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Influence of oxygen fugacity on the electrical conductivity of hydrous olivine: Implications for the mechanism of conduction



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ABSTRACT

Various mechanisms of hydrogen-assisted electrical conductivity predict different dependence of conductivity on oxygen fugacity. If a majority of hydrogen-related defects (i.e., two protons at M-site) carries the electrical current, then the conductivity will be independent of oxygen fugacity, whereas if a minority defect such as free proton carries the electrical conductivity, then the electrical conductivity will decrease with oxygen fugacity. We have determined the dependence of hydrogen-assisted electrical conductivity in olivine decreases with oxygen fugacity. This result supports a model where hydrogen-related defect with minor concentration (e.g., free proton or one hydrogen at M-site) carries most of the electric charge.

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1. Introduction

Importance of hydrogen in electrical conductivity in minerals such as olivine was pointed out by (Karato, 1990). This model was motivated by the reported high diffusion coefficient of hydrogen (Mackwell and Kohlstedt, 1990) and known appreciable solubility of hydrogen. Recent experimental studies have largely confirmed this model (e.g., Karato and Wang, 2013; Yoshino and Katsura, 2013), but some details are still debated.

One of the most important observations is the large discrepancies between the results of H-D isotope diffusion (Du Frane and Tyburczy, 2012) and electrical conductivity. (Karato, 2013) developed a theory of isotope diffusion and discussed that such discrepancies can be explained if electrical conductivity is due mostly to defects with small concentration (e.g., free proton) but high mobility whereas isotope diffusion is controlled by the diffusion of defects slow diffusing species. In contrast to this hybrid model of electrical conductivity, most of other researchers propose a model where all hydrogen atoms contribute equally to electrical conductivity (e.g., Du Frane and Tyburczy, 2012; Yoshino and Katsura, 2013). Although detailed discussions that favor the hybrid model has been published (e.g., Karato, 2006, 2013), it is useful to provide an additional experimental study to test these models.

In this paper, we present the results of an experimental study to determine the dependence of electrical conductivity of olivine on oxygen fugacity under hydrogen-rich conditions. Many experimental studies on electrical conductivity of dry olivine showed that electrical conductivity increases with oxygen fugacity (e.g., Duba and Nicholls, 1973; Schock et al., 1989). This is due to the fact that the concentration of the charge carrying species, i.e., ferric iron, increases with oxygen fugacity. However, the situation can be different when the charge carrier is proton-related defects. Indeed, a defect chemistry analysis on olivine (or other (Mg, Fe), Si bearing minerals) shows that the dependence of conductivity on oxygen fugacity will be different for hydrogen-assisted conductivity. In more detail, if hydrogen-assisted conductivity were due to the most abundant hydrogen-related defects, i.e., $(2H)_M^{\times}$ (two protons at M-site vacancy), then electrical conductivity will be independent of oxygen fugacity, whereas if electrical conductivity is due to minority defects such as H^{\bullet} (free proton) or H'_{M} (one proton at M-site vacancy), then electrical conductivity would decrease with oxygen fugacity (Nishihara et al., 2008). Therefore a study on the influence of oxygen fugacity will provide a clear means to distinguish among various mechanisms of hydrogen-assisted electrical conduction.

2. Experimental procedure

2.1. Sample preparation

The hydrous polycrystalline olivine samples were synthesized from San Carlos olivine powders with the grain size less than 5 μ m by hydrothermal annealing experiments at *P* = 4 GPa and

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T = 1473 K for 3 h. About 1.85 wt% of San Carlos orthopyroxene was added to buffer the oxide activity. Powder samples of olivine and orthopyroxene mixture were sealed into a metal capsule (water was added by the dehydration of talc + brucite). Both synthesis and conductivity measurements were performed using the same solid oxygen buffer (e.g., Mo-MoO₂, Ni-NiO and Re-ReO₂). Our previous experience indicates that under these environments, the oxidation conditions of samples are well buffered by the metal and its corresponding metal oxide reactions (e.g., Nishihara et al., 2006, 2008). The oxygen fugacity was controlled by varying the selected metal type in the electrodes, corresponding sleeves, and shielding cases. The shielding is made of metal foils such as Mo, Ni, Re, whose thickness is 25 µm. After the electrical conductivity measurements were completed, the recovered electrode/sample interfaces were examined by the scanning electron microscope in order to confirm the coexistence of metal and its corresponding metal oxide. We confirmed the co-existence of metal and corresponding metal oxide at the interface between electrode and the samples.

