



Collapse and revival of electromagnetic cascades in focused intense laser pulses



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ABSTRACT

We consider interaction of a high-energy electron beam with superstrong laser pulses. Nonlinear Compton scattering and electron–positron pair production by the emitted photons result in development of an electromagnetic “shower-type” cascade, which however collapses rather quickly due to energy losses by secondary particles. Nevertheless, the laser field accelerates the slowed down electrons and positrons, thus giving rise to development of electromagnetic cascade of another type (“avalanche-type”). This effect of cascade collapse and revival can be observed at the electron beam energy of the order of several GeV and laser pulses intensity $\sim 10^{24}$ W/cm². This means that it can be readily observed at the novel laser facilities which are either planned for the nearest future, or are already under construction. The proposed experimental setup provides the most realistic and promissory way to observe the “avalanche-type” cascades.

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1. Introduction

Amazing progress in elaboration of high-power laser technologies was demonstrated in recent decades. By now, laser intensity has already surpassed the level of 10^{22} W/cm² [1]. Recently, several projects aiming at intensities up to 10^{26} W/cm² were announced and are currently under active development, see, e.g., [2,3]. When realized, they would open unique possibilities to observe a variety of yet unexplored exotic QED effects of nonlinear interaction of electromagnetic radiation with matter and vacuum, see reviews [4,5] for the details.

Experimental tests of Intense Field QED were pioneered almost two decades ago at SLAC [6], where both the nonlinear Compton effect and the nonlinear Breit–Wheeler process have been observed for the first time. The forthcoming facilities [2,3] would allow to reproduce the SLAC-like experiments at a principally new level. Instead of rare events of photon emission and e^+e^- -pair creation, the long chains of sequential events of those processes (cascades) would be observed [7], see also [8].

The cascades discussed in Refs. [7,8] are very similar to the well studied air showers which are induced by cosmic rays in atmosphere [9,10]. For the sake of brevity we will call them hereinafter *S (Shower)-type cascades*. However, another mechanism of

cascade development is much more intriguing. The point is that a cascade may occur due to acceleration of charged particles by the laser field and can be seeded even by a particle initially at rest [11–14]. Such cascades have a lot of similarities with the Townsend discharge in gases [15], or the avalanche breakdown in insulators or semiconductors, see, e.g., [16]. We will call them here *A (Avalanche)-type cascades*.¹ It is worth noting that since an A-type cascade withdraws the energy necessary for its development from the field, it may result in substantial depletion of the laser field [18] and, very likely, even bound the maximally attainable value of laser intensity [13].

In the case of S-type cascades, creation of e^+e^- pairs and hard photons occurs at the expense of kinetic energy of the initial and secondary particles. The multiplication factor for such a cascade depends also on the laser field strength. The higher is the field strength, the higher are the probabilities of particle creation processes and, hence, the multiplication factor. Nevertheless, the S-type cascade collapses at any laser intensity when the energy of secondary particles decreases to some threshold value. On the other hand, occurrence of the A-type cascades does not necessary require high energy of a primary particle. However, the presence of a very strong electric field is absolutely necessary to ensure fast energy gain for a particle initially at rest [13].

The goal of this paper is to show that it is possible to observe both types of cascades in the same SLAC-like experiment at the

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¹ Other denominations for the two types of cascades were proposed in Ref. [17].

forthcoming facilities [2,3]. This could be possible if the time of collapse for an S-type cascade were small compared to duration time of the laser pulse τ_L . Then, the mechanism of acceleration of relatively slow particles, which did not have enough time to leave the focal region, could turn on giving start of an A-type cascade development. In such a way we could encounter revival of the cascade process.

2. Model and estimations

We will consider an electron beam colliding with the field of two counterpropagating circularly polarized intense focused laser pulses with frequency ω and duration $\tau_L \gg 1/\omega$, so that magnetic component of the field vanishes in the focus of the standing wave and the field there can be considered as a rotating electric field, compare [13,14]. As in Ref. [14], we assume that for both pulses $E, H \ll E_S$, where $E_S = \frac{m^2 c^3}{e \hbar} = 1.32 \times 10^{16}$ V/cm is the QED critical field, and parameter $\xi = e \sqrt{-A^\mu A_\mu} / mc \gg 1$, where A_μ is the field 4-potential. We assume that the direction of the electron beam coincides with the direction of one of the laser pulses and the initial energy of the electrons $\varepsilon_0 \gg mc^2$.

