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Measurement of the lifetime of tau-lepton

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Abstract

The lifetime of the τ -lepton is measured using the process $e^+e^- \rightarrow \tau^+\tau^-$, where both τ -leptons decay to $3\pi\nu_{\tau}$. The result for the mean lifetime, based on 711 fb⁻¹ of data collected with the Belle detector at the $\Upsilon(4S)$ resonance and 60 MeV below, is $\tau = (290.17 \pm 0.53(\text{stat.}) \pm 0.33(\text{syst.})) \cdot 10^{-15}$ s. The first measurement of the lifetime difference between τ^+ and τ^- is performed. The upper limit on the relative lifetime difference between positive and negative τ -leptons is $|\Delta \tau|/\tau < 7.0 \times 10^{-3}$ at 90% CL.

High precision measurements of the mass, lifetime and leptonic branching fractions of the τ -lepton can be used to test lepton universality [1, 2], which is assumed in the Standard Model. Among the recent experimental results that may manifest the violation of the lepton universality in the case of the τ -lepton, the combined measurement of the ratio of the branching fraction of W-boson decay to τv_{τ} to the mean branching fraction of W-boson decay to μv_{μ} and ev_{e} by the four LEP experiments stands out: $2\mathcal{B}(W \to \tau \nu_{\tau})/(\mathcal{B}(W \to \tau \nu_{\tau}))$ μv_{μ} + $\mathcal{B}(W \rightarrow e v_e)$ = 1.066 ± 0.025 [3], which differs from unity by 2.6 standard deviations. The present PDG value of the τ -lepton lifetime $(290.6 \pm 1.0) \cdot 10^{-15}$ s [4] is dominated by the results obtained in the LEP experiments [5, 6, 7, 8]. The results of the analysis described below are recently published in [9]. 1

In the following, we use symbols with and without an asterisk for quantities in the e^+e^- center-of-mass (CM) and laboratory frame, respectively. In the CM frame, τ^+ and τ^- leptons emerge back to back with the energy E^*_{τ} equal to the beam energy E^*_{beam} if we neglect the initial-(ISR) and final-state radiation (FSR). We determine the direction of the τ -lepton momentum in the CM frame as follows. If the neutrino mass is assumed to be zero for the hadronic decay $\tau \rightarrow Xv_{\tau}$ (X representing the

hadronic system with mass m_X and energy E_X^*), the angle θ^* between the momentum \vec{P}_X^* of the hadronic system and that of the τ -lepton is given by:

$$\cos\theta^* = \frac{2E_\tau^* E_X^* - m_\tau^2 - m_X^2}{2P_X^* \sqrt{(E_\tau^*)^2 - m_\tau^2}}.$$
 (1)

The requirement that the τ -leptons be back to back in the CM can be written as a system of three equations: two linear and one quadratic. For the components x^* , y^* , z^* of the unit vector \hat{n}^*_+ representing the direction of the positive τ -lepton, we write:

$$\begin{cases} x^* \cdot P_{1x}^* + y^* \cdot P_{1y}^* + z^* \cdot P_{1z}^* = |P_1^*| \cos \theta_1^* \\ x^* \cdot P_{2x}^* + y^* \cdot P_{2y}^* + z^* \cdot P_{2z}^* = -|P_2^*| \cos \theta_2^* \\ (x^*)^2 + (y^*)^2 + (z^*)^2 = 1 \end{cases}$$
(2)

where \vec{P}_1^* and \vec{P}_2^* are the momenta of the hadronic systems in the CM and $\cos \theta_i^*$ (i = 1, 2) are given by Eq. (1). Index 1 (2) is used for the positive (negative) τ -lepton.

There are two solutions for Eq. (2), so the direction \hat{n}^*_+ is determined with twofold ambiguity. In the present analysis, we take the mean vector of the two solutions of Eq. (2) as the direction of the τ -lepton in CM. The analysis of MC simulated events shows that there is no bias due to this choice.

