



Egyptian Mathematical Society
Journal of the Egyptian Mathematical Society

www.etms-eg.org
www.elsevier.com/locate/joems



ORIGINAL ARTICLE

Slip effects on a generalized Burgers' fluid flow between two side walls with fractional derivative



Shihao Han ^a, Liancun Zheng ^{a,*}, Xinxin Zhang ^b

^a School of Mathematics and Physics, University of Science and Technology Beijing, 30 Xueyuan Road, Haidian District, Beijing 100083, China

^b School of Mechanical Engineering, University of Science and Technology Beijing, 30 Xueyuan Road, Haidian District, Beijing 100083, China

Received 10 July 2014; accepted 9 October 2014

Available online 11 February 2015

KEYWORDS

Generalized Burgers' fluid;
Fractional calculus;
Pressure gradient;
Slip boundary

Abstract This paper presents a research for the 3D flow of a generalized Burgers' fluid between two side walls generated by an exponential accelerating plate and a constant pressure gradient, where the no-slip assumption between the exponential accelerating plate and the Burgers' fluid is no longer valid. The governing equations of the generalized Burgers' fluid flow are established by using the fractional calculus approach. Exact analytic solutions for the 3D flow are established by employing the Laplace transform and the finite Fourier sine transform. Furthermore, some 3D and 2D figures for the fluid velocity and shear stress are plotted to analyze and discuss the effects of various parameters.

2000 MATHEMATICAL SUBJECT CLASSIFICATION: 02.30.Jr; 02.30.Uu; 04.20.Jb

© 2015 Production and hosting by Elsevier B.V. on behalf of Egyptian Mathematical Society.

1. Introduction

Studies of non-Newtonian fluids have obtained more and more attention in the last few years, the main reason may be that non-Newtonian fluids (such as humans blood, polymer suspension, slurries, and oil) are widely exist in life, production and nature. It is well known that the fluid which satisfies the law

of Newton inner friction that the stress tensor is proportional to fluid velocity gradient, is called Newtonian fluid. However non-Newtonian fluids do not obey the law of Newton inner friction and show more complex rheological behaviors than Newtonian fluids. Among of all non-Newtonian fluids, viscoelastic fluid which has both viscous and elastic characteristics, is an important class. Various viscoelastic fluid models have been put forward to research the flow behaviors of the fluid. Among these models, the Maxwell fluid model [1,2], the second grade fluid model [3,4], the Oldroyd-B fluid model [5,6] and the Burgers' fluid model [7–9] are four typical viscoelastic fluid models. The constitutive equations of viscoelastic fluid with fractional derivative have been successfully used to describe viscoelastic characteristics, which are derived by replacing

* Corresponding author. Tel.: +86 1062332002.

E-mail address: liancunzheng@ustb.edu.cn (L. Zheng).

Peer review under responsibility of Egyptian Mathematical Society.



Production and hosting by Elsevier

the integer order time derivative of the classical constitutive equations with the fractional derivative [10–13]. Last few years, many attempts about the generalized Burgers' fluid flow with fractional derivative have been done. Some flows of generalized Burgers' fluid with fractional derivative in the porous media space were considered by Xue [14,15], Khan and Hayat [16] and Hayat et al. [17]. Khan [18–20] obtained the exact analytical solutions for the accelerated, rotating and oscillating flows of a fractional generalized Burgers' fluid. Hyder Ali Muttaqi Shah [21] discussed the Poiseuille and Couette flows of a fractional Burgers' fluid between two parallel plates. Liu et al. [22] investigated the radiation effects on the heat transfer of a fractional generalized Burgers' fluid and obtained the exact analytical solutions for velocity and temperature fields.

In the past few years, unsteady flow problems of viscoelastic fluids between two side walls have obtained considerable attention. The exact analytical solutions for the Maxwell fluid flow between two side walls due to a constant velocity plate was investigated by Hayat et al. [23]. Fetecau et al. [24] and Khan and Wang [25] investigated the flows of a generalized second-grade fluid induced by a accelerated plate between two side walls. Some exact analytical solutions for generalized Oldroyd-B fluid flows between two side walls due to a accelerated plate were established by Fetecau [26,27] and Hyder Ali Muttaqi Shah [28]. Zheng et al. [29] for the first time presented the 3D figures for the flow of a generalized Oldroyd-B fluid with fractional derivative generated by a constant pressure gradient. Moreover, boundary slip is found to play an important role in various engineering technological applications. Zheng et al. [30] firstly studied slip effects on magnetohydrodynamic flow of a generalized Oldroyd-B fluid and obtained the exact analytical solutions.

Based on the above mentioned works and in order to better describe the flow of generalized Burgers' fluid, we investigate slip effects on the 3D flow between two side walls generated by an exponential accelerating plate and a constant pressure gradient. The governing equations of the generalized Burgers' fluid are established by making use of the fractional calculus approach. The exact analytical solutions for the 3D flow are calculated by employing the Laplace transform and the finite Fourier sine transform. Furthermore, some figures are plotted to analyze and discuss the effects of various parameters.

