# Real spectral triples over noncommutative Bieberbach manifolds 

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#### Abstract

We classify and construct all real spectral triples over noncommutative Bieberbach manifolds, which are restrictions of irreducible, real, equivariant spectral triples over the noncommutative three-torus. We show that, in the classical case, the constructed geometries correspond exactly to spin structures over Bieberbach manifolds and the Dirac operators constructed for a flat metric.


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## 1. Introduction

Bieberbach manifolds are compact manifolds, which are quotients of the Euclidean space by a free, properly discontinuous and isometric action of a discrete group. The torus is the canonical example of a Bieberbach manifold, however, the first nontrivial example appears in dimension 2 and is a Klein bottle. The case $d=3$ has already been described in the seminal works of Bieberbach [1,2]. In this paper we work with the dual picture, looking at the suitable algebra of functions on the Bieberbach manifold (and their noncommutative counterparts) in terms of a fixed point subalgebra of the relevant dense subalgebra of the $C^{*}$ algebra of continuous functions on the three-torus and its corresponding noncommutative deformation $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right)$.

### 1.1. Noncommutative Bieberbach manifolds

In this section we shall briefly recall the description of three-dimensional noncommutative Bieberbach manifolds as quotients of the three-dimensional noncommutative tori by the action of a finite discrete group. For details we refer to [3], here we present the notation and the results. Out of 10 different Bieberbach three-dimensional manifolds (six orientable, including the three-torus itself and four nonorientable ones) only six have noncommutative counterparts.

[^0]Table 1
The action of finite cyclic groups on $A\left(T_{\theta}^{3}\right)$.

| Name | $\mathbb{Z}_{N}$ | Generator | Action of $\mathbb{Z}_{N}$ on $U, V, W$ |
| :--- | :--- | :--- | :--- |
| $\mathrm{B2}_{\theta}$ | $\mathbb{Z}_{2}$ | $h, h^{2}=\mathrm{id}$ | $h \triangleright U=-U, h \triangleright V=V^{*}, h \triangleright W=W^{*}$, |
| $\mathrm{B3}_{\theta}$ | $\mathbb{Z}_{3}$ | $h, h^{3}=\mathrm{id}$ | $h \triangleright U=e^{\frac{2}{3} \pi i} U, h \triangleright V=e^{-\pi i \theta} V^{*} W, h \triangleright W=V^{*}$, |
| $\mathrm{B4}_{\theta}$ | $\mathbb{Z}_{4}$ | $h, h^{4}=\mathrm{id}$ | $h \triangleright U=i U, h \triangleright V=W, h \triangleright W=V^{*}$, |
| $\mathrm{B6}_{\theta}$ | $\mathbb{Z}_{6}$ | $h, h^{6}=\mathrm{id}$ | $h \triangleright U=e^{\frac{1}{3} \pi i} U, h \triangleright V=W, h \triangleright W=e^{-\pi i \theta} V^{*} W$, |
| $\mathrm{N1}_{\theta}$ | $\mathbb{Z}_{2}$ | $h, h^{2}=\mathrm{id}$ | $h \triangleright U=U^{*}, h \triangleright V=-V, h \triangleright W=W$, |
| $\mathrm{N} 2_{\theta}$ | $\mathbb{Z}_{2}$ | $h, h^{2}=\mathrm{id}$ | $h \triangleright U=U^{*}, h \triangleright V=-V, h \triangleright W=W U^{*}$, |

Definition 1.1 (See [3]). Let $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right)$ be an algebra of smooth elements on a three-dimensional noncommutative torus, which contains the polynomial algebra generated by three unitaries $U, V, W$ satisfying relations,

$$
U V=V U, \quad U W=W U, \quad W V=e^{2 \pi i \theta} V W
$$

where $0<\theta<1$ is irrational. We define the algebras of noncommutative Bieberbach manifolds as the fixed point algebras of the following actions of finite groups $G$ on $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right)$, which are combined in Table 1.

We have shown in [3] that the actions of the cyclic groups $\mathbb{Z}_{N}, N=2,3,4,6$ on the noncommutative three-torus, as given in Table 1is free. The aim of this paper is to study and classify flat (i.e. restricted from flat geometries of the torus $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right)$ ) real spectral geometries over the orientable noncommutative Bieberbachs.

