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# Detailed contact pressure between wire rope and friction lining

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### ABSTRACT

The authors have studied the contact pressure between a wire rope and a friction lining block. The current study is crucial for the subsequent work on frictional heating under sliding conditions. The contact equations are solved on a fine grid, using multigrid techniques. Results from the numerical contact model show that the local maximum contact pressure is about 45 times higher than the average one obtained from the "plane method" formula. Two approximate analytical calculations are proposed that give a maximum contact pressure with a relative error of the order of 10% with only a small computational effort.

# 1. Introduction

Multi-rope friction hoists are widely used in vertical shafts of coal mines. Wire ropes are bent over a hoisting sheave which periphery is fixed with many friction lining blocks. Such a hoist utilizes friction between wire rope and friction lining to lift coal, minerals, equipment and personnel [1].

Under exceptional operating conditions, such as an emergency braking when lowering a heavy load, an accident of the wire rope sliding against the hoisting sheave can happen [2]. Once it occurs, the frictional heating causes the temperature on the groove's surface to rise, resulting in a Coefficient of Friction (COF) reduction of the contact pair. When the friction force becomes smaller than the heavy load, the sliding speed increases (possibility of a vicious circle with dramatic consequences).

Therefore, it is crucial to study the transient temperature field of a friction lining block under sliding conditions to determine the dominant factors influencing the temperature. In order to solve the transient thermal problem, it is necessary to first calculate the contact pressure on the surfaces at any time to determine the heat flux.

The pioneering work by Heller [3] presented a derivation of the contact pressure between wire rope and sheave's groove. Heller pointed out that the formula had been mentioned in the work of Drucker [4] but the source was not known. It has been widely used as the average contact pressure between wire rope and friction lining [5] as well as another formula presented in [6]. Both formulas are based on simplifying a wire rope into a cylinder. On the basis of this simplification, some researchers further considered the contact pressure distribution on a cross-section plane to be uniform to obtain a simple

analytical expression for the heat flux [7,8] or to be of cosine shape to obtain a conversion formula for the COF [9]. Recent work by Wang [10] similarly homogenized the rope into a cylinder and then established a 3d finite element contact model with dynamic tension being applied to the two ends of the rope. The work by Peng [1,2] simplified a 6-strand wire rope into a formation of six helically twisted cylinders and considered the contact pressure distribution (on the transverse rope section) between any of the three or two small male circles and the female circle to be of cosine shape.

The above-mentioned contact models are all far from precise, because of major simplifications of the wire rope shape. The objective of the current work is to build a more precise 2-dimensional contact model so as to obtain a more precise pressure distribution on the contact surfaces.

According to the Hertzian theory, contact pressure depends on the contact geometries and the distance between the two contact surfaces. However, the geometry of a stranded wire rope is very complicated. It contains single-helical, double-helical and, depending on the rope construction, even multi-helical wires.

There are many models of a stranded wire rope geometry, which involve deriving Coordinate Equations (CE) of helical wires' centroidal axes. Stein and Bert [11] were amongst the first researchers who considered double-helical wires in a rope. They presented the CE and the curvature equation for the ordinary lay double helix with a very brief derivation. Karamchetty [12] also derived CE of double helical wires, which, however, does not agree with those of Stein and Bert because they do not distinguish between Lang's lay and ordinary lay.

Lee [13,14] gave a comprehensive study into wire rope geometry and derived CE (using a local coordinate system referred to as the

