



# An adjusted temperature wall function for turbulent forced convective heat transfer for bluff bodies in the atmospheric boundary layer

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## ARTICLE INFO

### Article history:

Received 27 February 2011

Received in revised form

11 April 2011

Accepted 11 April 2011

### Keywords:

Wall function

Computational fluid dynamics

RANS

Convective heat transfer coefficient

Cube

Urban heat transfer

## ABSTRACT

Accurate convective heat transfer predictions are required in building engineering and environmental studies on urban heat islands, building energy performance, building-envelope durability or conservation and (natural) ventilation of buildings. When applying computational fluid dynamics (CFD) for these computationally-expensive studies at high-Reynolds numbers, wall functions are mostly used to model the boundary-layer region. In this study, an adjustment to the standard temperature wall function is proposed for forced convective heat transfer at surfaces of typical wall-mounted bluff bodies in turbulent boundary layers, such as the atmospheric boundary layer, at moderate to high Reynolds numbers. The methodology to determine this customised temperature wall function (CWF) from validated numerical data of CFD simulations using low-Reynolds number modelling (LRNM) is explained, where a logarithmic-law behaviour is found. The performance of this CWF is evaluated for several bluff-body configurations. Standard wall functions (SWFs) yield deviations of about 40% for the convective heat transfer coefficient, compared to LRNM. With the CWF however, these deviations are reduced to about 10% or lower. The CWF therefore combines increased (wall-function) accuracy for convective heat transfer predictions with the typical advantage of wall functions compared to LRNM, being a lower grid resolution in the near-wall region, which increases computational economy and facilitates grid generation. Furthermore, this CWF can be easily implemented in existing CFD codes, and is implemented in the commercial CFD code Fluent in this study.

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## 1. Introduction

Forced convective heat transfer at surfaces of wall-mounted bluff bodies in turbulent boundary layers at moderate to high Reynolds numbers ( $Re = 10^4$ – $10^7$ ) is of interest in many engineering applications, such as the evaluation of wind-induced convective heat losses from building surfaces or building components (e.g. solar collectors) in the atmospheric boundary layer (ABL). In building and urban engineering, convective heat transfer predictions are especially relevant for ABL flow applications on urban heat islands [1,2], building (component) energy performance [3,4], building-envelope durability or conservation [5,6] and (natural) ventilation of buildings [7]. Furthermore, they can be used

to estimate the convective moisture transfer from building surfaces, by using the heat and mass transfer analogy [8]. Convective moisture transfer is especially of interest for hygrothermal analysis of building envelopes and for urban applications involving evaporation of water from ponds, roof ponds, green roofs, green walls or surfaces which are wetted by (wind-driven) rain [9].

Convective heat transfer research for this type of flow problem is mainly performed by wind-tunnel experiments [10–14] and by computational fluid dynamics (CFD) studies [15–17] for both isolated bodies as well as arrays of bluff bodies. Compared to wind-tunnel experiments, CFD simulations have the advantage that usually a higher spatial resolution is obtained. With CFD also no restrictions are imposed regarding scaling and accessibility of certain surfaces, in contrast to wind-tunnel infrared thermography measurements [11] for example, which is especially important for more complex configurations. On the other hand, the applied numerical modelling approaches determine to a large extent the accuracy of CFD simulations.

