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Effects of sudden change in pitch angle on oscillating wind turbine airfoil performances



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ABSTRACT

In this paper, the aerodynamic performances of a wind turbine airfoil in sinusoidal pitching motion was numerically predicted and analyzed using the two-dimensional singularity method. A parametric study based on the sudden pitch angle rate variation, the reduced frequency and the pitching amplitudes have been investigated.

The results highlight the sudden change effect in pitch angle on the aerodynamic performances of the airfoil and the shedding wake. The aspect of lift and drag coefficient loops and hysteresis were customized regarding to the pure sinusoidal motion case. The dynamic stall phenomenon, which occurs during the rapid change in pitch angle, produces more drag forces influencing the wind turbines performance.

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1. Introduction

Pitching airfoils have several industrial applications such as wind energy, power extraction from flapping airfoil, micro-aerial vehicle, and many other domains. Various parameters have relevant influence on the pitching airfoil performance. They can be either natural such as the wind speed and direction, the wind turbulence, the structure instabilities in the airflow, or mechanical such as the pitching amplitude, the reduced frequency of the oscillation and the oscillating motion shape. This important number of parameters encourages researchers to investigate their impact in order to optimize the system design and inhance its performance.

The oscillating motion can be a sinusoidal or a non-sinusoidal shape, which has a relevant effect on the wake and the airfoil aerodynamic performance. Literary, many numerical and experimental investigations have been conducted to study the sinusoidal oscillation motion dynamics, either in a pure pitching motion, in a pure plunging motion or in a flapping (combined pitching/plunging) motion. In fact, in flapping airfoil investigation case, Anderson et al. [1], Lai and Platzer [17], Read et al. [27], Tuncer and Platzer [29], Jones et al. [12], Young and Lai [34], Ashraf et al. [4], La Mantia and Dabnichki [19–24] have investigated the thrust and the propulsive forces generated by the airfoil and their dependence on the reduced frequency, the pitching amplitude, the camber and thickness of the airfoil.

Recently, many numerical investigations have been made to study the sinusoidal pure pitching airfoil characteristics. Amiralei et al. [3] have used the OpenFOAM software based on the finite volume discretization scheme to solve the two-dimensional Navier-Stokes equations. They have studied the influence of unsteady parameters on pitching NACA0012 airfoil at low Reynolds number between 555 and 5000. Results shown that, the maximum lift coefficient, the hysteresis loops and the generated wake were affected by those parameters. Wang et al. [30] and Gharali et al. [7] have numerically investigated the dynamic stall of a pitching airfoil using the package AnsysFluent software solving Navier–Stokes equations with k ω -SST turbulence model. They have found that the ko-SST model has the capability to predict the aerodynamic loads when the dynamic stall phenomenon occurs. Moreover, this model can well capture the leading edge vortex structures generated by the airfoil during his movement. They have noticed also, that the lift coefficient under dynamic stall is more important than under the static stall. The high circulation of the leading edge vortex is the main responsible parameter of this increase.

However, the research works are limited and a few studies have been made in terms of non-sinusoidal pitching motion case. Koochesfahani [14] has experimentally studied the vortical patterns in the wake of sinusoidal and non-sinusoidal pitching airfoils. He observed the existence of an axial airflow in the cores of the wake vortices with magnitude presenting a linear dependency on the oscillation frequency and amplitude. Recently, the non-sinusoidal pitching airfoil has been used to study its effects on the thrust and the propulsion generated by a flapping foil. Hover et al. [11] have experimentally investigated the flapping airfoil perfor-

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Nomencl	\vec{V}	to	
	.1	$[(V^t)_i]_k$	ta
Α	the airtoil surface area	$(v_w)_k$	to
а	pitch axis location with respect to the $\frac{1}{2}$ chord		p
c	airfoil chord length	(x _{mi} , y _{mi}) co
C _T	thrust coefficient	x_m	х
Ср	pressure coefficient	y_m	У
Cl	lift coefficient	Create an	h 1
Cd	drag coefficient	Greek syl	ποοι
Cm	moment coefficient	α	att
C _x	force coefficient at x-direction	β_k	the
Cy	force coefficient at y-direction	φ	tot
C(k)	Theodorsen's transfer function	$\phi\infty$	fre
d_k	the shed panel length	ϕS	SO
E_{ij}^n	normal velocity component induced at the <i>i</i> th panel	ϕV	vo
-	control point by unit strength source distribution on	ϕW	wa
	the <i>j</i> th panel	$(\phi_{le})_k$	dis
E_{ii}^t	tangential velocity component induced at the <i>i</i> th panel	$(\phi_i)_k$	dis
- 5	control point by unit strength source distribution on	Γ_k	the
	the <i>j</i> th	γ	vo
\overline{F}	mean force applied on the airfoil in x direction	$(\gamma_{\rm sh})_{\rm k}$	vo
F_{ii}^n	normal velocity component induced at the <i>i</i> th panel	γ _w	the
.,	control point by unit strength vorticity distribution on	Ω(t)	an
	the <i>j</i> th panel	ω	OS
$(F_{i,N+1}^n)_k$	normal velocity component induced at the <i>i</i> th panel	ρ	is
1,N+1 K	control point by unit strength vorticity distribution on	σ	SO
	the shed vorticity panel at time t_k .	$\theta(t)$	pit
$F_{\cdot \cdot}^t$	tangential velocity component induced at the <i>i</i> th panel	θ_0	pit
1)	control point by unit strength vorticity distribution on	$\dot{\theta}$	ve
	the <i>i</i> th panel	$\ddot{\theta}$	ac
$(F^t,\ldots)_k$	tangential velocity component induced at the <i>i</i> th panel		
$i, N+1^{\prime \kappa}$	control point by unit strength vorticity distribution on	Subscript	
	the shed vorticity panel at time $t_{\rm h}$	ε	up
$(G^n_{i})_i$	normal velocity component induced at the <i>i</i> th panel		
(im)K	control point by unit strength <i>m</i> th core vortex at time	mance in d	iffor
	t.	obtain three	e dif
$(G^t)_i$	tangential velocity component induced at the <i>i</i> th panel	ric saw too	th w
(im ^k	control point by unit strength <i>m</i> th core vortex at time	reculte with	the
	t.	oct thrust a	ooffi
1.45	r_{κ}	function or	uol i
h(t)	plunging amplitude	Stroubal pu	uai
h0 İn	velocity due to the plunging motion	studied num	nomi
n h	acceleration due to the plunging motion	but the second	hima
n i	unit vectors in the airfoil fixed coordinate system (x, y)	by the com	Dina
;	unit vectors in the airfoil fixed coordinate system (x,y)	on the gene	rate
J k	reduced frequency	(INS) solver	wit
л 1	perimeter of the airfoil	Der. Ine stu	iay '
L N	total airfoil papale number	maximum a	angl
IN D	total alloli pallels liullidel	and pitchin	.g m
P	pressure at each mu-point panel	ping airtoil	wer

