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# A centerless representation of the Virasoro algebra associated with the unitary circular ensemble

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#### a r t i c l e i n f o

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## A B S T R A C T

We consider the 2-dimensional Toda lattice tau functions  $\tau_n(t, s; \eta, \theta)$  deforming the probabilities  $\tau_n(\eta, \theta)$  that a randomly chosen matrix from the unitary group  $U(n)$ , for the Haar measure, has no eigenvalues within an arc  $(\eta, \theta)$  of the unit circle. We show that these tau functions satisfy a centerless Virasoro algebra of constraints, with a boundary part in the sense of Adler, Shiota and van Moerbeke. As an application, we obtain a new derivation of a differential equation due to Tracy and Widom, satisfied by these probabilities, linking it to the Painlevé VI equation.

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## **1. Introduction**

Consider the group  $U(n)$  of  $n \times n$  unitary matrices, with the normalized Haar measure as a probability measure. The Weyl integral formula gives the induced density distribution on the eigenvalues of the matrices on the unit circle in the complex plane, and is given by

$$
\frac{1}{n!}|\Delta_n(z)|^2\prod_{k=1}^n\frac{dz_k}{2\pi iz_k};\quad z_k=e^{i\varphi_k}\quad\text{and}\quad\Delta_n(z)=\prod_{1\leq k
$$

Thus, for  $\eta$ ,  $\theta \in ]-\pi, \pi[$ , with  $\eta \leq \theta$ , the probability that a randomly chosen matrix from  $U(n)$  has no eigenvalues within an arc of circle  $(\eta, \theta) = \{z \in S^1 | \eta < \arg(z) < \theta\}$  is given by

$$
\tau_n(\eta,\theta)=\frac{1}{(2\pi)^n n!}\int_{\theta}^{2\pi+\eta}\cdots\int_{\theta}^{2\pi+\eta}\prod_{1\leq k
$$

Obviously, this probability depends only on the length  $\theta - \eta$ . All of this is well known and we refer the reader to [\[1\]](#page--1-0) for details. We shall denote by

$$
R(\theta) = -\frac{1}{2} \frac{d}{d\theta} \log \tau_n(-\theta, \theta), \tag{1.1}
$$

the logarithmic derivative of the probability that an arc of circle of length  $2\theta$  contains no eigenvalues of a randomly chosen unitary matrix.

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The starting motivation for our work was to understand a differential equation satisfied by the function  $R(\theta)$ 

<span id="page-1-1"></span>
$$
R(\theta)^{2} + 2 \sin \theta \cos \theta R(\theta) R'(\theta) + \sin^{2} \theta R'(\theta)^{2}
$$
  
=  $\frac{1}{2} \left( \frac{1}{4} \sin^{2} \theta \frac{R''(\theta)^{2}}{R'(\theta)} + \sin \theta \cos \theta R''(\theta) + (\cos^{2} \theta + n^{2} \sin^{2} \theta) R'(\theta) \right),$  (1.2)

obtained in [\[2\]](#page--1-1), from the point of view of the Adler–Shiota–van Moerbeke [\[3\]](#page--1-2) approach, in terms of Virasoro constraints. Introducing the 2-Toda time-dependent tau functions

<span id="page-1-2"></span>
$$
\tau_n(t,s;\,\eta,\theta) = \frac{1}{n!} \int_{[\theta,2\pi+\eta]^n} |\Delta_n(z)|^2 \prod_{k=1}^n \left( e^{j\sum\limits_{j=1}^n (t_j z_k^j + s_j z_k^{-j})} \frac{dz_k}{2\pi i z_k} \right),\tag{1.3}
$$

with  $z_k = e^{i\varphi_k}$ , deforming the probabilities  $\tau_n(\eta,\theta) = \tau_n(0,0;\eta,\theta)$ , we discover that they satisfy a set of Virasoro constraints indexed by *all* integers, decoupling into a boundary-part and a time-part

$$
\frac{1}{i}\left(e^{ik\theta} \frac{\partial}{\partial \theta} + e^{ik\eta} \frac{\partial}{\partial \eta}\right) \tau_n(t,s;\eta,\theta) = L_k^{(n)} \tau_n(t,s;\eta,\theta), \quad k \in \mathbb{Z}, i = \sqrt{-1},
$$

with the time-dependent operators  $L_k^{(n)}$  providing a centerless representation of the *full* Virasoro algebra, see Section [2](#page-1-0) [\(Theorem 2.2\)](#page--1-3) for a precise statement and the proof of the result.

