

Available online at www.sciencedirect.com

ScienceDirect

PROCEEDINGS SUPPLEMENTS

Nuclear Physics B (Proc. Suppl.) 253-255 (2014) 103-106

www.elsevier.com/locate/npbps

Towards a determination of the tau lepton dipole moments

M. Fael^{a,b}, L. Mercolli^c, M. Passera^d

^aDipartimento di Fisica e Astronomia, Università di Padova, 35131 Padova, Italy ^bInstitut für Theoretische Physik, Universität Zürich, CH-8057, Zürich, Switzerland ^cPrinceton University, Department of Astrophysical Sciences, Princeton, NJ, 08544, USA ^dIstituto Nazionale Fisica Nucleare, 35131 Padova, Italy

Abstract

The τ anomalous magnetic moment a_{τ} and electric dipole moment d_{τ} have not yet been observed. The present bounds on their values are of order 10^{-2} and $10^{-17}e$ cm, respectively. We propose to measure a_{τ} with a precision of $O(10^{-3})$ or better and improve the existing limits on d_{τ} using precise $\tau^- \to l^- \nu_{\tau} \bar{\nu}_l \gamma$ (l = e or μ) data from high-luminosity B factories. A detailed feasibility study of this method is underway.

1. Introduction

The very short lifetime of the τ lepton $(2.9 \times 10^{-13} \text{s})$ makes it very difficult to measure its electric and magnetic dipole moments. The present resolution on its anomalous magnetic moment a_{τ} is only of $O(10^{-2})$ [1], more than an order of magnitude larger than its Standard Model (SM) prediction $\frac{\alpha}{2\pi} \simeq 0.001$. In fact, while the SM value of a_{τ} is known with a tiny uncertainty of 5×10^{-8} [2], this short lifetime has so far prevented the determination of a_{τ} measuring the τ spin precession in a magnetic field, like in the electron and muon g-2 experiments. Instead, experiments focused on various highprecision measurements of τ pair production in highenergy processes, comparing the measured cross sections with the SM predictions. As these processes involve off-shell photons or taus in the $\tau \bar{\tau} \gamma$ vertices, the measured quantity is not directly a_{τ} .

A precise measurement of a_{τ} would offer an excellent opportunity to unveil new physics effects. Indeed, in a large class of theories beyond the SM, new contributions to the anomalous magnetic moment of a lepton l of mass m_l scale with m_l^2 . Therefore, given the large factor $m_{\tau}^2/m_{\mu}^2 \sim 283$, the g-2 of the τ is much more sensitive than the muon one to electroweak and new physics loop effects that give contributions proportional to m_l^2 . In these scenarios, the present discrepancy in the muon g-2 suggests a new-physics effect in a_{τ} of

 $O(10^{-6})$; several theories exist where this naïve scaling is violated and much larger effects are expected [3].

Ever since the very first discovery of CP violation in the 1960s there have been ongoing efforts to measure fundamental electric dipole moments (EDMs). Indeed, EDMs violate parity and time reversal; if CPT is a good symmetry, T violation implies CP violation and vice versa. The SM predictions for lepton EDMs are far too small to be seen by projected future experiments. Hence, the observation of a non-vanishing lepton EDM would be evidence for a CP-violating new physics effect [4]. The experimental determination or a stringent bound on the τ lepton EDM d_{τ} poses the same difficulties that one encounters in the case of a_{τ} : the τ 's lifetime is very short. Nonetheless, a CP violation signature arising from d_{τ} was searched for in the $e^+e^- \to \tau^+\tau^-$ reaction reaching a sensitivity to d_{τ} of $10^{-17}e$ cm [5].

After a brief review of the present experimental and theoretical status of the τ dipole moments, we discuss the possibility to probe them via precise measurements of the decays $\tau^- \to l^- \nu_\tau \, \bar{\nu}_l \, \gamma \, (l = e \, \text{or} \, \mu)$.

2. The anomalous magnetic moment of the tau

The SM prediction for a_{τ} is given by the sum of QED, electroweak (EW) and hadronic terms. The QED contribution has been computed up to three loops: $a_{\tau}^{\text{QED}} =$

 $117\,324\,(2)\times 10^{-8}$ (see [6] and references therein), where the uncertainty $\pi^2\ln^2(m_\tau/m_e)(\alpha/\pi)^4\sim 2\times 10^{-8}$ has been assigned for uncalculated four-loop contributions. The errors due to the uncertainties of the $O(\alpha^2)$ (5×10^{-10}) and $O(\alpha^3)$ terms (3×10^{-11}), as well as that induced by the uncertainty of α (8×10^{-13}), are negligible. The sum of the one- and two-loop EW contributions is $a_\tau^{\rm EW}=47.4(5)\times 10^{-8}$ (see [7, 2] and references therein). The uncertainty encompasses the estimated errors induced by hadronic loop effects, neglected two-loop bosonic terms and the missing three-loop contribution. It also includes the tiny errors due to the uncertainties in $M_{\rm top}$ and m_τ .

