



Performance evaluation of CO-OFDM systems based on electrical constant-envelope signals



Vinicius O.C. Dias^a, Ezequiel da V. Pereira^a, Helder R.O. Rocha^{b,*}, Marcelo E.V. Segatto^a, Jair A.L. Silva^a

^aLaboratório de Telecomunicações, Universidade Federal do Espírito Santo, Vitória-ES, Brazil

^bCentro Universitário Norte do Espírito Santo, São Mateus-ES, Brazil

ARTICLE INFO

Article history:

Received 16 January 2017

Revised 21 April 2017

Accepted 22 June 2017

Keywords:

Coherent detection

Constant-envelope

Orthogonal frequency division multiplexing

Phase modulation index

ABSTRACT

The influence of the electrical phase modulation index h in the performance of constant-envelope orthogonal frequency division multiplexing (CE-OFDM) in coherent detection optical systems is treated analytically and its range of validity examined by simulations. A compromise between h and subcarrier mapping is identified according to differences in sensitivity related to non-linearities inserted by the optical modulator. It is shown that the proposed scheme outperforms conventional coherent detection OFDM systems, which is strongly dependent on both phase and optical modulation indexes.

© 2017 Elsevier Inc. All rights reserved.

1. Introduction

The resilience towards fiber dispersion is the main attractive feature of orthogonal frequency division multiplexing (OFDM) signal processing in optical communications systems [1–3]. However, large peak-to-average power ratio (PAPR) is one of the main drawbacks in both direct and coherent detection optical orthogonal frequency division multiplexing (CO-OFDM) systems, as it not only limits the resolution of digital-to-analog converters and RF amplifiers, but also reduces the tolerance to the non-linearities introduced by optical modulators and optical fibers [1,4].

To reduce PAPR, an effective technique based on the transmission of constant-envelope (CE-OFDM) signals is a suitable solution to the aforementioned problems [5]. Proposed in [6] and experimentally demonstrated in [7,8] to combat nonlinear degradations in direct-detection optical OFDM systems, this electrical phase modulation based technique is, for the first time, introduced in coherent detection systems.

Unlike the approaches evaluated in [9,10], the intermediate electrical CE-OFDM signals of this solution are used to modulate the continuous wave (CW) laser source, employing a conventional one-branch Mach–Zehnder modulator (MZM).

Section 2 of this letter develops an analytical expression from which the influence of the electrical phase modulation index h

may be derived. Its validity is examined by means of numerical simulations in Section 3. A tradeoff between the index h and the subcarrier mapping is clarified from the simulation results described in this Section. It is shown that, in CO-OFDM systems that employ such multicarrier constant-envelope signals, high subcarrier modulation levels requires reduction in the aforementioned index due to non-linearities inserted by the MZM. Section 4 concludes the paper.

2. Theoretical model

Fig. 1 shows the CE-OFDM considered in this analysis. It is a modulation format where an electrical carrier is phase modulated by OFDM waveforms, which results in envelope signals with PAPR = 3 dB [6]. Therefore, a single OFDM signal with mean power given by σ_s^2 is written as

$$x(t) = C \sum_{k=0}^{N_s-1} \Re[X(k)] \cos(2\pi\Delta_f kt) - \Im[X(k)] \sin(2\pi\Delta_f kt) \quad (1)$$

with $\{X(k)\}_{k=1}^{N_s-1}$ the M -QAM data symbols, $T = \frac{N}{F_s}$ the OFDM symbol time, $N = 2N_s + 2$ the fast-Fourier transform length, F_s the sampling rate and C a constant, modulates the phase of a carrier resulting in a bandpass signal given by

$$c(t) = A \cos[2\pi f_0 t + \phi(t)] = A \cos[2\pi f_0 t + \theta_n + 2\pi h C_N x(t)], \quad (2)$$

where A is the signal amplitude, f_0 is the carrier frequency, $\phi(t)$ is the phase signal during the n -th signal interval $nT \leq t < (n+1)T$,

* Corresponding author.

E-mail addresses: helder.rocha@ufes.br (H.R.O. Rocha), jair.silva@ufes.br (J.A.L. Silva).