In their study of Painlevé equations satisfied (as functions of *x*) by integrals of Gessel's type *EU*(*n*)e *x* tr(*M*+*M*) , where the expectation  $E_{U(n)}$  refers to integration with respect to the Haar measure over the whole of  $U(n)$ , Adler and van Moerbeke [\[4\]](#page--1-4) found the  $s_l$ <sub>2</sub> subalgebra corresponding to  $k = -1, 0, 1$ , without boundary terms. The appearance of boundary terms and of a *full* centerless Virasoro algebra is to the best of our knowledge new. From this result, it is easy to obtain Eq. [\(1.2\),](#page-1-1) using the algorithmic method of [\[3\]](#page--1-2). Finally, similarly to a result by the first author and Semengue [\[5\]](#page--1-5) on the Jacobi polynomial ensemble, we show that  $R(\theta)$  is the restriction to the unit circle of a function  $r(z)$  defined in the complex plane, so that  $\sigma(z) = -i(z-1)r(z) - n^2z/4$  satisfies a special case of the Okamoto–Jimbo–Miwa form of the Painlevé VI equation. This will be explained in Section [3](#page--1-6) of the paper.

## <span id="page-1-0"></span>**2. A centerless representation of the Virasoro algebra**

The proof of the Virasoro constraints satisfied by the integral [\(1.3\)](#page-1-2) is a non-trivial adaptation of the self-similarity argument exploited in the case of the Gaussian ensembles, based on the invariance of the integrals with respect to translations, see [\[6\]](#page--1-7) and references therein. Here, we replace translations by appropriate rotations. More precisely, setting

$$
dl_n(t,s,z) = |\Delta_n(z)|^2 \prod_{\alpha=1}^n \left( e^{\sum_{j=1}^\infty (t_j z_\alpha^j + s_j z_\alpha^{-j})} \frac{dz_\alpha}{2\pi i z_\alpha} \right),\tag{2.1}
$$

with  $z_\alpha = e^{i\varphi_\alpha}$  and  $|\Delta_n(z)|^2 = \prod_{1 \le \alpha < \beta \le n} |z_\alpha - z_\beta|^2$ , we have the fundamental next proposition.

**Proposition 2.1.** *The following variational formulas hold*

$$
\frac{\mathrm{d}}{\mathrm{d}\varepsilon}\mathrm{d}I_n\left(z_\alpha\mapsto z_\alpha\mathrm{e}^{\varepsilon(z_\alpha^k-z_\alpha^{-k})}\right)\Big|_{\varepsilon=0}=\left(L_k^{(n)}-L_{-k}^{(n)}\right)\mathrm{d}I_n,\tag{2.2}
$$

$$
\frac{\mathrm{d}}{\mathrm{d}\varepsilon}\mathrm{d}I_n\left(z_\alpha \mapsto z_\alpha \mathrm{e}^{\mathrm{i}\varepsilon (z_\alpha^k + z_\alpha^{-k})}\right)\Big|_{\varepsilon=0} = \mathrm{i}\left(L_k^{(n)} + L_{-k}^{(n)}\right) \mathrm{d}I_n,\tag{2.3}
$$

*for all*  $k > 0$ *, with* 

$$
L_k^{(n)} = \sum_{j=1}^{k-1} \frac{\partial^2}{\partial t_j \partial t_{k-j}} + n \frac{\partial}{\partial t_k} + \sum_{j=1}^{\infty} j t_j \frac{\partial}{\partial t_{j+k}} - \sum_{j=k+1}^{\infty} j s_j \frac{\partial}{\partial s_{j-k}} - \sum_{j=1}^{k-1} j s_j \frac{\partial}{\partial t_{k-j}} - n k s_k, \quad k \ge 1,
$$
\n(2.4)

$$
L_0^{(n)} = \sum_{j=1}^{\infty} j t_j \frac{\partial}{\partial t_j} - \sum_{j=1}^{\infty} j s_j \frac{\partial}{\partial s_j},\tag{2.5}
$$

$$
L_{-k}^{(n)} = -\sum_{j=1}^{k-1} \frac{\partial^2}{\partial s_j \partial s_{k-j}} - n \frac{\partial}{\partial s_k} - \sum_{j=1}^{\infty} j s_j \frac{\partial}{\partial s_{j+k}} + \sum_{j=k+1}^{\infty} j t_j \frac{\partial}{\partial t_{j-k}} + \sum_{j=1}^{k-1} j t_j \frac{\partial}{\partial s_{k-j}} + n k t_k, \quad k \ge 1.
$$
 (2.6)

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