



# Large time behavior of solutions for hyperbolic balance laws

Hongjun Yu

School of Mathematical Sciences, South China Normal University, Guangzhou 510631, PR China

Received 8 March 2016

Available online 18 August 2016

## Abstract

We study the existence and the large time behavior of global solutions to the initial value problem for hyperbolic balance laws in  $n$  space dimensions with  $n \geq 3$  admitting an entropy and satisfying the stable condition. We first construct global existence of the solutions to such a system around a steady state if the initial energy is small enough. Then we show that  $k$ -order derivatives of these solutions approach a constant state in the  $L^p$ -norm at a rate  $O(t^{-\frac{1}{2}(k+\rho+\frac{n}{2}-\frac{n}{p})})$  with  $p \in [2, \infty]$  and  $\rho \in [0, \frac{n}{2}]$  provided that initially  $\|z_0\|_{\dot{B}_{2,\infty}^{-\rho}} < \infty$ , where  $\dot{B}_{2,\infty}^{-\rho}$  is a homogeneous Besov space. These decay results do not impose an additional smallness assumption on  $L^p$  norm of the initial data and we thus improve the results in [3,19]. We also show faster decay results in the sense that if  $\|\mathbf{P}z_0\|_{\dot{B}_{2,\infty}^{-\rho}} + \|(\mathbf{I} - \mathbf{P})z_0\|_{\dot{B}_{2,\infty}^{-\rho+1}} < \infty$  with  $\rho \in (\frac{n}{2}, \frac{n+2}{2}]$ ,  $k$ -order derivatives of the solutions approach a constant state in the  $L^p$ -norm at a rate  $O(t^{-\frac{1}{2}(k+\rho+1+\frac{n}{2}-\frac{n}{p})})$ .

© 2016 Elsevier Inc. All rights reserved.

## 1. Introduction

This work concerns the following  $N$ -component hyperbolic system of balance laws in  $n$  space dimensions:

E-mail address: [yuhj2002@sina.com](mailto:yuhj2002@sina.com).

$$w_t + \sum_{j=1}^n f^j(w)_{x_j} = g(w), \tag{1.1}$$

with the initial conditions

$$w(0, x) = w_0(x), \tag{1.2}$$

where  $t \geq 0$  and  $x \in \mathbf{R}^n$  with  $n \geq 3$ ,  $w = w(t, x) \in \mathcal{O}_w \subset \mathbf{R}^N$  with  $N \geq 1$  and  $\mathcal{O}_w$  is open and convex. Here  $f^j(w)$  ( $j = 1, \dots, n$ ) and  $g(w)$  are given  $N$ -vector valued smooth functions of  $w$  from  $\mathcal{O}_w$  into  $\mathbf{R}^N$ . We define

$$\tilde{\mathcal{N}} = \{\psi \in \mathbf{R}^N; \psi^T g(w) = 0 \text{ for any } w \in \mathcal{O}_w\}.$$

Here the superscript  $T$  denotes the transpose. We set the dimension of  $\tilde{\mathcal{N}}$  to be  $N_1$  with  $1 \leq N_1 \leq N - 1$  and  $\tilde{\mathcal{N}}^\perp$  to be the orthogonal complement of  $\tilde{\mathcal{N}}$ . Then we have  $g(w) \in \tilde{\mathcal{N}}^\perp$ .

Such nonlinear systems typically govern non-equilibrium processes in physics for media with hyperbolic response as, for example, in gas dynamics. They also arise in the numerical simulation of conservation laws by relaxation schemes, see [1,3,4,8,12,37] and the references cited therein.

To obtain the existence of solutions to (1.1) and (1.2) the entropy condition is needed as in [9,10,18,37]. Here we first give the definition of entropy [18,19].

**Definition 1.1.** Let  $\eta = \eta(w)$  be a smooth function defined in  $\mathcal{O}_w$ . Then  $\eta(w)$  is called an entropy of balance laws (1.1) if the following four statements hold:

- (a)  $\eta(w)$  is strictly convex in  $\mathcal{O}_w$  in the sense that the Hessian  $D_w^2 \eta(w)$  is positive definite for  $w \in \mathcal{O}_w$ .
- (b) The matrix  $D_w f^j(w)(D_w^2 \eta(w))^{-1}$  is symmetric for  $w \in \mathcal{O}_w$  and  $j = 1, \dots, n$ .
- (c) Let  $w \in \mathcal{O}_w$ . Then  $g(w) = 0$  holds if and only if  $(D_w \eta(w))^T \in \tilde{\mathcal{N}}$ .
- (d) For  $w \in \mathcal{O}_w$  satisfying  $g(w) = 0$ , the matrix  $D_w g(w)(D_w^2 \eta(w))^{-1}$  is symmetric and non-positive definite and its null space coincides with  $\tilde{\mathcal{N}}$ .

Let  $\eta = \eta(w)$  be defined as above and put  $u(w) = (D_w \eta(w))^T$ . It was shown in [18] that the mapping  $u = u(w)$  is a diffeomorphism from  $\mathcal{O}_w$  onto its range  $\mathcal{O}_u$ . Suppose that  $w = w(u)$  is the inverse mapping from  $\mathcal{O}_u$  onto  $\mathcal{O}_w$ . Putting  $w = w(u)$  in (1.1) we have

$$\bar{A}^0(u)u_t + \sum_{j=1}^n \bar{A}^j(u)u_{x_j} = H(u). \tag{1.3}$$

Here

$$\bar{A}^0(u) := D_u w(u), \quad H(u) := g(w(u)), \tag{1.4}$$

$$\bar{A}^j(u) := D_u f_j(w(u)) = D_w f^j(w(u))D_u w(u). \tag{1.5}$$

In addition we see that  $H(u) \in \tilde{\mathcal{N}}^\perp$  for any  $u \in \mathcal{O}_u$ . Let us define

$$\bar{L}(u) := -D_u H(u) = -(D_w g)(w(u))D_u w(u). \tag{1.6}$$

Download English Version:

<https://daneshyari.com/en/article/4609272>

Download Persian Version:

<https://daneshyari.com/article/4609272>

[Daneshyari.com](https://daneshyari.com)