For the case of a laser field of optical frequency and ultrarelativistic particles one can use probability rates for photon emission W_γ and e^+e^- -pair creation W_{cr} in the approximation of a locally constant crossed field [13]. Then, they are exclusively determined by the dynamical quantum parameters of participating particles

$$\chi_{e,\gamma} = \frac{e \hbar}{m^3 c^4} \sqrt{(p_0 \mathbf{E} + \mathbf{p} \times \mathbf{H})^2 - (\mathbf{pE})^2}, \quad (1)$$

and can be estimated as

$$W_{\gamma,cr} \sim \frac{\alpha m^2 c^4}{\hbar \varepsilon_{e,\gamma}} \chi_{e,\gamma}^{2/3}, \quad \chi_{e,\gamma} \gtrsim 1, \quad (2)$$

[19], where 4-momenta p^μ are: $p_e^\mu = (\varepsilon_e, \mathbf{p}_e)$ for electron or positron, and $k_\gamma^\mu = (\varepsilon_\gamma = \hbar \omega, \mathbf{k}_\gamma)$ for a photon, \mathbf{E} and \mathbf{H} are the local values of the electric and magnetic field. For small values of parameters (1), $W_\gamma \sim \frac{\alpha m^2 c^4}{\hbar \varepsilon_e} \chi_e$, while the probability of pair creation is suppressed exponentially $W_{cr} \sim \frac{\alpha m^2 c^4}{\hbar \varepsilon_\gamma} \chi_\gamma \exp(-8/3 \chi_\gamma)$.

Let us estimate duration of the S-type cascade τ_S first. The initial value of parameter (1) for a primary ultrarelativistic electron can be estimated as $\chi_0 \sim \frac{\varepsilon_0}{mc^2} \frac{E_0}{E_S}$, where E_0 is the peak value of the electric field strength in the focal region of the cumulative field of two colliding pulses. Any involved particle transforms into two in every event, so that multiplicity of the S-type cascade can be estimated as in Ref. [20],

$$2^n = \frac{\chi_0}{\chi_f}, \quad (3)$$

where n is the number of generations of secondary particles and χ_f is the value of parameter (1) for the final electrons. Thus, for τ_S we have $\tau_S \sim t_e n$, where $t_e \sim W_\gamma^{-1}(\varepsilon_0, \chi_0)$ is the mean lifetime of a primary electron with respect to hard photon emission. Finally,

$$\tau_S \sim \tau_C \frac{\varepsilon_0}{\alpha mc^2} \chi_0^{-2/3} \log_2 \left(\frac{\chi_0}{\chi_f} \right), \quad (4)$$

where $\tau_C = \hbar/mc^2$ is the Compton time. The primary electrons are supposed to be ultrarelativistic, so that $\chi_0 > 1$. Obviously, the S-type cascade collapses when $\chi_f < 1$. We will choose χ_f to be ~ 0.1 . At such χ_f , the development of the S-cascade cannot continue since the mean lifetime of a photon with respect to pair production, $t_\gamma \sim W_{cr}^{-1} \sim \exp(8/3 \chi_f)$, becomes exponentially large.

We are interested in the case when $n > 1$ and the collapse time of the S-type cascade is less than the laser pulse duration τ_L . Hence, the following requirements

$$t_{e,\gamma} < \tau_C \frac{\varepsilon_0}{\alpha mc^2} \chi_0^{-2/3} \log_2(10 \chi_0) < \tau_L, \quad (5)$$

should be respected. Let $\varepsilon_0 = 3$ GeV and $E_0 = 3.2 \times 10^{-3} E_S$. Then $\chi_0 \approx 20$, and the left condition in (5) is satisfied at the expense of large logarithm $\log_2(10 \chi_0) \approx 8$. The right condition in (5) is satisfied for $\tau_L \gtrsim 10$ fs.