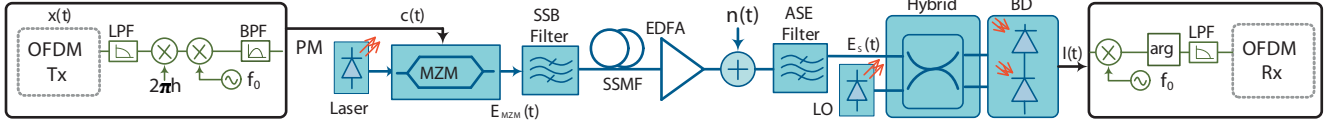


Fig. 1. CO-CE-OFDM system model. LPF: Low Pass Filter; BPF: Band Pass Filter; PM: Phase modulation; MZM: Mach-Zehnder Modulator; SSB: Single-Sideband; EDFA: Erbium-doped Fiber Amplifiers; SSMF: Standard Single-Mode Fiber; ASE: Amplified Spontaneous Emission; LO: Local Oscillator; BD: Balanced Photoreceiver.

θ_n is a memory term designed to make the modulation phase continuous (when $\theta_n = 0$ the modulation is memoryless), h is referred to as the electrical phase modulation index, and C_N a constant used to normalize the variance of the message signal $x(t)$ [6]. Continuous phase modulation (CPM) can be achieved by introducing memory in Eq. (2) making $\theta_n \neq 0$ [11]. Although the constant-envelope property and the compact signal spectrum of this class of CPM signaling, the memory increases the receiver complexity [5,12,13].

The bandwidth of the electrical signal $c(t)$ is usefully expressed as $B = \max(2\pi h, 1)B_W$, which is an RMS (root-mean-square) bandwidth lower bounded by the conventional OFDM signal bandwidth B_W [5]. A conservative bandwidth is estimated with this expression for small values of $2\pi h$. For $2\pi h > 1$ it measures at most 92% of the total bandwidth [14].

Assuming $\theta_n = 0$ and $C_N = 1$, for the sake of simplicity, the CE-OFDM signal becomes $c(t) = A \cos[w_0 t + s(t)]$, for $w_0 = 2\pi f_0$ and $s(t) = 2\pi h x(t)$. Therefore, the optical field at the output of a conventional single arm Mach-Zehnder modulator (MZM) characterized by its switching voltage V_π and biased by V_{bias} is a bandpass signal centered at frequency $w_c = 2\pi f_c$ given by

$$E_{MZM}(t) = \cos \left[\frac{\pi c(t)}{2V_\pi} - \frac{\pi V_{bias}}{2V_\pi} \right] \cdot \sqrt{2P} \cos(w_c t), \quad (3)$$

where P is the power of the CW laser signal fed into the MZM optical input. Biased at its null point ($V_{bias} = -V_\pi$) the MZM output signal can be rewritten as

$$\begin{aligned} E_{MZM}(t) &= \sin \left[\frac{\pi c(t)}{2V_\pi} \right] \cdot \sqrt{2P} \cos(w_c t) \\ &= \sqrt{2P_c} \cdot \sin[F(t)] \cdot \cos(w_c t), \end{aligned} \quad (4)$$

for $F(t) = \frac{\pi c(t)}{2V_\pi}$. After optical pre-amplification with gain G , the input signal into the balanced receiver can be written as

$$\begin{aligned} E_s(t) &= \sqrt{2GP} \cdot \sin[F(t)] \cdot \cos(w_c t) + n_i(t) \cos(w_c t) + n_q(t) \\ &\quad \times \sin(w_c t), \end{aligned} \quad (5)$$

for $n_i(t)$ and $n_q(t)$ the in-phase and quadrature component of the amplifier noise with power spectrum density given by $N_{ASE}/2$ and variance $\sigma_n^2 = N_{ASE} \cdot B_o$, with B_o the bandwidth of the optical filter at its output.

If we describe the optical field of the local oscillator (LO) light source as $E_{LO} = A_{LO} e^{jw_{LO}t + j\theta_{LO}}$, then $\Re\{E_{LO}\} = A_{LO} \cos(w_{LO}t)$ for E_{LO} its amplitude, w_{LO} its frequency and θ_{LO} the laser phase noise, which is zeroed for the sake of simplicity. Because of the intrinsic heterodyne detection of the CO-CE-OFDM system model, we assume that electrical synchronous demodulation is used to estimate the phase noise in order to benefit the phase demodulation [15].

