



Blow-up criterion for the 3D compressible non-isentropic Navier–Stokes equations without thermal conductivity



Yachun Li^a, Junru Xu^b, Shengguo Zhu^{b,*}

^a Department of Mathematics and Key Lab of Scientific and Engineering Computing (MOE), Shanghai Jiao Tong University, Shanghai 200240, PR China

^b Department of Mathematics, Shanghai Jiao Tong University, Shanghai 200240, PR China

ARTICLE INFO

Article history:

Received 27 February 2015

Available online 12 June 2015

Submitted by D.M. Ambrose

Keywords:

Navier–Stokes

Strong solutions

Vacuum

Blow-up criterion

ABSTRACT

In this paper, we prove a blow-up criterion in terms of the density ρ and the pressure P for the strong solutions with vacuum to Cauchy problem of the 3D compressible non-isentropic Navier–Stokes equations without thermal conductivity. More precisely, we show that the strong solution exists globally if the norm $\|(\rho, P)\|_{L^\infty([0, T] \times \mathbb{R}^3)}$ is bounded from above.

© 2015 Elsevier Inc. All rights reserved.

1. Introduction

The motion of a viscous compressible fluid in \mathbb{R}^3 can be described by the compressible Navier–Stokes equations

$$\begin{cases} \rho_t + \operatorname{div}(\rho u) = 0, \\ (\rho u)_t + \operatorname{div}(\rho u \otimes u) + \nabla P = \operatorname{div} \mathbb{T}, \\ (\rho e)_t + \operatorname{div}(\rho e u) + P \operatorname{div} u - \kappa \Delta \theta = \operatorname{div}(u \mathbb{T}) - u \operatorname{div} \mathbb{T}. \end{cases} \quad (1.1)$$

In this system, $x \in \mathbb{R}^3$ is the spatial coordinate, $t \geq 0$ is the time, ρ is the mass density, $u = (u^1, u^2, u^3) \in \mathbb{R}^3$ is the velocity vector of fluids, e is the specific internal energy, the constant κ is the thermal conductivity coefficient, P is the pressure satisfying

$$P = (\gamma - 1)\rho e = R\rho\theta, \quad (1.2)$$

* Corresponding author.

E-mail addresses: ycli@sjtu.edu.cn (Y.C. Li), 543022824@qq.com (J.R. Xu), zhushengguo@sjtu.edu.cn (S.G. Zhu).

where θ is the absolute temperature, γ is the adiabatic exponent and R is a positive constant; \mathbb{T} is the viscous stress tensor given by

$$\mathbb{T} = 2\mu D(u) + \lambda \operatorname{div} u \mathbb{I}_3, \quad (1.3)$$

where \mathbb{I}_3 is the 3×3 identity matrix, $D(u) = \frac{\nabla u + (\nabla u)^\top}{2}$ is the deformation tensor, μ is the shear viscosity coefficient, $\lambda + \frac{2}{3}\mu$ is the bulk viscosity coefficient, μ and λ are both real constants satisfying

$$\mu > 0, \quad 3\lambda + 2\mu \geq 0, \quad (1.4)$$

which ensure the ellipticity of the Lamé operator L defined by

$$Lu = -\operatorname{div} \mathbb{T} = -\mu \Delta u - (\lambda + \mu) \nabla \operatorname{div} u. \quad (1.5)$$

When $\kappa = 0$, from (1.2), system (1.1) can be written as

$$\begin{cases} \rho_t + \operatorname{div}(\rho u) = 0, \\ (\rho u)_t + \operatorname{div}(\rho u \otimes u) + \nabla P = \operatorname{div} \mathbb{T}, \\ P_t + u \cdot \nabla P + \gamma P \operatorname{div} u = (\gamma - 1)(\operatorname{div}(u \mathbb{T}) - u \operatorname{div} \mathbb{T}). \end{cases} \quad (1.6)$$

This paper is aimed at giving a blow-up criterion of strong solutions to the Cauchy problem of system (1.6) with the following initial data

$$(\rho, u, P)|_{t=0} = (\rho_0(x), u_0(x), P_0(x)), \quad x \in \mathbb{R}^3 \quad (1.7)$$

and the far field behavior

$$(\rho, u, P)(t, x) \rightarrow (0, 0, 0) \quad \text{as } |x| \rightarrow +\infty, \quad t > 0. \quad (1.8)$$

Throughout this paper, we adopt the following simplified notations for the standard homogeneous and inhomogeneous Sobolev space:

$$\begin{aligned} \|f\|_s &= \|f\|_{H^s(\mathbb{R}^3)}, & |f|_p &= \|f\|_{L^p(\mathbb{R}^3)}, & \|(f, g)\|_X &= \|f\|_X + \|g\|_X, \\ D^{k,r} &= \{f \in L^1_{loc}(\mathbb{R}^3) : |\nabla^k f|_r < +\infty\}, & D^k &= D^{k,2}, & |f|_{D^{k,r}} &= \|f\|_{D^{k,r}(\mathbb{R}^3)}, \\ D^1_0 &= \{f \in L^6(\mathbb{R}^3) : |\nabla f|_2 < \infty\}, & |f|_{D^1_0} &= \|f\|_{D^1_0(\mathbb{R}^3)}. \end{aligned}$$

A detailed study on homogeneous Sobolev spaces can be found in [4].

As has been observed in Theorem 3 of Cho–Kim [3], in which the existence of unique local strong solution for system (1.6) was proved, in order to make sure that the Cauchy problem (1.6)–(1.8) with vacuum is well-posed, some compatibility condition on the initial data (ρ_0, u_0, P_0) was proposed to compensate the lack of a positive lower bound of the initial mass density ρ_0 .

Theorem 1.1. (See [3].) *If the initial data (ρ_0, u_0, P_0) satisfy*

$$(\rho_0, P_0) \in H^1 \cap W^{1,q}, \quad \rho_0 \geq 0, \quad P_0 \geq 0, \quad u_0 \in D^1_0 \cap D^2, \quad (1.9)$$

and the compatibility condition

$$Lu_0 + \nabla P_0 = \sqrt{\rho_0} h, \quad \text{for some } h \in L^2, \quad (1.10)$$

Download English Version:

<https://daneshyari.com/en/article/4615110>

Download Persian Version:

<https://daneshyari.com/article/4615110>

[Daneshyari.com](https://daneshyari.com)