- r
 distance between the *j*th panel and the respective midpoint of the *i*th panel

 R_t
 the pitching rate
- Strouhal numberSparametric variable
- *S* length of the airfoil panels
- $S_{\rm sh}$ length of the shed panel
- $S_{\rm w}$ length of the wake panels
- time step
- u(t) translation velocity at x direction
- $(u_w)_k$ total velocity components in x-directions of the shed panel at the airfoil fixed coordinate system. v_{∞} free stream velocity
- V_{∞} free stream velocity v(t) translation velocity at y direction
- *V*_{stream} velocity in the fixed coordinate system

\overline{V}	total velocity of the airfoil	
$[(V^t)_i]_k$	tangential velocities at time step t_k of the <i>i</i> th panel	
$(v_w)_k$	total velocity components in y-directions of the shed	
	panel at the airfoil fixed coordinate system	
(x _{mi} , y _{mi})	coordinates of panel mid-point	
x_m	x coordinate of <i>m</i> th core vortex at time t_k	
y _m	y coordinate of <i>m</i> th core vortex at time t_k	
Greek symbols		
α	attack angle	
β_k	the shed panel inclination	
, к ф	total potential	
φm	free stream potential	
φω dS	source potential	
фV	vortex potential	
φW	wake potential	
$(\phi_{le})_k$	disturbance potential at the airfoil leading edge	
$(\phi_i)_k$	disturbance potential at the i th panel control point	
Γ_k	the airfoil total circulation	
γ	vortex distribution strength on each airfoil panel	
$(\gamma_{\rm sh})_{\rm k}$	vortex strength of shed panel	
γ _w	the wake vortex strength	
Ω(t)	angular velocity	
ω	oscillation frequency	
ρ	is the flow density	
σ	source distribution strength on each airfoil panel	
θ(t)	pitching function	
θ_0	pitching amplitude	
$\dot{ heta}$	velocity due to the pitching motion	
$\ddot{ heta}$	acceleration due to the pitching motion	
Subscript		
ε	upstream panels mid points	

ent cases of combined pitching and plunging motions to ferent shapes of attack angles: a square wave, a symmetave and a cosine function. They compared the obtained case of ordinary attack angle. They found that the highicient and efficiency were obtained in the case of cosine to $C_{\rm T} = 0.624$ and 64%, respectively, for $\alpha_{\rm max} = 20^{\circ}$ and er St = 0.35. In the same context, Xiao and Liao [32] have cally the effect of a cosine function attack angle obtained tion of pitching and plunging NACA0012 airfoil motion d thrust using unsteady two-dimensional Navier–Stokes h a compressible flow at a low free-stream Mach numwas based on the variation of the Strouhal number, the e of attack and the phase angle between the plunging otions. A good efficiency and performance of the flape observed when they change the pitching motion from a sinusoidal function to a non-sinusoidal function composed by a cosine function and arctangent function relating the plunging velocity and the free stream velocity. To enlarge their study, Xiao et al. [33] have investigated the effects of the non-sinusoidal pitching motion on the energy extraction of an oscillating airfoil used in a bio mimic energy generator. A trapezoidal pitching motion was combined with the plunging motion to compare the obtained efficiency with that of sinusoidal pitching motion case. For that, a parametric trapezoidal function was used. It had shown that a remarkable increase in the extracted power and the total output efficiency for S > 1. They have mentioned also the existence of an optimal pitching profile that can be determined to obtain an output power coefficient equal to 63% and output efficiency equal to 50%. Lu et al. [15] have numerically investigated the effects of the large amplitude, the non-sinusoidal motion and the airfoil camber on the energy extracted from pitching airfoil. They applied the finite volume method using the CFD package CFX in two-dimensional incompressible flow case.

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