The beating of the optical signal and the LO at the 180° optical hybrid results in two photocurrents outputs from the balanced photodiodes (PD) defined as

$$\begin{aligned} I_1(t) &= R \left| \Re \left\{ \frac{1}{\sqrt{2}} (E_s + E_{LO}) \right\} \right|^2 \\ I_2(t) &= R \left| \Re \left\{ \frac{1}{\sqrt{2}} (E_s - E_{LO}) \right\} \right|^2, \end{aligned} \quad (6)$$

with R denoting the PD responsivity and \Re the real number set. Hence, the $I_2(t)$ current can be expressed as

$$\begin{aligned} I_2(t) &= R \left\{ 2GP \sin^2(F(t)) \cos^2(w_c t) + \right. \\ &\quad + n_i^2 \cos^2(w_c t) + n_q^2 \sin^2(w_c t) + \\ &\quad + A_{LO}^2 \cos^2(w_{LO} t) + \\ &\quad + 2[\sqrt{2GP} \sin(F(t)) \cos^2(w_c t) n_i + \\ &\quad + \sqrt{2GP} \sin(F(t)) \cos(w_c t) \sin(w_c t) + \\ &\quad + n_i n_q \sin(w_c t) \cos(w_c t) - \\ &\quad \left. - A_{LO} \cos(w_{LO} t) (\sqrt{2GP} \sin(F(t)) \cos(w_c t) + \right. \\ &\quad \left. + n_i \cos(w_c t) + n_q \sin(w_c t))] \right\}. \end{aligned} \quad (7)$$

Using the same mathematical description in $I_1(t)$ current, the resulting photon electrical current can be expressed as follows

$$\begin{aligned} I(t) &= I_1 - I_2 \\ &= 4RA_{LO} \cos(w_{LO} t) [\sqrt{2GP} \sin(F(t)) \cos(w_c t) + \\ &\quad + n_i \cos(w_c t) + n_q \sin(w_c t)]. \end{aligned} \quad (8)$$

Assuming that the LO is frequency locked to the received signal ($w_{LO} = w_c$), the current at the balanced receiver output can be written as

$$\begin{aligned} I(t) &= 4RA_{LO} [\sqrt{2GP} \sin(F(t)) \cos^2(w_c t) + n_i \cos^2(w_c t) \\ &\quad + n_q \sin(w_c t) \cos(w_c t)] \\ &= 2RA_{LO} [\sqrt{2GP} \sin(F(t)) (1 + \cos(2w_c t)) + n_i (1 + \cos(2w_c t)) \\ &\quad + n_q \sin(2w_c t)]. \end{aligned} \quad (9)$$

After low-pass filtering to filter out the high-order items, this signal is approximated to

$$\begin{aligned} I(t) &= 2RA_{LO} [\sqrt{2GP} \sin(F(t)) + n_i] \\ &\approx 2RA_{LO} \left[\sqrt{2GP} \left(\frac{\pi c(t)}{2V_\pi} \right) + n_i \right], \end{aligned} \quad (10)$$

considering the first-order Taylor's expansion that allows $\sin[F(t)] = \sin \left(\frac{\pi c(t)}{2V_\pi} \right) \approx \frac{\pi c(t)}{2V_\pi}$. Substituting $F(t)$ and Eq. (2) into (10), the received bandpass signal becomes

$$\begin{aligned} I(t) &= \frac{RA_{LO} \sqrt{2GP} \pi}{V_\pi} \cdot A \cos[w_0 t + s(t)] + 2RA_{LO} \cdot n_i \\ &= K_1 \cos[w_0 t + s(t)] + K_2 n_i, \end{aligned} \quad (11)$$

for $K_1 = \frac{RA_{LO} \sqrt{2GP} \pi A}{V_\pi}$ and $K_2 = 2RA_{LO}$. After multiplying this passband signal by the electrical carrier $\cos(w_0 t + \frac{\pi}{2})$, for downconversion, the expression

$$\begin{aligned} I(t) &\approx \frac{K_1}{2} \left[\cos \left(2w_0 t + \frac{\pi}{2} + s(t) \right) + \cos \left(\frac{\pi}{2} - s(t) \right) \right] + K_2 n_i \\ &\quad \times \cos \left(w_0 t + \frac{\pi}{2} \right) \end{aligned} \quad (12)$$

is obtained. Neglecting the high frequency contributions and considering $\sin(s(t)) \approx s(t) = 2\pi h x(t)$ after first-order Taylor expansion this photocurrent becomes

Download English Version:

<https://daneshyari.com/en/article/4957089>

Download Persian Version:

<https://daneshyari.com/article/4957089>

[Daneshyari.com](https://daneshyari.com)