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Nomenclature			
$A$	slope	$y_P^+$	dimensionless wall (normal) distance of cell centre point $P$ of wall-adjacent cell, $u_\tau y_P / \nu$
$B$	intercept	$y_P^*$	dimensionless wall (normal) distance of cell centre point $P$ of wall-adjacent cell
$c_p$	specific heat capacity of air (J/kg.K)	$y_T^*$	lower limit for temperature wall function
$C_\mu$	coefficient used in the k- $\epsilon$ turbulence models (0.09)	$y^+$	dimensionless wall (normal) distance, $u_\tau y / \nu$
$E$	constant in wall-function expression (9.793)	$y^*$	dimensionless wall (normal) distance
$H$	cube height (m)	$z$	height above the ground (m)
$k$	turbulent kinetic energy ( $\text{m}^2/\text{s}^2$ )	$z_0$	aerodynamic roughness length (m)
$k_P$	turbulent kinetic energy in cell centre point $P$ of wall-adjacent cell ( $\text{m}^2/\text{s}^2$ )	<b>Greek symbols</b>	
$P$	cell centre point of wall-adjacent cell	$\epsilon$	turbulence dissipation rate ( $\text{m}^2/\text{s}^3$ )
$P_f$	empirically-determined coefficient (function of $Pr_t$ )	$\kappa$	von Karman constant (0.4187)
$Pr$	molecular Prandtl number	$\mu$	dynamic viscosity of air (kg/ms)
$Pr_t$	turbulent Prandtl number	$\nu$	kinematic viscosity of air ( $\text{m}^2/\text{s}$ )
$Pr_{t,CWF}$	turbulent Prandtl number used by CWF	$\rho$	density of air ( $\text{kg}/\text{m}^3$ )
$q_{c,w}$	convective heat flux at the surface ( $\text{W}/\text{m}^2$ )	$\tau_w$	wall shear stress ( $\text{kg}/\text{ms}^2$ )
$Re$	Reynolds number	$\phi$	incidence angle ( $^\circ$ )
$T$	air temperature ( $^\circ\text{C}$ )	<b>Abbreviations</b>	
$T_P$	temperature in cell centre point $P$ of wall-adjacent cell ( $^\circ\text{C}$ )	ABL	atmospheric boundary layer
$T_P^*$	dimensionless temperature in cell centre point $P$ of wall-adjacent cell	CFD	computational fluid dynamics
$T_{ref}$	reference temperature ( $^\circ\text{C}$ )	CHTC	convective heat transfer coefficient
$T_w$	wall temperature ( $^\circ\text{C}$ )	CWF	customised wall function
$T^+$	dimensionless temperature, $\rho c_p u_\tau (T_w - T) / q_{c,w}$	DNS	direct numerical simulation
$T^*$	dimensionless temperature	LES	large-eddy simulation
$u_{ABL}^*$	ABL friction velocity (m/s)	LRNM	low-Reynolds number modelling
$u_\tau$	boundary-layer friction velocity, $(\tau_w / \rho)^{1/2}$ (m/s)	RANS	Reynolds-averaged Navier-Stokes
$U$	mean air speed (m/s)	SWF	standard wall function
$U_{10}$	mean air speed in the undisturbed flow at a height of 10 m above the ground (m/s)	WF	wall function
$U^+$	dimensionless mean air speed, $U / u_\tau$	<b>Subscripts</b>	
$U^*$	dimensionless mean air speed	LRNM	low-Reynolds number modelling
$y$	distance (normal) from the wall (m)	P	cell centre point $P$ of wall-adjacent cell
$y_P$	distance (normal) of cell centre point $P$ of wall-adjacent cell to the wall (m)	WF	wall function

At high Reynolds numbers ( $Re \approx 10^6 - 10^7$ ), CFD computations for this type of flow are often performed with (steady) Reynolds-averaged Navier-Stokes (RANS) combined with wall functions (WFs) [18], especially for complex configurations, for example in large-scale environmental studies [6,7,19–22]. WFs are used here to take care of the boundary-layer region, instead of low-Reynolds number modelling (LRNM), mainly for reasons of computational economy: at high Reynolds numbers, the use of WFs, which model the flow quantities in the boundary-layer region by calculating them by means of semi-empirical functions, avoids an extremely high grid resolution here, which would be required for LRNM, where the boundary layer is resolved explicitly. Both boundary-layer modelling approaches are described more in detail below (section 3).

The standard formulation for WFs [23], referred to as standard wall functions (SWFs), however has two main limitations:

1. SWFs impose strict requirements to the computational grid: the cell centre point  $P$  of the wall-adjacent cell has to be located outside of the viscosity-affected region (viscous sublayer and buffer layer), i.e.  $y_P^+ > 30$ . On the other hand, the point  $P$  has to be located close enough to the wall for it to be in the logarithmic layer, i.e.  $y_P^+ < 500$  [24], and to ensure a sufficiently high grid resolution in the boundary layer. The dimensionless wall distance of point  $P$  ( $y_P^+$ ) is defined as:

$$y_P^+ = \frac{\sqrt{\tau_w} y_P}{\nu} \quad (1)$$

where  $y_P$  is the distance (normal) from the cell centre point  $P$  of the wall-adjacent cell to the wall,  $\tau_w$  is the shear stress at the wall,  $\rho$  is the air density and  $\nu$  is the kinematic viscosity of air. For complex flows, these  $y_P^+$  criteria are difficult to achieve throughout the entire computational domain, especially if automated grid generation is used, by which WFs are sometimes used outside of their specified validity range.

2. The wall-function concept of SWFs is based on a universal behaviour of the boundary layer, in terms of velocity, turbulence and temperature profiles. These SWFs are derived for wall-attached boundary layers under so-called equilibrium conditions, i.e. small pressure gradients, local equilibrium between generation and dissipation of turbulent energy and a constant (uniform) shear stress and heat flux in the near-wall region [25]. This wall-function concept breaks down for more complex flows, such as flow around bluff bodies where the boundary layer does not remain attached to the surface. SWFs can therefore lead to inaccurate predictions, especially for wall friction and convective heat transfer [25,26].

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