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Iteration theory of *L*-index and multiplicity of brake orbits

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Abstract

In this paper, we first establish the Bott-type iteration formulas and some abstract precise iteration formulas of the Maslov-type index theory associated with a Lagrangian subspace for symplectic paths. As an application, we prove that there exist at least $[\frac{n}{2}] + 1$ geometrically distinct brake orbits on every C^2 compact convex symmetric hypersurface Σ in \mathbb{R}^{2n} satisfying the reversible condition $N\Sigma = \Sigma$. Furthermore, if all brake orbits on this hypersurface are nondegenerate, then there are at least *n* geometrically distinct brake orbits in every bounded convex symmetric domain in \mathbb{R}^n . Furthermore, if all brake orbits in this domain are nondegenerate, then there are at least *n* geometric case, we give a positive answer to the Seifert conjecture of 1948 under a generic condition.

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Keywords: Brake orbit; Maslov-type index; Bott-type iteration formula; Convex symmetric domain

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1. Introduction

Our aim of this paper is twofold. We first establish an iteration theory of the Maslov-type index associated with a Lagrangian subspace of $(\mathbb{R}^{2n}, \omega_0)$ for symplectic paths starting from identity. The Bott-type iteration formulas and some abstract precise iteration formulas are obtained here. Then as the application of this theory, we consider the brake orbit problem on a fixed energy hypersurface of the autonomous Hamiltonian systems. The multiplicity results are obtained in this paper.

1.1. Main results for the brake orbit problem

Let $V \in C^2(\mathbf{R}^n, \mathbf{R})$ and h > 0 such that $\Omega \equiv \{q \in \mathbf{R}^n \mid V(q) < h\}$ is nonempty, bounded, open and connected. Consider the following fixed energy problem of the second order autonomous Hamiltonian system

$$\ddot{q}(t) + V'(q(t)) = 0, \quad \text{for } q(t) \in \Omega,$$
(1.1)

$$\frac{1}{2} \left| \dot{q}(t) \right|^2 + V(q(t)) = h, \quad \forall t \in \mathbf{R},$$
(1.2)

$$\dot{q}(0) = \dot{q}\left(\frac{\tau}{2}\right) = 0,\tag{1.3}$$

$$q\left(\frac{\tau}{2}+t\right) = q\left(\frac{\tau}{2}-t\right), \qquad q(t+\tau) = q(t), \quad \forall t \in \mathbf{R}.$$
(1.4)

A solution (τ, q) of (1.1)–(1.4) is called a *brake orbit* in Ω . We call two brake orbits q_1 and $q_2 : \mathbf{R} \to \mathbf{R}^n$ geometrically distinct if $q_1(\mathbf{R}) \neq q_2(\mathbf{R})$.

We denote by $\mathcal{O}(\Omega)$ and $\tilde{\mathcal{O}}(\Omega)$ the sets of all brake orbits and geometrically distinct brake orbits in Ω respectively.

Let $J = \begin{pmatrix} 0 & -I \\ I & 0 \end{pmatrix}$ and $N = \begin{pmatrix} -I & 0 \\ 0 & I \end{pmatrix}$ with *I* being the identity in \mathbb{R}^n . Suppose that $H \in C^2(\mathbb{R}^{2n} \setminus \{0\}, \mathbb{R}) \cap C^1(\mathbb{R}^{2n}, \mathbb{R})$ satisfying

$$H(Nx) = H(x), \quad \forall x \in \mathbf{R}^{2n}.$$
(1.5)

We consider the following fixed energy problem

$$\dot{x}(t) = JH'(x(t)), \tag{1.6}$$

$$H(x(t)) = h, \tag{1.7}$$

$$x(-t) = Nx(t), \tag{1.8}$$

$$x(\tau + t) = x(t), \quad \forall t \in \mathbf{R}.$$
(1.9)

A solution (τ, x) of (1.6)–(1.9) is also called a *brake orbit* on $\Sigma := \{y \in \mathbb{R}^{2n} \mid H(y) = h\}$.

Remark 1.1. It is well known that via

$$H(p,q) = \frac{1}{2}|p|^2 + V(q), \qquad (1.10)$$

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