



# Convective heat transfer at high subcritical pressures in tubes with and without flow obstacles



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## ABSTRACT

Heat transfer measurements to CO<sub>2</sub>-cooled tubes with and without flow obstacles at high subcritical pressures were obtained at the University of Ottawa's supercritical test facility. Tests were performed for vertical upward flow in a directly heated 8 mm ID tube for a pressure range from 6.6 to 7.2 MPa, a mass flux range from 115 to 1000 kg m<sup>-2</sup> s<sup>-1</sup>, and a wall heat flux range from 1 to 122 kW m<sup>-2</sup>. Cases with spacing between obstacles equal to 250, 500 and 1000 mm were investigated. The results are presented in plots of wall temperature and heat transfer coefficient vs. bulk fluid enthalpy. The effects of flow conditions and obstacle geometry and spacing on high subcritical heat transfer have been discussed. Obstacles were found to suppress nucleate boiling, increase the CHF and lower the maximum post-CHF temperature. The maximum increases in CHF and in post-dryout heat transfer coefficient were observed for the obstacles having a spacing of 250 mm. Comparisons of the experimental data with prediction methods for single-phase heat transfer, CHF and post-dryout heat transfer are also presented.

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## 1. Introduction

During startup or postulated accidents, fuel elements of supercritical water-cooled reactors (SCWR) can experience cooling at high sub-critical pressures instead of supercritical ones. Therefore, accurate prediction of heat transfer and critical heat flux (CHF) characteristics at high subcritical pressures is essential for SCWR design. Knowledge of the effects of flow obstructions, such as fuel rod spacers, under such conditions is also important.

Several early studies have investigated experimentally heat transfer to flow of water in upward vertical, uniformly heated round tubes at high subcritical pressures; these include the studies by Schmidt (1960), Swenson et al. (1962), Miropol'skiy (1962), Bishop et al. (1965) and Herkenrath et al. (1967). Renewed interest on this topic was received in recent years. Some of the main results of the relevant studies will be summarized in the following paragraphs.

Schmidt (1960) performed experiments at high subcritical and supercritical pressures in the range from 17 to 30 MPa for water flows in vertical and horizontal tubes with 8 mm ID. He noted a significant decrease in CHF as the pressure approached the critical value. As a result of this, the CHF location and the maximum

temperature location shifted towards lower enthalpies, as pressure was increased under constant wall heat flux. Further increases of the system pressure above the critical pressure at the same heat flux and mass flux led to a gradual disappearance of a maximum in the wall temperature profile at a pressure of 30 MPa ( $P/P_c = 1.36$ ).

Swenson et al. (1962) investigated heat transfer to upwards vertical flow of water in a stainless steel tube with 10.4 mm ID at a pressure of 20.7 MPa, mass fluxes of 994 and 1356 kg m<sup>-2</sup> s<sup>-1</sup>, and inlet specific enthalpies in the range 1760–2220 kJ kg<sup>-1</sup>. They also investigated the case of an internally ribbed carbon-steel tube having an ID of 10.4 mm. Their results showed that, by comparison to the bare tube, the ribbed tube suppressed CHF and maintained nucleate boiling at higher bulk enthalpies. The film boiling data obtained from the bare tube were described within 10% by the modified form  $Nu = 0.076Re^{0.8}Pr^{0.4}$  ( $Nu$ ,  $Re$  and  $Pr$  are, respectively, the Nusselt, Reynolds and Prandtl numbers) of the Dittus and Boelter (1930) correlation.

