Contents lists available at ScienceDirect



Journal of Mathematical Analysis and Applications

www.elsevier.com/locate/jmaa

On regularity for the axisymmetric Navier–Stokes equations



 ^a School of Finance, Southwestern University of Finance and Economics, Chengdu 611130, Sichuan, PR China
 ^b School of Mathematics and Information Sciences, North China University of Water Resources and

^a School of Mathematics and Information Sciences, North China University of Water Resources a Electric Power, Zhengzhou, 450011, PR China
^c Department of Mathematics, Zhejiang Normal University, Jinhua 321004, Zhejiang, PR China

ARTICLE INFO

Article history: Received 9 March 2015 Available online 6 January 2016 Submitted by J. Guermond

Keywords: Navier–Stokes equations Axi-symmetric flow Blow-up criterion

ABSTRACT

In this paper, we establish a regularity criterion for the Navier–Stokes system with axisymmetric initial data. It is proved that if the local axisymmetric smooth solution u satisfies $\|u^{\theta} \mathbf{1}_{r \leq \zeta}\|_{L^{\alpha}((0,T^*);L^{\beta})} < \infty$ for any given $\varsigma > 0$, with $\frac{2}{\alpha} + \frac{3}{\beta} = 1$, $4 \leq \beta \leq 6$, then the strong solution keeps smoothness up to time T^* .

© 2015 Elsevier Inc. All rights reserved.

1. Introduction

We study the following classic 3D incompressible Navier–Stokes equations in the whole space:

$$\begin{cases} \partial_t u + (u \cdot \nabla u)u + \nabla p = \nu \Delta u, \\ \nabla \cdot u = 0, \\ u(x, t = 0) = u_0, \end{cases}$$
(1.1)

where $u(x,t) \in \mathbb{R}^3$ and $p(x,t) \in \mathbb{R}$ denote the unknowns, velocity and pressure respectively, while ν denotes the viscous coefficient of the system.

A lot of works have been devoted to study the above system, but global well-posedness for (1.1) with arbitrary large initial data is still a challenging open problem (see [2,4,9]).

Here, we are concerned with (1.1) with axisymmetric initial data. If u_0 is axisymmetric in system (1.1), then the solution u(x,t) of system (1.1) is also axisymmetric [14,8]. So, it is convenient to write u(x,t) as the following form:

* Corresponding author.



CrossMark

E-mail address: wei.zhiqiang@yahoo.com (Z. Wei).

$$u(x,t) = u^r(r,z,t)e_r + u^\theta(r,z,t)e_\theta + u^z(r,z,t)e_z$$

where e_r , e_{θ} and e_z are the standard orthonormal unit vectors in cylindrical coordinate system

$$e_r = \left(\frac{x_1}{r}, \frac{x_2}{r}, 0\right) = (\cos\theta, \sin\theta, 0),$$

$$e_\theta = \left(-\frac{x_2}{r}, \frac{x_1}{r}, 0\right) = (-\sin\theta, \cos\theta, 0),$$

$$e_z = (0, 0, 1),$$

with $r = (x_1^2 + x_2^2)^{\frac{1}{2}}$.

By direct computation, it is easy to have the following relations.

$$\nabla = (\partial_{x_1}, \partial_{x_2}, \partial_z)^T = \partial_r e_r + \frac{\partial_\theta}{r} e_\theta + \partial_z e_z,$$
$$\Delta = \nabla \cdot \nabla = \frac{1}{r} \partial_r (r \partial_r) + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{\partial^2}{\partial z^2},$$

and

$$\frac{\partial e_r}{\partial \theta} = e_\theta, \qquad \frac{\partial e_\theta}{\partial \theta} = -e_r$$

Accordingly, system (1.1) can be rewritten equivalently as the following form:

$$\begin{cases} \frac{\tilde{D}}{Dt}u^r - \nu(\partial_r^2 + \partial_z^2 + \frac{1}{r}\partial_r - \frac{1}{r^2})u^r - \frac{(u^\theta)^2}{r} + \partial_r p = 0, \\ \frac{\tilde{D}}{Dt}u^\theta - \nu(\partial_r^2 + \partial_z^2 + \frac{1}{r}\partial_r - \frac{1}{r^2})u^\theta + \frac{u^r u^\theta}{r} = 0, \\ \frac{\tilde{D}}{Dt}u^z - \nu(\partial_r^2 + \partial_z^2 + \frac{1}{r}\partial_r)u^z + \partial_z p = 0, \\ u|_{t=0} = u_0^r \cdot e_r + u_0^\theta \cdot e_\theta + u_0^z \cdot e_z, \end{cases}$$
(1.2)

where $\frac{\tilde{D}}{Dt}$ denotes the material derivative

$$\frac{\ddot{D}}{Dt} = \partial_t + u^r \partial_r + u^z \partial_z.$$

If $u^{\theta} = 0$ (so-called without swirl), Ukhovskii and Yudovich [14] (see also [8]) proved the existence of generalized solutions, uniqueness and regularity. When $u^{\theta} \neq 0$ (with swirl), it is very complicated and difficult. For recent progress, one can find results on regularity criteria or global existence with small initial data in [8,10–12,1,8,15]. The question of the regularity criteria to the axisymmetric Navier–Stokes equations with swirl in dependence on the θ -component of the velocity for the first time was studied in [10] and [11]; the former deals with suitable weak solutions and the latter with the Cauchy problem. An important result was proved in [12], where it was shown that if

$$\left\|u^{\theta}\right\|_{L^{\alpha}((0,T);L^{\beta})} < \infty, \text{ with } \frac{2}{\alpha} + \frac{3}{\beta} < 1, \beta > 4,$$

then the solution is smooth up to time T.

The fact that the only singularity may lay on the z-axis follows from [10], so very recently in [15], the following regularity criterion was established:

$$|u^{\theta} 1_{r \le \varsigma}||_{L^{\alpha}((0,T);L^{\beta})} < \infty, \text{ with } \frac{2}{\alpha} + \frac{3}{\beta} = 1, \text{ with } (\alpha, \beta) = (4,6),$$
 (1.3)

where $\varsigma > 0$ is given.

Download English Version:

https://daneshyari.com/en/article/4614434

Download Persian Version:

https://daneshyari.com/article/4614434

Daneshyari.com