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# On sound scattering by rigid edges and wedges in a flow, with applications to high-lift device aeroacoustics



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#### ABSTRACT

Exact analytical solutions for the scattering of sound by the edge of a rigid half-plane and by a rigid corner in the presence of a uniform flow are considered in this work, for arbitrary source and observer locations. Exact Green's functions for the Helmholtz equation are first reviewed and implemented in a quiescent propagation space from reference expressions of the literature. The effect of uniform fluid motion is introduced in a second step and the properties of the field are discussed for point dipoles and quadrupoles. The asymptotic regime of a source close to the scattering edge/wedge and of an observer far from it in terms of acoustic wavelengths is derived in both cases. Its validity limits are assessed by comparing with the exact solutions. Typically the asymptotic directivity is imposed by Green's function but not by the source itself. This behaviour is associated with a strong enhancement of the radiation with respect to what the source would produce in free field. The amplification depends on the geometry, on the source type and on the source distance to the edge/wedge. Various applications in aeroacoustics of wall-bounded flows are addressed, more specifically dealing with high-lift device noise mechanisms, such as trailing-edge or flap side-edge noise. The asymptotic developments are used to highlight trends that are believed to play a role in airframe noise.

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#### 1. Introduction

The role of solid surfaces in aeroacoustics is twofold. Firstly they participate in the generation of sound by direct interaction with flows. Secondly they redistribute the sound radiated by sources possibly located elsewhere. In that sense they act as either sources of sound or scattering obstacles and can be artificially classified as active or passive surfaces, respectively. The physical understanding and the modeling of both aspects are key issues in the definition of noise reduction strategies in many engineering applications. Pure sound scattering is usually investigated using the classical theory of diffraction in linear acoustics. Aerodynamic sound production in the presence of solid surfaces can be formulated from the standpoint of linear acoustics by resorting to the acoustic analogy. According to Lighthill's original statement [1] and related

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#### Nomenclature

Italic symbols

		$\mathbf{x} = (r \ \theta)$	<i>z</i> ) cylindrical observer coordinates for an edge
$A_0$	Kutta-correction factor	$\mathbf{x}_0 = (r_0)$	$\theta_0$ $z_0$ ) cylindrical source coordinates for
С	chord length	0 (.0,	an edge
<i>c</i> <sub>0</sub>	speed of sound	$\mathbf{x} = (\mathbf{x}_1, \mathbf{x}_2)$	$(x_2, x_3)$ observer Cartesian coordinates for
D	directivity factor	$\mathbf{n} = (n_1, n_2)$	a corner
Ε	Fresnel integral	$\mathbf{x}_{a} = (v)$	$v_{-}v_{-}$ ) source Cartesian coordinates for
F	complex function for the 2D half-plane	$\mathbf{A}_{0} = 0_{1},$	a corner
	Green's function	$\mathbf{X} = (\hat{r}, \Theta)$	$(\phi)$ observer spherical coordinates for a corner
F	dipole strength vector	$\mathbf{x}_0 = (\hat{r}_0, \hat{r}_0)$	$(\Theta_0, \phi_0)$ source spherical coordinates for
G	Green's functions for the Helmholtz equation	5 ( <del>6</del> ,	a corner
$G_K$	correction to Green's function for the Kutta	$X_0, X = x$	$x_0/\beta, x/\beta$ modified streamwise coordinates for
C	Creen's functions with uniform flow		an edge
$G_{M_0}$	2D asymptotic Creen's function in free field	$X_1, Y_1 =$	$x_1/\beta, y_1/\beta$ modified streamwise coordinates
G∞ h	flap side edge thickness		for a corner
п U	Harviside function		
н г ц(1)	Passal and Hankal functions of the first kind	Greek sy	mbols
$J_{\nu}, \Pi_{\nu}$	acoustic wayanumbar		
$\kappa = \omega/c_0$	acoustic wavenumber	$\alpha_d$	dipole inclination angle
$\kappa = \kappa/\rho$	dimensionless coredunamis wavenumber	0 /1	$M^2$ compressibility parameter
$\kappa_1$	modified Rescal functions	$p = \sqrt{1}$	-M <sub>0</sub> compressionity parameter
$\kappa_n$	unstoady lift on a flan	$\varepsilon_m$	constant in Green's function for rigid wedge
ι 1	size of guadrilatoral elements	$\theta_1$	angle in the correction to Green's function for
l <sub>max</sub>	Size of quadrinateral elements		the Kutta condition
w <sub>0</sub>		$\theta_0, \theta$	corrected spherical source and observer angles
р D	dinale strongth	λ	acoustic wavelength
P D <sup>u</sup>		$\phi^{2D}, \phi^{3D}$	two-or three-dimensional acoustic potentials
$P'_{\nu}$	general Legendre functions	$\Phi$	wedge aperture angle
q	monopole strengtn	$\varphi$	projection angle for the half-plane Green's
Q _	quadrupole strength tensor		function
$(r,\theta,z)$	corrected observer cylindrical coordinates	$\rho_0$	fluid density
$(r_0, \theta_0, z_0)$	) corrected source cylindrical coordinates	ω	angular frequency
$r_{>} = ma$	$(r, r_0)$ maximum distance		
$r_{<} = mi$	$n(r, r_0)$ minimum distance	Subscrip	ts/superscripts symbols
$r, r_{1,2}, r_1$	,2 corrected 3D and 2D scattering distances		
$S_{1,2}, S_{1,2}$	3D and 2D integral bounds in convected	к. т	summation indices
	Green's functions	K	Kutta-condition correction
$S_0$	convection-corrected distance	Mo	flow-corrected quantity
$T_{i,j}$	quadrupole strength components	(1/2)	half-plane
$U_0$	flow speed	(1/2)	

works by Howe [2] and Ffowcs Williams and Hall [3], for instance, unsteady flow patterns interacting with solid surfaces can be interpreted as equivalent quadrupoles distributed in the fluid, the direct sound of which is scattered by the surfaces. This view is developed in some asymptotic theories of high-lift device noise [4,5]. According to Ffowcs Williams and Hawkings' statement of the analogy [6], surfaces explicitly involved in noise generation can also be mathematically interpreted as equivalent sources of lower orders. A priori the latter point of view is well suited for active surfaces and the former for passive surfaces. But the distinction is questionable when two bodies in close vicinity of each other are embedded in a disturbed flow region. The present analysis is dedicated to the high-lift devices that are deployed when the wing of an aircraft operates in approach and landing conditions. The source and scattering surfaces are implicitly assumed to be well separated, which only covers a part of the complete physics in most cases of interest. For instance a deployed flap (Fig. 1) can be interpreted as a distribution of equivalent sources, the sound of which is scattered by the main part of the wing. The underlying mechanism partly contributes to the airframe noise that also includes landing-gear associated sources. Airframe noise in itself is recognized as a major contributor to the total noise of an aircraft at approach, essentially because the modern high by-pass ratio engines are much quieter at idle power.

Going into the details, High-Lift Device (HLD) noise involves distributed sources along the leading or trailing edges of the wing ((2) in Fig. 1), the slat and the flap ((1) and (4) in Fig. 1), as well as sources that concentrate around the span ends of wing elements such as flap side-edges ((3) in Fig. 1) or slat corners. In the present study two contributions in which sound

for

coordinates

 $\mathbf{x} = (x, y, z)$  observer Cartesian coordinates for an edge

 $\mathbf{x}_0 = (x_0, y_0, z_0)$  source Cartesian

an edge

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