# Spiral slits map and its inverse of bounded multiply connected regions 

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#### Abstract

This paper presents a boundary integral equation method for computing numerical conformal mapping of bounded multiply connected region onto a spiral slits region. The method is an extension of the author's method for computing the circular slits map of bounded multiply connected regions (see A.W.K. Sangawi et al. (2012) [16]). Several numerical examples are given to prove the effectiveness of the proposed methods.


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## 1. Introduction

A boundary integral equation method for mapping a bounded multiply connected region onto a circular slits region has been presented in [16]. The classical books [14,1] described the five classical canonical slit regions. The spiral slits region includes as special cases the canonical slits regions considered in Koebe [6], DeLillo et al. [4] and Nasser [11]. This region has been studied for bounded and unbounded multiply connected regions in [11] by reformulating the conformal mapping as a Riemann-Hilbert problem. The main idea of this paper is to modify a previously known construction of conformal mapping to a circular slit region, to produce instead a region whose $j$ th boundary component is a logarithmic spiral

$$
\operatorname{Im} e^{-\mathrm{i} \theta_{j}} \log w=\mathrm{constant},
$$

where the oblique angles $\theta_{j}$ are prescribed in advance. In this paper, we extend the result of [16] by presenting a boundary integral equation method for numerical conformal mappings from a bounded multiply connected region onto the spiral slits region. The proposed method is based on linear boundary integral equation method which is constructed from a boundary relationship satisfied by an analytic function on a bounded multiply connected region.

The plan of the paper is as follows: Section 2 presents some notations and auxiliary materials. Section 3 presents a method to calculate the piecewise real functions $h_{j}$ and $v_{j}$. Section 4 presents a derivation of integral equation related to $S^{\prime}$. In Section 5 , we give some examples to illustrate our boundary integral equation method. Finally, Section 6 presents a short conclusion.

## 2. Notations and auxiliary material

A bounded multiply connected region $\Omega$ of connectivity $M+1$. The boundary $\Gamma$ consists of $M+1$ smooth Jordan curves $\Gamma_{j}, j=0,1, \ldots, M$ as shown in Fig. 1.

[^0]

Fig. 1. Mapping of the bounded multiply connected region onto a spiral slits region.

The curve $\Gamma_{j}$ is parametrized by a $2 \pi$-periodic twice continuously differentiable complex function $z_{j}(t)$ with non-vanishing first derivative

$$
z_{j}^{\prime}(t)=d z_{j}(t) / d t \neq 0, \quad t \in J_{j}=[0,2 \pi], \quad j=0,1, \ldots, M .
$$

The total parameter domain $J$ is the disjoint union of $M+1$ intervals $J_{0}, \ldots, J_{M}$. We define a parametrization $z(t)$ of the whole boundary $\Gamma$ on $J$ by

$$
\begin{equation*}
z(t)=z_{j}(t), \quad t \in J_{j}, \quad j=0,1, \ldots, M \tag{1}
\end{equation*}
$$

The unknown functions $S(t)$ and $R(t)$ (which will be defined in Section 4) will be given for $t \in J$ by

$$
S(t)=S_{j}(t) \text { and } R(t)=R_{j}(t), \quad t \in J_{j}, j=0,1, \ldots, M
$$

Suppose that $c(z), Q(z)$ and $H(z)$ are complex-valued functions defined on $\Gamma$ such that $c(z) \neq 0, H(z) \neq 0, Q(z) \neq 0$ and $\overline{H(z)} /(T(z) Q(z))$ satisfies the Hölder condition on $\Gamma$. Then the interior relationship is defined as follows:

A complex-valued function $P(z)$ is said to satisfy the interior relationship if $\mathrm{P}(\mathrm{z})$ is analytic in $\Omega$ and satisfies the non-homogeneous boundary relationship

$$
\begin{equation*}
P(z)=c(z) \frac{\overline{T(z) Q(z)}}{\overline{G(z)}} \overline{P(z)}+\overline{H(z)}, \quad z \in \Gamma \tag{2}
\end{equation*}
$$

where $G(z)$ analytic in $\Omega$, Hölder continuous on $\Gamma$, and $G(z) \neq 0$ on $\Gamma$. The boundary relationship (2) also has the following equivalent form:

$$
\begin{equation*}
G(z)=\overline{c(z)} T(z) Q(z) \frac{P(z)^{2}}{|P(z)|^{2}}+\frac{G(z) H(z)}{\overline{P(z)}}, \quad z \in \Gamma \tag{3}
\end{equation*}
$$

The following theorem gives an integral equation for an analytic function satisfying the interior non-homogeneous boundary relationship (2) or (3), [16].

Theorem 1. Let $U$ and $V$ be any complex-valued functions that are defined on $\Gamma$. If the function $P(z)$ satisfies the interior nonhomogeneous boundary relationship (2) or (3), then

$$
\begin{equation*}
\frac{1}{2}\left[V(z)+\frac{U(z)}{\overline{T(z) Q(z)}}\right] P(z)+\mathrm{PV} \int_{\Gamma} K(z, w) P(w)|d w|+c(z) U(z)\left[\sum_{a_{j} \in \Omega} \operatorname{Res}_{w=a_{j}} \frac{P(w)}{(w-z) G(w)}\right]^{\text {conj }}=-U(z) \overline{L_{R}^{-}(z)}, \quad z \in \Gamma \tag{4}
\end{equation*}
$$

where

$$
\begin{align*}
& K(z, w)=\frac{1}{2 \pi \mathrm{i}}\left[\frac{c(z) U(z)}{c(w)(\bar{w}-\bar{z}) \overline{Q(w)}}-\frac{V(z) T(w)}{w-z}\right]  \tag{5}\\
& L_{R}^{-}(z)=\frac{-1}{2} \frac{H(z)}{Q(z) T(z)}+\mathrm{PV} \frac{1}{2 \pi \mathrm{i}} \int_{\Gamma} \frac{\overline{c(z)} H(w)}{\overline{c(w)}(w-z) Q(w) T(w)} d w . \tag{6}
\end{align*}
$$

The symbol "conj" in the superscript denotes complex conjugate, while the minus sign in the superscript denotes limit from the exterior.

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