It was shown in Refs. [13,14] that a charged particle acquires very large transverse, with respect to the direction of the laser pulses propagation, momentum under the action of rotating electric field, and the dynamical parameter χ of the particle gains the increment $\Delta \chi \sim 1$ within a small fraction $t_{acc} \sim \frac{E_S}{E_0} \sqrt{\frac{mc^2}{\hbar \omega}} \tau_C$ of the rotation period, $\omega t_{acc} \ll 1$. If trajectory of the particle contorts so strongly that it can emit a hard photon in the direction transverse with respect to its initial propagation, which by-turn can create a pair, the A-type cascade occurs. In our case, the initial particle has a large longitudinal momentum and accordingly large value of parameter χ . Therefore a noticeable contortion of the particle trajectory necessary for development of the A-type cascade will occur only if $\chi \sim 1$.

Let us explain this issue in more details. The momentum of a particle moving in the field of two colliding laser pulses can be represented as $\mathbf{p}(t) = \mathbf{p}_\parallel(t) + \mathbf{p}_\perp(t)$. The longitudinal component of the momentum $\mathbf{p}_\parallel(t)$ varies only due to emission of photons, and its characteristic variation time is $t_\parallel \sim \tau_S \lesssim \tau_L$. Variation of the transverse momentum component $\mathbf{p}_\perp(t)$ is determined by the rotating electric field $\mathbf{E}(t) = \{E_0 \cos(\omega t), E_0 \sin(\omega t), 0\}$,

$$\dot{\mathbf{p}}_\perp(t) = e \mathbf{E}(t),$$

and thus its characteristic variation time is $t_\perp \sim 1/\omega$. We assume that a charged particle arrives at the focal region at the moment $t = 0$ with zero transverse momentum, $\mathbf{p}_\perp(0) = 0$. Then taking into account the relation $t_\parallel/t_\perp \sim \omega \tau_L \gg 1$, we can calculate $\chi(t)$ according to Eq. (1) and get the formula

$$\chi(t) \approx \sqrt{\chi_\parallel^2(t) + (\Delta \chi_\perp(t))^2}, \quad (6)$$

valid within the time interval $\Delta t \lesssim 1/\omega$. Here $\chi_\parallel = \chi(0) = \frac{E_0}{E_S} \sqrt{1 + \frac{p_\parallel^2}{m^2}}$ has the meaning of the dynamical parameter in the absence of accelerating field, and $\Delta \chi_\perp = 2\xi \frac{E_0}{E_S} \sin^2(\frac{\omega t}{2})$, compare to Eq. (9b) in Ref. [14]. $\Delta \chi_\perp$ gains the value ~ 1 within $t_{acc} \ll 1/\omega$. This can essentially contort the trajectory and thus give start to an A-type cascade only if $\chi_\parallel \sim 1$, as it is seen from Eq. (6).

So, the A-type cascade will retard with respect to the S-type cascade start and the delay time τ_R is determined by Eq. (4) with $\chi_f = 1$. Consequently, we can estimate the duration of the A-type cascade τ_A as $\tau_A \sim \tau_L - \tau_R$. Certainly, the conditions $t_{acc} \ll t_{e,\gamma} < \tau_A$ must be respected. It is easy to check that they hold for the values of ε_0 , E_0 and τ_L adopted above.

3. Numerical simulations

For more accurate analysis we use Monte-Carlo simulations. We assume e^- and e^+ are moving classically between the acts of emission and solve relativistic equations of motion numerically. Time-points of the quantum events are determined in a manner similar to Refs. [22,21]. The code was tested by simulating cascades initiated by high-energy electrons in a constant homogeneous transverse magnetic field and by simulating cascades caused by electron seeded at rest in the uniformly rotating homogeneous electric field. The results are in reasonable agreement with the cascade profiles from Refs. [23] and [14] respectively. It is worth noting that in our simulations we use exact expressions for the probabilities of photon emission and e^-e^+ pair creation in a constant crossed

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