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## Conjugate symplecticity of second-order linear multi-step methods

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#### Abstract

We review the two different approaches for symplecticity of linear multi-step methods (LMSM) by Eirola and Sanz-Serna, Ge and Feng, and by Feng and Tang, Hairer and Leone, respectively, and give a numerical example between these two approaches. We prove that in the conjugate relation  $G_3^{\lambda\tau} \circ G_1^{\tau} = G_2^{\tau} \circ G_3^{\lambda\tau}$  with  $G_1^{\tau}$  and  $G_3^{\tau}$  being LMSMs, if  $G_2^{\tau}$  is symplectic, then the *B*-series error expansions of  $G_1^{\tau}$ ,  $G_2^{\tau}$  and  $G_3^{\tau}$  of the form  $G^{\tau}(Z) = \sum_{i=0}^{+\infty} (\tau^i/i!) Z^{[i]} + \tau^{s+1} A_1 + \tau^{s+2} A_2 + \tau^{s+3} A_3 + \tau^{s+4} A_4 + \mathrm{O}(\tau^{s+5})$  are equal to those of trapezoid, mid-point and Euler forward schemes up to a parameter  $\theta$  (completely the same when  $\theta = 1$ ), respectively, this also partially solves a problem due to Hairer. In particular we indicate that the second-order symmetric leap-frog scheme  $Z_2 = Z_0 + 2\tau J^{-1} \nabla H(Z_1)$  cannot be conjugate-symplectic via another LMSM. © 2006 Elsevier B.V. All rights reserved.

Keywords: Linear multi-step method; Step-transition operator; B-series; Conjugate relation; Symplecticity

#### 1. First approach for symplectic multi-step methods

It is well-known that for a Hamiltonian system

$$\frac{\mathrm{d}Z}{\mathrm{d}t} = J^{-1}\nabla H(Z), \quad Z = [z_1, \dots, z_{2n}]^{\top} \in \mathbb{R}^{2n},\tag{1}$$

where  $J = [J_{ij}] = \begin{bmatrix} O_n & I_n \\ -I_n & O_n \end{bmatrix}$ ,  $\nabla$  stands for the gradient operator, and  $H : \mathbb{R}^{2n} \to \mathbb{R}^1$  is a smooth function (*Hamiltonian function*), its phase flow  $\{g^t | t \in \mathbb{R}\}$  is a one-parameter group of symplectic transformations [1]. The symplecticity of  $g^t : \mathbb{R}^{2n} \to \mathbb{R}^{2n}$  means

$$\left[ \frac{\partial g^t(Z)}{\partial Z} \right]^{\top} J \left[ \frac{\partial g^t(Z)}{\partial Z} \right] = J$$
 (2)

for any  $Z \in \mathbb{R}^{2n}$  and any  $t \in \mathbb{R}$ .

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It is easy to check that Eq. (2) is equivalent to [6]

$$(g^t)^*\omega = \omega, \tag{3}$$

where  $\omega = \frac{1}{2} \sum_{a,b=1}^{2n} J_{ab} \, dz_a \wedge dz_b = \sum_{1 \leqslant a < b \leqslant 2n} J_{ab} \, dz_a \wedge dz_b = \sum_{c=1}^n dz_c \wedge dz_{n+c}$ . More generally, if J becomes K(Z) where  $K(Z) = (K_{ab})$  is an antisymmetric, nondegenerate  $2n \times 2n$  matrix

$$\frac{\partial k_{ab}}{\partial z_c} + \frac{\partial k_{bc}}{\partial z_a} + \frac{\partial k_{ca}}{\partial z_b} = 0, \quad 1 \leqslant a, b, c \leqslant 2n, \tag{4}$$

then (1) becomes the general Hamiltonian system

$$\frac{\mathrm{d}Z}{\mathrm{d}t} = K^{-1}(Z)\nabla H(Z), \quad Z = [z_1, \dots, z_{2n}]^{\top} \in \mathbb{R}^{2n},\tag{5}$$

and the phase flow  $\{\hat{g}^t | t \in \mathbb{R}\}\$  of (5) is a one-parameter group of K-symplectic transformations [6,9]:

$$\left[\frac{\partial \hat{g}^{t}(Z)}{\partial Z}\right]^{T} K(\hat{g}^{t}(Z)) \left[\frac{\partial \hat{g}^{t}(Z)}{\partial Z}\right] = K(Z). \tag{6}$$

Furthermore Eq. (6) is equivalent to

$$(\hat{g}^t)^* \hat{\omega} = \hat{\omega},\tag{7}$$

where  $\hat{\omega} = \frac{1}{2} \sum_{a,b=1}^{2n} K_{ab} dz_a \wedge dz_b = \sum_{1 \leqslant a < b \leqslant 2n} K_{ab} dz_a \wedge dz_b$ .

A numerical scheme compatible with (5) is said to be *K*-symplectic if its step-transition operator  $G^{\tau}: \mathbb{R}^{2n} \to \mathbb{R}^{2n}$ is a K-symplectic transformation for any stepsize  $\tau$ . In particular, the mid-point rule  $G_{mp}^{\tau}: Z \to \widetilde{Z}$  (see [6])

$$Z_1 - Z_0 = \tau J^{-1} \nabla H \left( \frac{Z_1 + Z_0}{2} \right) \tag{8}$$

is a second-order symplectic scheme for the standard Hamiltonian system (1).

The symplecticity of compatible linear m-step method (LMSM) for Hamiltonian system (1)

$$\sum_{k=0}^{m} \alpha_k Z_k = \tau \sum_{k=0}^{m} \beta_k J^{-1} \nabla H(Z_k), \quad \left(\sum_{k=0}^{m} \beta_k \neq 0\right)$$
 (9)

is first studied under the consideration of transformations in the higher dimensional manifold  $\mathbb{R}^{2mn}$ .

For the special case m = 2 for example, for the second-order leap-frog scheme

$$Z_2 = Z_0 + 2\tau J^{-1} \nabla H(Z_1), \tag{10}$$

Ge and Feng [11] rewrote (10) into

and showed that the mapping  $[Z_1^\top, Z_0^\top]^\top \xrightarrow{\Phi} [Z_2^\top, Z_1^\top]^\top$  preserves the general symplectic structure related to  $\begin{bmatrix} O_{2n} & J_{2n} \\ J_{2n} & O_{2n} \end{bmatrix}$ . More generally, Eirola and Sanz-Serna [5] have shown that if one-leg method (see [13,16] for details)

$$\sum_{k=0}^{m} \alpha_k Z_k = \tau J^{-1} \nabla H \left( \sum_{k=0}^{m} \beta_k Z_k \right)$$
 (12)

is symmetric (i.e.,  $\alpha_{m-k} = -\alpha_k$ ,  $\beta_{m-k} = \beta_k$ ,  $0 \le k \le m$ ) and irreducible, then the transformation  $(Z_0^\top, \dots, Z_{m-1}^\top)^\top \longrightarrow (Z_1^\top, \dots, Z_m^\top)^\top$  in the higher dimensional manifold  $\mathbb{R}^{2mn}$  is symplectic with respect to the general structure  $\Lambda \otimes J$ , where  $\Lambda$  is an  $m \times m$  symmetric matrix defined by the coefficients  $\alpha_k$ ,  $\beta_k$ ,  $0 \le k \le m$ , it is  $\begin{bmatrix} 0 & 2 \\ 2 & 0 \end{bmatrix}$  for the leap-frog scheme (10).

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