## 2. Spectral triples over Bieberbachs

Since each noncommutative Bieberbach algebra is a subalgebra of the noncommutative torus, a restriction of the spectral triple over the latter to the subalgebra, gives a generic spectral triple over a noncommutative Bieberbach manifold, which might be, however, reducible. By restriction of a spectral triple $(\mathcal{A}, \pi, \mathscr{H}, D, J)$ to a subalgebra, $\mathcal{B} \subset \mathcal{A}$, we understand the triple $\left(\mathscr{B}, \pi, \mathscr{H}^{\prime}, D^{\prime}, J^{\prime}\right)$ where $\pi^{\prime}$ is the restriction of $\pi$ to $\mathscr{B}, \mathscr{H}^{\prime} \subset \mathscr{H}$ is a subspace invariant under the action of $\mathscr{B}, D$ and $J$, so that $D^{\prime}, J^{\prime}$ are their restrictions to $\mathscr{H}^{\prime}$ (note that in the even case this must apply also to $\gamma$ ).

In what follows we shall show that, in fact, each spectral triple over Bieberbach is a restriction of a spectral triple over the torus, first showing that each spectral triple over Bieberbach can be lifted to a noncommutative torus.

### 2.1. The lift and the restriction of spectral triples

Lemma 2.1. Let $\left(B N_{\Theta}, \mathscr{H}, J, D\right)$ be a real spectral triple over a noncommutative Bieberbach manifold $B N_{\theta}$. Then, there exists a spectral triple over a three-torus, such that this triple is its restriction.
Proof. In [3] we showed that the crossed product algebra $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right) \rtimes \mathbb{Z}_{N}$ is isomorphic to the matrix algebra of the noncommutative Bieberbach manifold algebra:

$$
\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right) \rtimes \mathbb{Z}_{N} \sim \mathrm{BN}_{\theta} \otimes M_{N}(\mathbb{C})
$$

First, let us recall that any spectral triple $\mathcal{A}, \mathcal{H}, D, J$ could be lifted to a spectral triple over $\mathcal{A} \otimes M_{n}(\mathbb{C})$. Indeed, if we take $\mathscr{H}^{\prime}=\mathscr{H} \otimes M_{n}(\mathbb{C})$ with the natural representation $\pi^{\prime}(a \otimes m)(h \otimes M)=\pi(a) h \otimes m M$, the diagonal Dirac operator and $J^{\prime}(h \otimes M)=J h \otimes U M^{\dagger} U^{\dagger}$, for an arbitrary unitary $U \in M_{n}(\mathbb{C})$ it is easy to see that we obtain again a real spectral triple. Applying this to the case of $\mathrm{B} N_{\theta}$, and identifying $\mathrm{B} N_{\theta} \otimes M_{N}(\mathbb{C})$ using the above isomorphism we obtain a spectral triple over $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right) \rtimes \mathbb{Z}_{N}$. As $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right)$ is a subalgebra of $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right) \rtimes \mathbb{Z}_{N}$ by restriction we obtain, in turn, a spectral triple over a three-dimensional noncommutative torus. In fact, it is easy to see that we obtain a spectral triple, which is equivariant with respect to the action of a $\mathbb{Z}_{N}$ group. Clearly, the fact that we have a representation of the crossed product algebra is just a rephrasing of the fact that we have a $\mathbb{Z}_{N}$-equivariant representation of $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right)$. By definition, the Dirac operator lifted from the spectral triple over $\mathrm{B} N_{\theta}$ commutes with the group elements which are identified as matrices in $M_{N}(\mathbb{C})$. A little care is required to show that the lift of $J$ would properly commute with the generator of $\mathbb{Z}_{N}$. However, since the lift of $J$ involves a matrix $U$, it is sufficient to use a matrix, which provides a unitary equivalence between the generator $h$ and its inverse $h^{-1}$ of $\mathbb{Z}_{N}$ in $M_{N}(\mathbb{C})$. Simple computation shows that the following $U, U_{00}=1, U_{k l}=\delta_{k, N-l}$ for $k, l=1, \ldots, N-1$ is the one providing the equivalence.

Hence, by this construction, we obtain a real, $\mathbb{Z}_{N}$-equivariant spectral triple over $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right)$. It is easy to see that the original spectral triple over $B N_{\theta}$ is a restriction of the constructed spectral triple by taking the invariant subalgebra of $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right)$, the $\mathbb{Z}_{N}$-invariant subspace of $\mathscr{H}$ and the restriction of $D$ and $J$.

Remark 2.2. The procedure described above does not necessarily provide the canonical (equivariant) Dirac operator over $\mathcal{A}\left(\mathbb{T}_{\theta}^{3}\right)$. Indeed, even a simple example of $\mathcal{A}\left(\mathbb{T}^{1}\right)$ shows that the lifted Dirac operator differs from the fully equivariant one by a bounded term. There may exist, however, a fully equivariant triple, so that its restriction is the same triple we started with.

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