

A generalized regularity criterion for 3D Navier–Stokes equations in terms of one velocity component

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Abstract

In this article, we study the global regularity of the 3D Navier–Stokes equations in terms of one velocity component. In particular, we establish a new version of regularity criterion in the framework of the anisotropic Lebesgue space. It is a generalization of the notable work of C. Cao and E.S. Titi (2011) [3].

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1. Introduction

In the present paper, we address sufficient conditions for the regularity of the Leray–Hopf weak solution of the Navier–Stokes equations

$$\begin{cases} \frac{\partial u}{\partial t} - \nu \Delta u + (u \cdot \nabla)u + \nabla p = 0, & \text{in } \mathbb{R}^3, t \geq 0, \\ \nabla \cdot u = 0, & \text{in } \mathbb{R}^3, t \geq 0, \\ u(x, 0) = u_0, & \text{in } \mathbb{R}^3, \end{cases} \quad (1.1)$$

where $u = (u_1, u_2, u_3) : \mathbb{R}^3 \times \mathbb{R}^+ \rightarrow \mathbb{R}^3$ is the velocity field, $p : \mathbb{R}^3 \times \mathbb{R}^+ \rightarrow \mathbb{R}$ is a scalar pressure and u_0 is the initial velocity field, $\nu > 0$ is the viscosity. We set $\nabla_h = (\partial_{x_1}, \partial_{x_2})$ as the

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horizontal gradient operator and $\Delta_h = \partial_{x_1}^2 + \partial_{x_2}^2$ as the horizontal Laplacian, and Δ and ∇ are the usual Laplacian and the gradient operators, respectively. Here we use the classical notation

$$(u \cdot \nabla)v = \sum_{i=1}^3 u_i \partial_{x_i} v_k, \quad (k = 1, 2, 3), \quad \nabla \cdot u = \sum_{i=1}^3 \partial_{x_i} u_i,$$

and for the sake of simplicity, we denote ∂_{x_i} by ∂_i .

We set

$$\mathcal{V} = \{\phi : \text{the 3D vector valued } C_0^\infty \text{ functions and } \nabla \cdot \phi = 0\},$$

which will form the space of test functions. Let H and V be the closure spaces of \mathcal{V} in L^2 under L^2 -topology, and in H^1 under H^1 -topology, respectively. We denote by L^q and H^m the usual L^q -Lebesgue and Sobolev spaces, respectively, and by

$$\|\varphi\|_{L^q} = \left(\int_{\mathbb{R}^3} |\varphi(x)|^q dx_1 dx_2 dx_3 \right)^{\frac{1}{q}}, \quad \text{for every } \varphi \in L^q(\mathbb{R}^3),$$

we also denote by $L^q(\mathbb{R}_{x_2, x_3}^2, L^p(\mathbb{R}_{x_1}))$ the anisotropic Lebesgue space, and by

$$\|\varphi\|_{L_{x_1}^p L_{x_2, x_3}^q} = \left(\int_{\mathbb{R}^2} \left(\int_{\mathbb{R}} |\varphi(x)|^p dx_1 \right)^{\frac{q}{p}} dx_2 dx_3 \right)^{\frac{1}{q}}, \quad \text{for every } \varphi \in L^q(\mathbb{R}_{x_2, x_3}^2, L^p(\mathbb{R}_{x_1})).$$

We first recall some of the classical definitions and results on weak solutions of the Navier–Stokes equations.

Definition 1.1. A weak solution of (1.1) on $[0, T]$ (or $[0, \infty)$ if $T = \infty$) is a function $u : [0, T] \rightarrow L^2(\mathbb{R}^3)$ in the class

$$u \in C_w([0, T]; L^2(\mathbb{R}^3)) \cap L_{loc}^2(0, T; H^1(\mathbb{R}^3)),$$

satisfying

$$(u(t), \varphi(t)) + \int_0^t \{-(u, \partial_t \varphi) + \nu(\nabla u, \nabla \varphi) + (u \cdot \nabla u, \varphi)\} ds = (u_0, \varphi(0)), \quad (1.2)$$

for all $t \in [0, T]$ and all test functions $\varphi \in C_0^\infty([0, T] \times \mathbb{R}^3)$ with $\nabla \cdot \varphi = 0$. Here (\cdot, \cdot) stands for L^2 -inner product, and C_w signifies continuity in the weak topology.

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