2.2. Sample characterization

In order to determine the water content of the samples, the infrared spectra of samples were obtained at wavenumbers from 1000 to 4000 cm⁻¹ both before and after each experiment. The measurements were made using a Fourier transform infrared spectroscopy (FT-IR) spectrometer (BIORAD, Varian 600 UMA). Doubly polished samples with a thickness of ~100 μ m were prepared for the IR analysis. The IR absorption of samples was measured by unpolarized radiation with a mid-IR light source, a KBr beam splitter and an MCT detector with a 150 × 150 μ m² aperture. 256 scans were accumulated for each sample. The infrared spectra of the acquired samples are shown in Fig. 1. We used the (Paterson, 1982) calibration to determine the water content from FT-IR absorption using,

$$C_{\rm OH} = \frac{B_i}{150\xi} \int \frac{K(\nu)}{(3780 - \nu)} d\nu$$
(1)

where $C_{\rm OH}$ is the molar concentration of hydroxyl (ppm wt H₂O or H/10⁶ Si), B_i is the density factor (4.39 × 10⁴ cm H/10⁶ Si), ξ is the orientation factor (1/3), and K(v) is the absorption coefficient in



Fig. 1. The representative FT-IR spectra of the recovered olivine aggregates after electrical conductivity measurements for the wavenumber range of $3000-4000 \text{ cm}^{-1}$.

cm⁻¹ at wavenumber v in cm⁻¹. The integration was made from 3000 to 3750 cm⁻¹. If another calibration such as (Bell et al., 2003) is used, the water content will be larger by a factor of ~3.

The water content of the original sample is less than $8H/10^6$ Si (0.0001 wt% H₂O). After the olivine was annealed under the hydrous environment, one sample containing different water contents of ~180–280 ppm wt was obtained. The water loss during the electrical conductivity measurements of hydrous sample was less than 8% of the total water.

2.3. Impedance spectroscopy measurements

The experimental sample assembly is shown in Fig. 2. Pressure was generated by eight cubic tungsten carbide anvils $(26 \times 26 \times 26 \text{ mm}^3)$ with an 18 mm truncated edge length. Pressure calibrations were conducted using the phase transitions of coesite to stishovite (Zhang et al., 1996) (9.5 GPa and 1573 K). In order to avoid the influence of adsorbed water on the measurement of electrical conductivity, sample assembly parts including MgO octahedral pressure medium, MgO and Al₂O₃ insulation tubes were heated to 1223 K for 15 h prior to each experiment (Mo ring and a sample itself were not heat treated). In order to control the oxygen fugacity of the sample chamber and reduce the leakage currents, a metal foil shield (e.g., Mo, Ni and Re) was placed between a sample and an MgO insulation tube. The metal used for shielding purpose is the same as the metal used for a capsule. A disk-shaped sample (Φ 1.6 × 0.4 mm) was placed between two parallel metal electrodes that were surrounded by alumina rings. The temperature was measured by a $W_{5\%}$ Re– $W_{26\%}$ Re thermocouple that is attached to another side of metal electrode (metal electrode is in direct contact with a sample, but the thermo-couple is not). The experimental errors of the pressure and temperature gradient were estimated to be no more than 0.5 GPa and 10 K, respectively (absolute error in pressure estimate can be larger but the relative error is \sim 0.5 GPa or less). The errors in the electrical conductivity measurement through the impedance fitting were estimated to be less than 4%.

The pressure was first raised at the rate of ~0.9 GPa/h to a designated pressure. Under a constant pressure condition, temperature was raised at the rate of ~50 K/min to the preset value and the impedance spectroscopy measurements were performed at various temperatures. After the temperature reached to an each value for a constant pressure condition, the ZPlot program of a Solartron-1260 Impedance/Gain-phase analyzer was run to determine the complex impedance for the frequency range of $f = 10^{-2}-10^{6}$ Hz for a sinusoidal alternating current of signal voltage of 1.0 V. The impedance semi-circle arc of high frequency



Fig. 2. The experimental setup for electrical conductivity measurements at high pressure and temperature.

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