2. Basic governing equations

The constitutive relationship for an incompressible generalized Burgers' fluid with fractional derivative is [16–22]

$$\mathbf{T} = -p\mathbf{I} + \mathbf{S}, \quad (1 + \lambda_1^z \tilde{D}_t^z + \lambda_2^z \tilde{D}_t^{2z})\mathbf{S} = \mu\lambda_3^\beta \tilde{D}_t^\beta \mathbf{A} \quad (1)$$

In the above equations, $-p\mathbf{I}$ indicates the indeterminate spherical stress, \mathbf{T} is the Cauchy stress tensor, \mathbf{S} the extra stress tensor, \mathbf{A} represents the first Rivlin–Ericksen tensor, μ indicates dynamic viscosity, λ_1 the relaxation time, λ_2 is the new material parameter, $\lambda_3 (< \lambda_1)$ the retardation time, α and β ($0 \leq \alpha \leq \beta \leq 1$) are the fractional calculus parameters, \tilde{D}_t^α denotes the upper convected time fractional derivative and is defined by [16–22]

$$\tilde{D}_t^\alpha \mathbf{S} = D_t^\alpha \mathbf{S} + (\mathbf{V} \cdot \nabla) \mathbf{S} - \mathbf{L} \mathbf{S} - \mathbf{S} \mathbf{L}^T \quad (2)$$

$$\tilde{D}_t^{2\alpha} \mathbf{S} = \tilde{D}_t^\alpha (\tilde{D}_t^\alpha \mathbf{S}) \quad (3)$$

where \mathbf{V} denotes the fluid velocity, \mathbf{L} denotes the velocity gradient, D_t^p denotes the fractional calculus operator of order p with respect to t and is defined by [31]

$$D_t^p f(t) = \frac{1}{\Gamma(1-p)} \frac{d}{dt} \int_0^t \frac{f(\tau)}{(t-\tau)^p} d\tau, \quad 0 \leq p \leq 1, \quad (4)$$

where $\Gamma(\cdot)$ is the Gamma function.

The motion equation in the absence of body force is the following form [18]

$$\rho \frac{d\mathbf{V}}{dt} = \nabla \cdot \mathbf{T} \quad (5)$$

where ρ denotes the fluid density, ∇ is the gradient operator.

We assume the motion of the fluid is unidirectional, the velocity field of the fluid may be

$$\mathbf{V} = [u(y, z, t), 0, 0] \quad (6)$$

where $u(y, z, t)$ denotes the fluid velocity in the x -axis direction. Substituting Eq. (6) into (1), (5) and considering initial condition $\mathbf{S}(y, z, 0) = 0$, we obtain the following equations

$$\rho \frac{\partial u}{\partial t} = -\frac{\partial p}{\partial x} + \frac{\partial S_{xy}}{\partial y} + \frac{\partial S_{xz}}{\partial z} \quad (7)$$

$$(1 + \lambda_1^z D_t^z + \lambda_2^z D_t^{2z}) S_{xy} = \mu(1 + \lambda_3^\beta D_t^\beta) \frac{\partial u}{\partial y} \quad (8)$$

$$(1 + \lambda_1^z D_t^z + \lambda_2^z D_t^{2z}) S_{xz} = \mu(1 + \lambda_3^\beta D_t^\beta) \frac{\partial u}{\partial z} \quad (9)$$

where $\partial p / \partial x$ is the pressure gradient along x -axis. Eliminating S_{xy} and S_{xz} between Eqs. (7)–(9), we get the governing equation of the flow problem

$$(1 + \lambda_1^z D_t^z + \lambda_2^z D_t^{2z}) \frac{\partial u(y, z, t)}{\partial t} = -\frac{1}{\rho} (1 + \lambda_1^z D_t^z + \lambda_2^z D_t^{2z}) \frac{\partial p}{\partial x} + \nu(1 + \lambda_3^\beta D_t^\beta) \left(\frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) u(y, z, t) \quad (10)$$

where $\nu = \mu / \rho$ denotes kinematic viscosity.

3. Flow between two side walls

In the following sections, we investigate an incompressible generalized Burgers' fluid with fractional derivative between two side walls occupying the full space above the plate which is perpendicular to the side walls. The Burgers' fluid begins to move due to the exponential accelerating plate with a motion of the speed e^{-t} and a constant pressure gradient $A = -1/\rho \cdot \partial p / \partial x$ in the x -axis direction. The slip between the exponential accelerating plate and the Burgers' fluid is considered. The governing equation of the fluid is Eq. (10) and the corresponding initial and boundary conditions are

$$u(y, z, 0) = u_t(y, z, 0) = u_{tt}(y, z, 0) = 0, \quad (y > 0, 0 \leq z \leq h) \quad (11)$$

$$u(0, z, t) = e^{-t} + \theta \frac{\partial u(0, z, t)}{\partial y}, \quad (t \geq 0, 0 \leq z \leq h) \quad (12)$$

$$u(y, z, t) = u_y(y, z, t) = 0, \quad (y, t \geq 0; z = 0, h) \quad (13)$$

$$u_y(y, z, t) \rightarrow 0, \quad (y \rightarrow \infty, t \geq 0; 0 < z < h) \quad (14)$$

where θ is the slip coefficient, h denotes the distance between the side walls, and Eq. (14) is the natural boundary condition of the flow problem.

Download English Version:

<https://daneshyari.com/en/article/483785>

Download Persian Version:

<https://daneshyari.com/article/483785>

[Daneshyari.com](https://daneshyari.com)