Each direction \hat{n}_{\pm}^* is converted to a four-momentum using the e^{\pm} beam energy and the τ mass. Both fourmomenta are then boosted into the laboratory frame,

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Figure 1: The schematic view of the $\tau^+\tau^-$ event in the laboratory frame.

each passing through the corresponding τ decay vertex \vec{V}_i that is determined by the three pion-daughter tracks (see Fig. 1). We approximate the trajectory of τ leptons in the magnetic field of the Belle detector with a straight line. Due to the finite detector resolution, these straight lines do not intersect at the $\tau^+\tau^-$ production point. The three-dimensional separation between these lines is characterized by the distance dl between the two points $(\vec{V}_{01} \text{ and } \vec{V}_{02})$ of closest approach. The typical size of dl is ~ 0.01 cm. For the production point of each τ -lepton, we take the points \vec{V}_{01} and \vec{V}_{02} . The flight distance l_1 (l_2) of the τ^+ (τ^-) in the laboratory frame is defined as the distance between the points \vec{V}_1 and \vec{V}_{01} (\vec{V}_2 and \vec{V}_{02}). The proper time t (the product of the speed of light and the decay time of τ -lepton) for the positive τ -lepton is equal to the distance l_1 divided by its relativistic kinematic factor $\beta\gamma$ in the laboratory frame: $t_1 = l_1/(\beta \gamma)_1$. The corresponding parameter for the negative τ -lepton is $t_2 = l_2/(\beta \gamma)_2$.

The analysis presented here is based on the data collected with the Belle detector [10, 11] at the KEKB asymmetric-energy e^+e^- (3.5 on 8 GeV) collider [12, 13] operating at the $\Upsilon(4S)$ resonance and 60 MeV below. The total integrated luminosity of the data sample used in the analysis is 711 fb⁻¹. Two inner detector configurations were used. A 2.0 cm beampipe and a 3-layer silicon vertex detector (SVD1) were used for the first sample of 157 fb⁻¹, while a 1.5 cm beampipe, a 4-layer silicon detector (SVD2) and a small-cell inner drift chamber were used to record the remaining 554 fb⁻¹ [14].

The following requirements are applied for the selection of the $\tau^+\tau^-$ events where both τ -leptons decay into three charged pions and a neutrino: there are exactly six charged pions with zero net charge and there are no other charged tracks; the thrust value of the event in the CM frame is greater than 0.9; the square of the transverse momentum of the 6π system is required to be greater than $0.25 \, (\text{GeV}/c)^2$ to suppress the $e^+e^- \rightarrow e^+e^-6\pi$ two-photon events: the mass $M(6\pi)$ of the 6π system should fulfill the requirement $4 \text{ GeV}/c^2 <$ $M(6\pi) < 10.25 \,\mathrm{GeV}/c^2$ to suppress other background events; there should be three pions (triplet) with net charge equal to ± 1 in each hemisphere (separated by the plane perpendicular to the thrust axis in the CM), the pseudomass (see the definition in Ref. [15]) of each triplet of pions must be less than $1.8 \,\text{GeV}/c^2$ and each pion-triplet vertex-fit quality must satisfy $\chi^2 < 20$; the discriminant D of Eq. (2) should satisfy D > -0.05(with slightly negative values arising from experimental uncertainties; if this happens, we use D = 0 when solving the equation); the distance of closest approach must satisfy dl < 0.03 cm to reject events with large uncertainties in the reconstructed momenta and vertex positions. All of these selection criteria are based on a study of the signal and background Monte Carlo (MC) simulated events. For the signal MC sample, we use $\tau^+\tau^-$ events produced by the KKMC generator [16] with the mean lifetime $\langle \tau \rangle = 87.11 \,\mu\text{m}$ that are then fed to the full detector simulation based on GEANT 3 [17]. These events are passed through the same reconstruction procedures as for the data. For the background estimation, we use the MC samples of events generated by the EVTGEN program [18], which correspond to the one-photon annihilation diagram $e^+e^- \rightarrow q\bar{q}$, where $q\bar{q}$ are $u\bar{u}$, $d\bar{d}$, $s\bar{s}$ (uds events), $c\bar{c}$ (charm events), and $e^+e^- \rightarrow \Upsilon(4S) \rightarrow B^+B^-, B^0\bar{B^0}$ (beauty events). All these events are passed through the full detector simulation and reconstruction procedures. The statistics in these MC samples are equivalent to the integrated luminosity of the data, *i.e.*, the number of events of a given category is equal to the product of the integrated luminosity of the data and the expected cross section from theory. For the estimation of the background from the process $\gamma \gamma \rightarrow hadrons$ ($\gamma \gamma$ events), we use events generated by PYTHIA [19] that are subjected to the aforementioned simulation and reconstruction procedures. In addition to the above MC samples, we also use two $e^+e^- \rightarrow \tau^+\tau^-$ MC samples, generated by KKMC, where both τ -leptons are forced to decay into three charged pions and a neutrino. The mean lifetimes for these two samples are 84 and 90 μ m, which are about 10 σ below and above the PDG value. These two samples are also passed through the same detector simulation and reconstruction procedures.

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