Bishop et al. (1965) performed heat transfer experiments in water flowing in tubes made of hastelloy C and stainless steel 316 with 5.08 and 2.54 mm ID. Their pressure  $P$  was in the range 16.69–21.51 MPa ( $P/P_c = 0.76$ – $0.98$ ), the mass flux  $G$  was in the range 678–3390 kg m<sup>-2</sup> s<sup>-1</sup>, and the exit quality varied from 7 to 100%. With the exception of results at their lowest mass flux of 678 kg m<sup>-2</sup> s<sup>-1</sup>, their film boiling data were described within 19% by the expression  $Nu = 0.0193Re^{0.8}Pr^{1.23}(\rho_g/\rho_b)^{0.68}(\rho_g/\rho_l)^{0.068}$ ,

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## Nomenclature

CHF	critical heat flux, $\text{kW m}^{-2}$	$T_s$	saturated temperature, $^{\circ}\text{C}$
$\text{CHF}_{bare}$	critical heat flux in the bare tube geometry, $\text{kW m}^{-2}$	$T_w$	corrected wall temperature, $^{\circ}\text{C}$
$\text{CHF}_{corr}$	predicted CHF using CHF correlation of Appendix B, $\text{kW m}^{-2}$	$T_{w,o}$	tube outer wall temperature (measured by thermocouple), $^{\circ}\text{C}$
$\text{CHF}_{nfc}$	normalized predicted CHF using CHF correlation of Appendix B at nominal values of $P_o/P_c$ and $G_o$ , $\text{kW m}^{-2}$	$U_g$	streamwise gas velocity, $\text{m s}^{-1}$
$\text{CHF}_{obs}$	critical heat flux in the obstructed tube geometry, $\text{kW m}^{-2}$	$X$	vapour quality
$D$	inside diameter of the tube, mm	$X_{th}$	thermodynamic quality
$f$	friction factor	$z$	axial distance from the unheated test section inlet, mm
$G$	mass flux, $\text{kg m}^{-2} \text{s}^{-1}$	$z_{do}$	boiling crisis (CHF conditions) location, mm
$H$	heat transfer coefficient, $\text{kW m}^{-2} \text{K}^{-1}$	$z_h$	axial distance from the heated test section inlet, mm
$h_b$	specific bulk enthalpy, $\text{kJ kg}^{-1}$		
$h_{fg}$	latent heat, $\text{kJ kg}^{-1}$	<i>Greek symbols</i>	
$h_{in}$	specific enthalpy at inlet, $\text{kJ kg}^{-1}$	$\Delta Q_e$	power, W
$h_{out}$	specific outlet enthalpy, $\text{kJ kg}^{-1}$	$\Delta z_{CHF}$	axial distance between the CHF location and the downstream edge of the nearest upstream obstacle, mm
$k$	thermal conductivity, $\text{kW m}^{-1} \text{K}^{-1}$	$\Delta z_d$	axial distance from the downstream end of an obstacle, mm
$L_h$	heated length, mm	$\varepsilon$	blockage ratio, %
$\dot{m}$	mass flow rate, $\text{kg s}^{-1}$	$\mu_b$	bulk dynamic viscosity, Pa s
Nu	Nusselt number	$\mu_g$	gas dynamic viscosity, Pa s
$\text{Nu}_b$	bulk Nusselt number	$\rho_b$	fluid density at bulk temperature, $\text{kg m}^{-3}$
$\text{Nu}_{bare}$	Nusselt number in a bare tube	$\rho_g$	gas density at saturation temperature, $\text{kg m}^{-3}$
$\text{Nu}_{obs}$	Nusselt number in an obstructed tube	$\rho_l$	liquid density at saturation temperature, $\text{kg m}^{-3}$
$P$	system pressure, MPa	$\sigma$	surface tension, $\text{N m}^{-1}$
$P_c$	critical pressure, MPa	<i>Acronyms</i>	
$\text{Pr}_b$	bulk Prandtl number	DPT	differential pressure transducer
$Q_e$	nominal electrical power supplied to the test section, kW	HTC	heat transfer coefficient
$q$	heat flux, $\text{kW m}^{-2}$	ID	tube inner diameter
$q_{loss}$	heat flux to surroundings (heat losses), $\text{kW m}^{-2}$	LHS	left hand side
$q^v$	volumetric heat generation, $\text{kW m}^{-3}$	OD	tube outer diameter
$\text{Re}_b$	bulk Reynolds number	PS	power supply
$r_i$	test section inner radius, mm	PT	pressure transducer
$r_o$	test section outer radius, mm	RHS	right hand side
$T_b$	bulk temperature, $^{\circ}\text{C}$	SCUOL	University of Ottawa supercritical $\text{CO}_2$ loop
$T_{in}$	test section inlet temperature, $^{\circ}\text{C}$	SCWR	super-critical water reactor
$T_{Pc}$	pseudo critical temperature at supercritical pressures, $^{\circ}\text{C}$		

