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# Wavelet coorbit spaces viewed as decomposition spaces



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#### ABSTRACT

In this paper we show that the Fourier transform induces an isomorphism between the coorbit spaces defined by Feichtinger and Gröchenig of the mixed, weighted Lebesgue spaces  $L_v^{p,q}$  with respect to the quasi-regular representation of a semi-direct product  $\mathbb{R}^d \rtimes H$  with suitably chosen dilation group H, and certain decomposition spaces  $\mathcal{D}(\mathcal{Q}, L^p, \ell_u^q)$ (essentially as introduced by Feichtinger and Gröbner) where the localized "parts" of a function are measured in the  $\mathcal{F}L^p$ -norm.

This equivalence is useful in several ways: It provides access to a Fourier-analytic understanding of wavelet coorbit spaces, and it allows to discuss coorbit spaces associated to different dilation groups in a common framework. As an illustration of these points, we include a short discussion of dilation invariance properties of coorbit spaces associated to different types of dilation groups.

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#### 1. Introduction

There exist several methods in the literature for the construction of higher-dimensional wavelet systems. A rather general class of constructions follows the initial inception of the continuous wavelet transform in [21] and uses the language of group representations [26,1,16,24]: Picking a suitable matrix group  $H \leq \text{GL}(\mathbb{R}^d)$ , the **dilation group**, one defines the associated semidirect product  $G = \mathbb{R}^d \rtimes H$ . This group acts on  $L^2(\mathbb{R}^d)$  via the (unitary) **quasi-regular representation**  $\pi$  defined by

$$(\pi(x,h)f)(y) = |\det(h)|^{-1/2} f(h^{-1}(y-x)), \quad (x,h) \in \mathbb{R}^d \times H.$$

The associated continuous wavelet transform of a signal  $f \in L^2(\mathbb{R}^d)$  is then obtained by picking a suitable mother wavelet  $\psi \in L^2(\mathbb{R}^d)$ , and letting

$$W_{\psi}f: G \to \mathbb{C} , \ (x,h) \mapsto \langle f, \pi(x,h)\psi \rangle .$$
 (1.1)

A wavelet  $\psi$  is called **admissible** if the operator  $W_{\psi}$  is (a multiple of) an isometry as a map into  $L^2(G, \mu_G)$ , where  $\mu_G$  denotes a left Haar measure on G. By definition we thus have for admissible vectors  $\psi$  that

$$\forall f \in \mathcal{L}^{2}(\mathbb{R}^{d}) : ||f||_{2}^{2} = \frac{1}{C_{\psi}} \cdot \int_{H} \int_{\mathbb{R}^{d}} |W_{\psi}f(x,h)|^{2} \, \mathrm{d}x \, \frac{\mathrm{d}h}{|\mathrm{det}(h)|}$$

alternatively expressed in the weak-sense inversion formula

$$f = \frac{1}{C_{\psi}} \cdot \int_{H} \int_{\mathbb{R}^d} W_{\psi} f(x,h) \cdot \pi(x,h) \psi \, \mathrm{d}x \, \frac{\mathrm{d}h}{|\mathrm{det}\,(h)|} \, .$$

An alternative construction of wavelet systems, with somewhat less structure but higher design flexibility, is the semi-discrete approach described as follows: Pick a discretely labeled quadratic partition of unity  $(\hat{\psi}_i)_{i \in I}$  in frequency domain, i.e. a family of functions satisfying

$$\forall_{a.e.}\xi \in \mathbb{R}^d : \sum_{i \in I} \left| \widehat{\psi}_i(\xi) \right|^2 = 1$$
(1.2)

and consider the system of all translates of the inverse Fourier transforms  $\psi_i = \mathcal{F}^{-1}(\widehat{\psi}_i)$ . This system is a (continuously labeled) tight frame, i.e.

$$\forall f \in \mathcal{L}^2(\mathbb{R}^d) : \|f\|_2^2 = \sum_{i \in I} \int_{\mathbb{R}^d} |\langle f, L_x \psi_i \rangle|^2 \, \mathrm{d}x,$$

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