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Nomenclature		Р	dimensionless pressure
		$R_{1x}, R_{2x}$	radius of curvature in $x$ direction of wire and groove
$a_e$	semi-major axis of contact ellipse	$R_{1y}, R_{2y}$	radius of curvature in $y$ direction of wire and groove
$b_e$	semi-minor axis of contact ellipse	$R_g$	groove radius
b	half-width of Hertzian line contact	$R_s$	strand helical radius
$d_{wo}$	diameter of outermost wires	Rw	wire helical radius
D	nominal diameter of wire rope	$R_{x}$	reduced radius of curvature in $x$ direction
E'	reduced modulus of elasticity	$R_y$	reduced radius of curvature in $y$ direction
i	index of an outermost wire in a strand	R <sup>;</sup>	reduced radius of curvature in elliptical contact
j	index of a strand in a rope	w	2d load
$k_e$	ratio of contact semi-axes	x	coordinate in direction of wire rope length
l <sub>wij</sub>	coordinate equations of the centroidal axis of an outer-	X	dimensionless coordinate in direction of wire rope length
	most wire	y	coordinate perpendicular to direction of wire rope length
l <sub>wij_b</sub>	coordinate equations of the outermost bus bar of an	Y	dimensionless coordinate perpendicular to direction of
	outermost wire		wire rope length
$L_{g}$	length of groove	$\beta s$	lay angle of strand in a rope
$n_x$	number of grid points in x direction	βw	lay angle of outer wire in a strand
$n_{\rm y}$	number of grid points in $y$ direction	δ	maximum deformation of Hertzian line contact
p	pressure	$\Delta d$	grid spacing
$p_h$	maximum Hertzian pressure of line contact		

Serret [15] and Frenet [16] frame) of double helical wires for stranded wire ropes in regular lay and in Lang's lay. Wang [17] developed a generalized mathematical method describing the wire geometry in any round-strand wire rope and gave a detailed example for deriving the vector equations of the double helices and single helices. On the basis of this work, Wang [18] gave CE for double-helical wires considering the wire rope's self-twist coefficient which is a technological parameter measuring the self rotation of the centerline of a strand in the process of forming a wire rope; Usabiaga et al. [19] derived the parametric equations of double-helical wires in a deformed case.

When a straight stranded wire rope is wound around a drum or bent over a sheave, the double helical wires take on the form of a triple helix. Their CE can be obtained by analogy. Those in a matrix form for a specific stranded wire rope were directly presented in [14]. Based on this work, Ma [20] derived the CE of the triple-helical helices in a 7\*7-Wire Strand Core wire rope which is bent over a sheave. After the CE of the double helices and single helices in a stranded wire rope are obtained, the modeling of the surfaces of the wire rope can be accomplished using CAD software [21,22] but seldom in analytical expressions.

Based on the above-mentioned studies on deriving CE of doublehelical wires, the paper first obtains the CE of the outmost wires in a 6\*7+FC wire rope. Because of the recursive property of the derivation presented by Lee [14], the tangential vector and normal vector of a double helix at any point can be obtained without singularity in the first or second derivative, on the basis of which, the present paper manages to analytically describe the surfaces of the outmost wires in a discrete form. To reduce the computational cost from the precise model of the wire rope geometry, an efficient multigrid solver is proposed to obtain the contact pressure. Multigrid methods are iterative acceleration methods based on the finite difference method. They solve a problem on a sequence of grids of different mesh size. They are very suitable for large scale problems as their efficiency permits one to use millions of points in the contact area [23]. The current study provides the basis for subsequent work on frictional heating of the contact under sliding conditions.

#### 2. Theory

## 2.1. CE of the centroidal axes of the outermost wires

The geometry of a  $6 \times 7 + FC$  wire rope in left-hand regular lay is shown in Fig. 1. The CE of the centroidal axes of the outermost wires are expressed in vector form as:

$$\mathbf{l}_{wij} = [R_s \theta_s \tan \beta_s - R_w \cos \beta_s \sin(k_{ws} \theta_s + \alpha_{wi});$$
  

$$\cos(-\theta_s + \alpha_{sj})(R_s - R_w \sin(k_{ws} \theta_s + \alpha_{wi})) + R_w \sin \beta_s \sin(-\theta_s + \alpha_{sj})$$
  

$$\sin(k_{ws} \theta_s + \alpha_{wi});$$
  

$$\sin(-\theta_s + \alpha_{sj})(R_s - R_w \sin(k_{ws} \theta_s + \alpha_{wi})) - R_w$$
  

$$\sin \beta_s \cos(-\theta_s + \alpha_{sj})\sin(k_{ws} \theta_s + \alpha_{wi}); ] (i, j = 1, 2, 3, 4, 5, 6)$$
(1)

where i denotes the index of an outermost wire in a strand and j that of a strand in a rope.



Fig. 1. Schematic of a 6×7+FC wire rope and a friction lining block.

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