Similarly to the case of the muon g-2, the leadingorder hadronic contribution to a_{τ} is obtained via a dispersion integral of the total hadronic cross section of the e^+e^- annihilation (the role of low energies is very important, although not as much as for a_{μ}). The result of the latest evaluation, using the whole bulk of experimental data below 12 GeV, is $a_{\tau}^{\text{HLO}} = 337.5(3.7) \times 10^{-8}$ [2]. The hadronic higher-order (α^3) contribution a_{τ}^{HHO} can be divided into two parts: $a_{\tau}^{\text{HHO}} = a_{\tau}^{\text{HHO}}(\text{vp}) + a_{\tau}^{\text{HHO}}(\text{lbl})$. The first one, the $O(\alpha^3)$ contribution of diagrams containing hadronic self-energy insertions in the photon propagators, is $a_{\tau}^{\text{HHO}}(\text{vp}) = 7.6(2) \times 10^{-8}$ [8]. Note that naïvely rescaling the corresponding muon g-2 result by a factor m_{τ}^2/m_{μ}^2 leads to the incorrect estimate $a_{\tau}^{\text{HHO}}(\text{vp}) \sim -28 \times 10^{-8}$ (even the sign is wrong!). Estimates of the light-by-light contribution $a_{\tau}^{\text{HHO}}(\text{lbl})$ obtained rescaling the corresponding one for the muon g-2 by a factor m_{τ}^2/m_{μ}^2 fall short of what is needed – this scaling is not justified. The parton-level estimate of [2] is $a_{\tau}^{\text{HHO}}(\text{lbl}) = 5(3) \times 10^{-8}$, a value much lower than those obtained by naïve rescaling. Adding up the above contributions one obtains the SM prediction [2]

$$a_{\tau}^{\rm SM} = a_{\tau}^{\rm QED} + a_{\tau}^{\rm EW} + a_{\tau}^{\rm HLO} + a_{\tau}^{\rm HHO} = 117\,721\,(5) \times 10^{-8}. \eqno(1)$$

Errors were added in quadrature.

The present PDG limit on the τ lepton g-2 was derived in 2004 by the DELPHI collaboration from $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$ total cross section measurements at \sqrt{s} between 183 and 208 GeV at LEP2 (the study of a_τ via this channel was proposed in [9]): $-0.052 < a_\tau < 0.013$ at 95% confidence level [1]. Reference [1] also quotes the result in the form:

$$a_{\tau} = -0.018(17). \tag{2}$$

This limit was obtained comparing the experimentally measured values of the cross section to the SM values, assuming that possible deviations were due to non-SM values of a_{τ} . Therefore, Eq. (2) is actually a limit on

the non-SM contribution to a_{τ} . We refer to [10] for a concise review of older, less stringent limits on a_{τ} , derived from data samples of $e^+e^- \to \tau^+\tau^-$ (for photon squared four-momentum q^2 up to $(37 \text{ GeV})^2$) [11] and $Z^0 \to \tau^+\tau^-\gamma$ (where $q^2 = 0$ but the radiating τ is not on-shell) [12, 13, 14, 15]. Other indirect bounds on a_{τ} were studied in [16].

Comparing Eqs. (1) and (2) (their difference is roughly one standard deviation), it is clear that the sensitivity of the best existing measurements is still more than one order of magnitude worse than needed. The possibility to improve these limits is certainly not excluded. Future high-luminosity B factories such as Super-KEKB [17] offer new opportunities to improve the determination of the τ magnetic properties. The authors of [18] proposed to determine the Pauli form factor $F_2(q^2)$ of the τ via $\tau^+\tau^-$ production in e^+e^- collisions at the Υ resonances with sensitivities possibly down to $O(10^{-5})$ or even better. In [19] the reanalysis of various measurements of the cross section of the process $e^+e^- \rightarrow \tau^+\tau^-$, the transverse τ polarization and asymmetry at LEP and SLD, as well as the decay width $\Gamma(W \to \tau \nu_{\tau})$ at LEP and Tevatron allowed the authors to set a model-independent limit on new physics contributions.

$$-0.007 < a_{\tau}^{NP} < 0.005, \tag{3}$$

a bound stronger than that in Eq. (2). This analysis, like earlier ones, was performed without radiative corrections, but the authors checked that the inclusion of initial-state radiation did not affect significantly the obtained bounds.

3. The τ lepton EDM

In the SM, the only sources of CP violation are the CKM-phase and a possible θ -term in QCD. A fundamental EDM can only be generated at the three-loop level [20], yielding a SM contribution which is far below experimental capabilities. Models for physics beyond the SM generally induce large contributions to lepton and neutron EDMs, and although there has been no experimental evidence for an EDM so far, there is considerable hope to gain new insights into the nature of CP violation through this kind of experiments.

The current 95% confidence level limits on the EDM of the τ lepton are given by

$$-2.2 < \text{Re}(d_{\tau}) < 4.5 \ (10^{-17} \ e \text{ cm}),$$

-2.5 < Im(d_{\tau}) < 0.8 \ (10^{-17} \ e \ cm);

they were obtained by the Belle collaboration [5] following the analysis of [21] for the impact of an effective

Download English Version:

https://daneshyari.com/en/article/1845738

Download Persian Version:

https://daneshyari.com/article/1845738

<u>Daneshyari.com</u>