where the bulk density was calculated from the gas and liquid densities as  $\rho_b = \rho_g(\alpha) + \rho_l(1 - \alpha)$  and the void fraction  $\alpha$  was calculated by the homogenous model.

Chen et al. (2015) conducted CHF tests for subcooled water at high subcritical pressures in tubes made of Inconel 625 having ID of 4.62, 7.98 and 10.89 mm. The operating conditions were:  $P = 1\text{--}20$  MPa ( $P/P_c = 0.05\text{--}0.90$ ),  $G = 454\text{--}4055$   $\text{kg m}^{-2} \text{s}^{-1}$ , and inlet subcooling  $T_s - T_{in} = 53\text{--}361$  K. Their data showed that CHF was inversely proportional to  $D^{0.35}$ , where  $D$  is the internal diameter of the tube. The authors correlated their CHF data as  $\text{CHF} = Cq_s$ , where  $q_s = (h_s - h_{in})GD/(4L_h)$  was the heat flux that resulted in saturation temperature at the end of the heated length, and  $C = \text{Min} [2350(1 - 0.0307P)(G(H_s - H_{in}))^{-0.35}(D/0.008)^{-0.35}, 1.0]$ . This correlation predicted the high subcritical data with an error ranging from  $-15$  to  $20\%$ .

The experiments in water require expensive high pressure test facilities operating at high temperatures and powers. A more convenient approach for investigating high-pressure subcritical heat transfer is by means of using modelling fluids. Recent studies using refrigerant R-134a as a modelling fluid include the ones by Hong et al. (2004), Vijayarangan et al. (2005), and Chun et al. (2007).

Hong et al. (2004) investigated CHF occurrence for R-134a upward flows in an internally heated vertical annular channel made of Inconel 601. They performed CHF tests at both steady

and transient pressures through the critical pressure. The operating conditions were as follows:  $P = 0.63\text{--}3.98$  MPa ( $P/P_c = 0.16\text{--}0.98$ ),  $G = 501, 1002$  and  $1502$   $\text{kg m}^{-2} \text{s}^{-1}$ , and inlet subcooling  $T_s - T_{in} = 14.5, 24.8, 34.5$ , and  $44.3$  K. For steady pressures, the CHF value decreased as the pressure approached the critical pressure. Mass flux and inlet subcooling had small effects on the CHF at near-critical pressures. The sharp decrease in the CHF near the critical pressure was attributed to the dramatic decrease in the latent heat  $h_{fg}$ . During transient tests with the pressure being reduced from supercritical to high subcritical values, the CHF occurred as soon as the pressure dropped below the critical value.

Vijayarangan et al. (2005) conducted CHF experiments in a vertical stainless steel tube having an ID of 12.7 mm and cooled by R-134a. The operating conditions were as follows:  $P = 1\text{--}4$  MPa ( $P/P_c = 0.25\text{--}0.98$ ),  $G = 200\text{--}2000$   $\text{kg m}^{-2} \text{s}^{-1}$ , and exit quality in the range  $0.17\text{--}0.94$ . They compared their CHF data with predictions based on the look-up table of Groeneveld et al. (1996) and the correlation of Katto and Ohno (1984). The prediction methods were in good agreement with experimental values for low pressures ( $P/P_c < 0.5$ ). At high reduced pressures ( $P/P_c = 0.6\text{--}0.94$ ), however, the prediction error varied between  $20$  and  $250\%$  for the correlation of Katto and Ohno (1984) and between  $-50\%$  and  $350\%$  for the look-up table. Vijayarangan et al. (2005) modified the correlation of Katto and Ohno (1984) to the following form

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