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A new modal-based approach for modelling the bump foil structure in the simultaneous solution of foil-air bearing rotor dynamic problems

M.F. Bin Hassan, P. Bonello*

School of Mechanical, Aerospace and Civil Engineering, University of Manchester, UK

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ABSTRACT

Recently-proposed techniques for the simultaneous solution of foil-air bearing (FAB) rotor dynamic problems have been limited to a simple bump foil model in which the individual bumps were modelled as independent spring-damper (ISD) subsystems. The present paper addresses this limitation by introducing a modal model of the bump foil structure into the simultaneous solution scheme. The dynamics of the corrugated bump foil structure are first studied using the finite element (FE) technique. This study is experimentally validated using a purpose-made corrugated foil structure, including bump interaction and foil inertia, can be represented by a modal model comprising a limited number of modes. This full foil structure modal model (FFSMM) is then adapted into the rotordynamic FAB problem solution scheme, instead of the ISD model. Preliminary results using the FFSMM under static and unbalance excitation conditions are proven to be reliable by comparison against the corresponding ISD foil model results and by cross-correlating different methods for computing the deflection of the full foil structure. The rotor-bearing model is also validated against experimental and theoretical results in the literature.

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1. Introduction

A FAB (Fig. 1(a)) supports the shaft by means of a cushion of air bounded by a flexible foil structure. The introduction of the foil structure resolves the problems associated with the very tight radial clearance required by a plain air bearing. With a FAB, the hydrodynamic air film pressure generated as the shaft turns pushes the foil boundary away, allowing the shaft to become completely airborne [1]. The dynamics of FAB turbomachinery are governed by the interaction between the rotor, air films and the foil structures, and exhibit nonlinear effects. The governing system of equations comprises the state (i.e. time-based first order differential) equations of the air films, foils and rotor and can be expressed in state-space form as [2,3]:

$$\mathbf{S}' = \mathbf{\chi}(\tau, \mathbf{S})$$

(1)

where ()' denotes differentiation with respect to the time variable τ , **s** is the state vector, containing the n_{state} state variables (respectively relating to the air film, foil and rotor domains) and χ is a vector of nonlinear functions of the state variables.

* Corresponding author. *E-mail address:* philip.bonello@manchester.ac.uk (P. Bonello).

