean energy of the unfolded spectrum. Defining the vector $\mathbf{e} = e_1, e_2, \ldots, e_m$, where e_i corresponds to the energy of the channel i, the mean energy is calculated by

$$\langle E\rangle = \mathbf{e}^t\mathbf{s}'A^{-1} \quad (5)$$

and its variance is obtained through

$$\sigma^2_{\langle E\rangle} = \mathbf{V}_{\langle E\rangle} = \mathbf{e}^t\mathbf{V}_{\mathbf{s}'}\mathbf{e}A^{-2} \quad (6)$$

where A is the sum of all elements of $\mathbf{s}'$.

3. The detector's response

The CdTe detector used in this study was made by Amptek (model: XR-100T-CdTe). This detector has a crystal with area of 3 mm^2 and thickness of 1 mm, operates at bias of 400 V, and has beryllium window with thickness of 250 μm. Additional information relevant for the detector modeling were taken from Amptek's documentation [6].

CdTe detectors exhibit spectral distortions mainly because of hole trapping, whose main effect is the presence of tails at the low-energy side of the full absorption peaks. The data acquisition system of XR-100T-CdTe includes the option of rise time discrimination (RTD). When operating with RTD, the slow rise time pulses are rejected and the low-energy tails are suppressed [7].

The detector's response function was calculated using the GEANT4 Monte Carlo toolkit [4] version 8.0.p01. For each interaction inside the CdTe crystal, two quantities were extracted: the deposited energy and the interaction position. Using this information, the charge collection efficiency can be calculated with the Hecht equation [8]

$$\eta(x) = [\lambda_h(1 - e^{-x/\lambda_h}) + \lambda_e(1 - e^{-(D-x)/\lambda_e})]D^{-1} \quad (7)$$

where D is the crystal thickness, x is the distance of the interaction position to the cathode, λ_h and λ_e are the drift length for holes and electrons, whose values (0.36 and 24 cm, respectively) were determined from measured spectra of ^{133}Ba, ^{241}Am and ^{152}Eu without RTD.

Fig. 1 shows ^{133}Ba spectra measured with and without collimation and RTD. The effect of RTD is evident from the tail suppression at the low-energy side of the full-energy peaks. After including a collimator, one observes suppression of background and lowering of a step-shaped structure at the low-energy peaks. Simulations including the collimator showed that this structure has large contribution from interactions that occur at the borders of the CdTe crystal. Comparison between simulated and experimental data, showed that the intensities of the low-energy photopeaks could be reproduced only after the inclusion of a dead layer near the cathode.

The best agreement between simulated and experimental data was achieved by calculating the charge collection efficiency with Eq. (7) for almost the whole crystal, but applying a continuous cut to represent dead layers in the regions near the cathode and the borders of the CdTe crystal. Complementary error functions were combined with the Hecht equation to produce continuous drop of the charge collection efficiency at the periphery of the detector: $f_i(t) = 0.5\{\mathrm{erfc}[(t-t_{ci})\alpha_i/l_i]\}$, where t_{ci} is the middle point of the dead layer corresponding to the coordinate t, l_i is the thickness of the dead layer and α_i is a parameter associated with the inclination of the curve at the cut position. Thicknesses of dead layers were 15 and 20 μm for cathode and borders, respectively. Electronic noise (120 eV) and statistical fluctuations (Fano factor = 0.1) were included as a Gaussian dispersion in the detector's response function.

Since the holes are slower than electrons, the principal effect of RTD is the rejection of slow pulses generated beyond a certain distance from the cathode [6]. Comparison between simulated and experimental spectra of ^{241}Am, ^{133}Ba and ^{152}Eu revealed that rejecting events generated at depths beyond 140 μm from the cathode is adequate to describe the efficiency up to 60 keV. The non-linear behavior of the employed RTD method, which shows dependence on the pulse amplitude, was taken into account

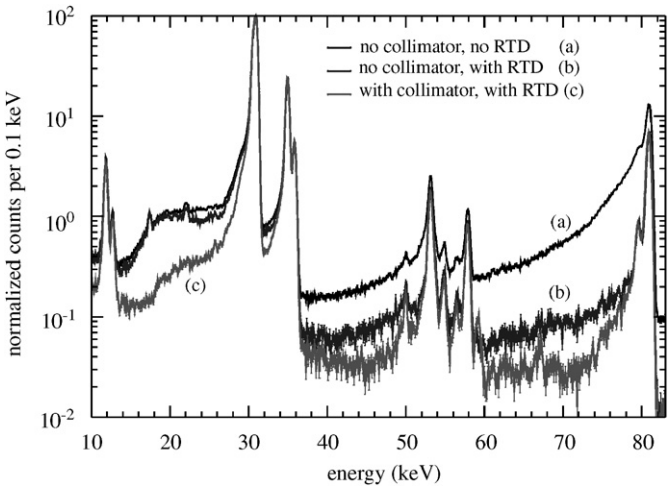

Fig. 1. ^{133}Ba spectra measured with a CdTe detector using different configurations.