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Nuclear Physics A 957 (2017) 85-98



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## Analysis of the $D_{s3}^*(2860)$ as a D-wave $c\bar{s}$ meson with QCD sum rules

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#### Abstract

In this article, we assign the  $D_{s3}^*(2860)$  to be a D-wave  $c\bar{s}$  meson, and study the mass and decay constant of the  $D_{s3}^*(2860)$  with the QCD sum rules by calculating the contributions of the vacuum condensates up to dimension-6 in the operator product expansion. The predicted mass  $M_{D_{s3}^*} = (2.86 \pm 0.10)$  GeV is in excellent agreement with the experimental value  $M_{D_{s3}^*} = (2860.5 \pm 2.6 \pm 2.5 \pm 6.0)$  MeV from the LHCb collaboration. The present prediction supports assigning the  $D_{s3}^*(2860)$  to be the D-wave  $c\bar{s}$  meson. © 2016 Published by Elsevier B.V.

Keywords:  $D_{s3}^{*}(2860)$ ; QCD sum rules

### 1. Introduction

In 2006, the BaBar collaboration observed the  $D_{sJ}^*(2860)$  meson with the mass (2856.6 ± 1.5±5.0) MeV and the width (48±7±10) MeV in decays to the final states  $D^0K^+$  and  $D^+K_S^0$  using 240 fb<sup>-1</sup> of data recorded by the BaBar detector at the PEP-II asymmetric-energy  $e^+e^-$  storage rings at the Stanford Linear Accelerator Center [1]. In 2009, the BaBar collaboration confirmed the  $D_{sJ}^*(2860)$  in the  $D^*K$  channel using 470 fb<sup>-1</sup> of data recorded by the BaBar detector, and measured the ratio R among the branching fractions [2],

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http://dx.doi.org/10.1016/j.nuclphysa.2016.08.003 0375-9474/© 2016 Published by Elsevier B.V.

$$R = \frac{\operatorname{Br}\left(D_{sJ}^{*}(2860) \to D^{*}K\right)}{\operatorname{Br}\left(D_{sJ}^{*}(2860) \to DK\right)} = 1.10 \pm 0.15 \pm 0.19 .$$
(1)

The observation of the decays  $D_{sJ}^*(2860) \rightarrow D^*K$  rules out the  $J^P = 0^+$  assignment, the possible assignments are the 1<sup>3</sup>D<sub>3</sub>  $c\bar{s}$  meson [3–10], the  $c\bar{s}$ - $cn\bar{s}\bar{n}$  mixing state [11], the dynamically generated  $D_1(2420)K$  bound state [12], etc.

In 2014, the LHCb collaboration observed a structure at 2.86 GeV with significance of more than 10 $\sigma$  in the  $\overline{D}^0 K^-$  mass spectrum in the Dalitz plot analysis of the decays  $B_s^0 \to \overline{D}^0 K^- \pi^+$ using a data sample corresponding to an integrated luminosity of 3.0 fb<sup>-1</sup> of pp collision data recorded by the LHCb detector, the structure contains both spin-1 ( $D_{s1}^{*-}(2860)$ ) and spin-3 ( $D_{s3}^{*-}(2860)$ ) components, which can be assigned to be the  $J^P = 1^-$  and  $3^-$  members of the 1D family [13,14]. The measured masses and widths are  $M_{D_{s3}^*} = (2860.5 \pm 2.6 \pm 2.5 \pm 6.0)$  MeV,  $M_{D_{s1}^*} = (2859 \pm 12 \pm 6 \pm 23)$  MeV,  $\Gamma_{D_{s3}^*} = (53 \pm 7 \pm 4 \pm 6)$  MeV, and  $\Gamma_{D_{s1}^*} = (159 \pm 23 \pm 27 \pm 72)$  MeV, respectively. Furthermore, the LHCb collaboration obtained the conclusion that the  $D_{sJ}^*(2860)$  observed by the BaBar collaboration in the inclusive  $e^+e^- \to \overline{D}^0 K^- X$  production and by the LHCb collaboration in the  $pp \to \overline{D}^0 K^- X$  processes consists of at least two particles [2,15].

If we assign the  $D_{sJ}^*(2860)$  to be the  $1^3D_3$  state or the  $D_{s3}^*(2860)$ , the ratio R from the leading order heavy meson effective theory [3], the  ${}^3P_0$  model [4,9,16] and the relativized quark model [17] cannot reproduce the experimental value  $R = 1.10 \pm 0.15 \pm 0.19$  [2]. In Ref. [18], we assign the  $D_{s3}^*(2860)$  and  $D_{s1}^*(2860)$  to be the  $1^3D_3$  and  $1^3D_1$   $c\bar{s}$  states, respectively, study their strong decays with the heavy meson effective theory by including the chiral symmetry breaking corrections. We can reproduce the experimental value  $R = 1.10 \pm 0.15 \pm 0.19$  with suitable hadronic coupling constants if the chiral symmetry breaking corrections are large. The preferred assignment is  $D_{sJ}^*(2860) = D_{s3}^*(2860)$ , while the assignment  $D_{sJ}^*(2860) = D_{s1}^*(2860)$  is not excluded.

According to the predictions of the potential models [19], the masses of the 1D  $c\bar{s}$  states are about 2.9 GeV. It is reasonable to assign the  $D_{s1}^*(2860)$  and  $D_{s3}^*(2860)$  to be the  $1^3D_1$  and  $1^3D_3$  $c\bar{s}$  states, respectively. We can obtain further support by calculating the mass of the  $D_{s3}^*(2860)$ based on the QCD sum rules. The QCD sum rules is a powerful theoretical tool in studying the ground state hadrons and has given many successful descriptions of the masses, decay constants, form-factors and hadronic coupling constants, etc. [20,21]. There have been many works on the spin-parity  $J^P = 0^{\pm}$ ,  $1^{\pm}$  heavy-light mesons with the full QCD sum rules [22,23] (and references therein), while the works on the  $J^P = 2^+$  are few [24,25], the  $J^P = 3^-$  heavy-light mesons are only studied with the QCD sum rules combined with the heavy quark effective theory [26]. In this article, we assign the  $D_{s3}^*(2860)$  to be a D-wave  $c\bar{s}$  meson, study the mass and decay constant of the  $D_{s3}^*(2860)$  with the full QCD sum rules in details by calculating the contributions of the vacuum condensates up to dimension-6 in the operator product expansion.

The article is arranged as follows: we derive the QCD sum rules for the mass and decay constant of the  $D_{s3}^*(2860)$  in Sect. 2; in Sect. 3, we present the numerical results and discussions; and Sect. 4 is reserved for our conclusions.

#### 2. QCD sum rules for the $D_{s3}^*(2860)$ as a D-wave meson

In the following, we write down the two-point correlation function  $\Pi_{\mu\nu\rho\alpha\beta\sigma}(p)$  in the QCD sum rules,

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