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		и	radius of rotor unbalance
		u	$n_{\text{bumps}} \times 1$ vector of the circumferential dis-
O′	differentiation with respect to τ		placements at the bump apexes
В	foil width (along $z_{\rm f}$ -axis)	w	foil deflection in radial direction
С	undeformed radial clearance of FAB	Ŵ	=w/c
D	diagonal damping matrix (Eq. (10))	W	$n_{\rm bumps} \times 1$ vector containing the radial dis-
Е	Young's Modulus		placements at the bump apexes
$f_{\rm res}$	fundamental resonance of single isolated	x, y, z	Cartesian coordinate system (Fig. 1)
	bump in Hz	$x_{\rm J}, y_{\rm J}$	Cartesian displacements of journal centre J
$\mathbf{f_p}$ F_x , F_y	vector of air pressure forces on bumps		relative to (fixed) bearing centre
F_x , F_y	Cartesian FAB forces	$x_{\rm f}, y_{\rm f}, z_{\rm f}$	Cartesian coordinate system local to the foil
1,]	generic degrees of freedom		structure (Fig. 2)
ñ	air film gap divided by c	$\alpha_{ij}(\omega)$	receptance FRF relating generic degrees of
H _{xf}	foil modal matrix whose columns are $\phi_{\mathbf{x_f}}^{(r)}$,		freedom <i>i</i> and <i>j</i>
	$r = 1n_{\rm f}$	Δ	diagonal matrix of squares of foil natural cir-
H _{yf}	foil modal matrix whose columns are $\phi_{\mathbf{y}_{\mathbf{f}}}^{(r)}$,		cular frequencies
	$r = 1n_{\rm f}$	$\varepsilon_{x,y}$	$x_j/C, y_j/C$
$k_{\rm bump}$	stiffness of single independent bump		$\overline{2}$
K	FE stiffness matrix of foil structure	ε	$=\sqrt{\varepsilon_x^2+\varepsilon_y^2}$
l _b	length of one bump along $x_{\rm f}$ -axis (Fig. 2)	ξ	$=z_{\rm f}/R$
L	axial length of FAB	ζ_{f_r}	viscous damping ratio of foil mode no. <i>r</i>
m _r	half-mass of rotor	θ	angular local bearing coordinate (Fig. 1)
М	FE mass matrix of foil structure	$\phi_i^{(r)}$	mass-normalised modal displacement of foil
n _{bumps}	number of bumps	(r)	in mode no. <i>r</i> at degree of freedom <i>i</i>
$n_{ m f}$	number of foil structure modes considered	$\boldsymbol{\varphi}_{\mathbf{x_f}}^{(r)}$	$n_{\text{bumps}} \times 1$ mass-normalised eigenvector in
n _{state}	total number of state variables in Eq. (1) .		mode no. <i>r</i> containing the circumferential
N_z, N_θ	number of points of FD grid along ξ , θ	\cdot (r)	displacements at the apexes of the bumps
	directions	$\mathbf{\Phi}_{\mathbf{y_f}}^{(r)}$	$n_{\text{bumps}} \times 1$ mass-normalised eigenvector in
	absolute air pressure at (ξ, θ) for FAB		mode no. <i>r</i> containing the radial displace-
p _a , p̃	atmospheric pressure, p/p_a respectively		ments at the apexes of the bumps
$\mathbf{q}_{\mathbf{f}}$	$n_{\rm f} \times 1$ vector of modal coordinates	γ	attitude angle
r	counter for foil modes	Λ	bearing number [2,3].
R	undeformed radius of FAB	μ	air viscosity
$\mathbf{R}_{\mathbf{x_fy_f}}(0)$	static receptance matrix relating the elements	ρ	density of foil material
- (-)	in u with those in $-\mathbf{f}_{\mathbf{p}}$	η	damping loss factor of foil structure
$\mathbf{R}_{\mathbf{y_f}\mathbf{y_f}}(0)$	static receptance matrix relating the elements	$\psi(\xi, \theta, \tau)$	
	in $-\mathbf{w}$ with those in $-\mathbf{f}_{\mathbf{p}}$	τ	non-dimensional time $(=\Omega t/2)$
S	vector of state variables (Eq. (1))	Ω	rotational speed in rad/s
S	static load in y direction per bearing	ω	general circular frequency in rad/s
Sp	pitch of bumps (Fig. 2)	ω_{f_r}	natural circular frequency of foil mode no. <i>r</i> in
t	time in seconds		rad/s
t _b	thickness of foil	χ	nonlinear vector function of s in Eq. (1).

Prior to the work by the co-author [2,3], the simultaneous solution of the state equations in (1) was avoided due to the computational burden involved. As explained in [2,3], this was typically done by approximating the subset of state equations relating to the air film as algebraic equations, effectively decoupling them from the rest of the equations. A non-simultaneous ("time-lagging") solution scheme was applied wherein the *current* value of the local air film thickness, and its rate of change with time, were estimated based on the foil and rotor displacements at the *previous* time step (e.g. as used in [4]). The work in [2,3] has shown that an efficient simultaneous solution is possible in both the time and frequency domains, which preserves the state-space format of Eq. (1) (and thus the true simultaneous coupling of its components), alleviating the potential convergence issues suffered by decoupled integration schemes [5].

The basic strategy for the simultaneous time-domain solution introduced in [2] has been independently verified by other researchers in [6] and successfully used to correlate with experimental results in [7], [5] and [3]. However, the simultaneous solution schemes in [2,3,5–7] are based on a simplified spring-damper model of the foil structure (as described below). The present paper addresses this simplification by devising a means of including more advanced features of the foil (namely, bump interaction and inertia) into the simultaneous solution scheme. The foil structure considered will be of the "first generation" type shown in Fig. 1(a), consisting of a single corrugated strip (comprising a series of bumps) underneath a top foil covering.

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